

Ionospheric *D* Region: VLF-Measured Electron Densities Compared With Rocket-Based FIRI-2018 Model

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Key Points:

- Very low frequency (VLF)-measured *D* region electron number densities are compared with the rocket-based Faraday-International Reference Ionosphere (FIRI)-2018 model over heights 60–90 km
- Average VLF-measured nighttime electron number densities agree very well with the rocket-based FIRI-2018 model (at heights 75–90 km)
- By day there is reasonable agreement at low latitudes (at ~60–75 km), but at high mid-latitudes FIRI lacks galactic cosmic ray effects

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Abstract Ground-based very low frequency (VLF) radio propagation in the Earth-ionosphere waveguide has enabled extensive electron number densities in the *D* region of the Earth's ionosphere to be determined, by day typically below heights of 70–80 km and by night in the height range ~75–90 km. Many rocket-based electron density measurements have also been reported in the literature from ~60 km upwards using current probes, and radio propagation at a few MHz between the rocket and ground. Recently these rocket measurements have been summarized, and supplemented with *D* region production-loss modeling, giving rise to a near global model named FIRI-2018 (Faraday-International Reference Ionosphere) which provides electron number densities as functions of height, latitude (<60°), solar zenith angle and F10.7 cm solar flux. These rocket-based electron density values are here compared with corresponding values from VLF measurements, by day at a low-latitude (~20°) and a high mid-latitude (~55°), and by night mainly at mid-latitudes. At night the average agreement (over 75–90 km) is remarkably good. By day, at low latitude the agreement is also fairly good (in the common height range ~60–75 km), with the changes with solar zenith angle being moderately comparable. For daytime high mid-latitudes, the agreement is less satisfactory, particularly at the lowest common altitudes, with the VLF measurements showing the expected effects of cosmic rays much more than the rocket-based values. Overall, we find that the *D* region description in the FIRI-2018 model is a significant advance on the earlier International Reference Ionosphere (IRI-2016) model.

1. Introduction

The *D* region is the lowest part of the Earth's ionosphere, found at heights below ~90 km. In quiet times, at non-polar latitudes, its free electrons normally extend down to ~75 km by night and down to ~55 km by day. Production of electrons by day is typically dominated by ionization of the minor atmospheric constituent nitric oxide (NO) by direct solar Lyman- α radiation above ~65–70 km, and by cosmic galactic rays ionizing all neutral constituents below this height. At night, the most significant ionizing source is likely to be the indirect solar Lyman- α reradiated by the neutral hydrogen in the Earth's geocorona (Banks & Kockarts, 1973). Production is essentially balanced by recombination and loss processes. However, these latter processes are not yet sufficiently well understood to enable calculation of reliable absolute electron number densities, and so measurements are needed to create empirical *D* region ionization models. At these low (*D* region) heights, satellite measurements are not practical because the air density is too high creating too much drag, while the electron number densities are normally too low (10^6 – 10^9 m⁻³) to allow sufficient reflection with a conventional ionosonde or an incoherent scatter radar. Two techniques have dominated experimental electron number density measurements in the *D* region at least at non-polar latitudes: (a) ground-based very low frequency (VLF) radio propagation and (b) in situ rocket measurements.

The VLF radio propagation technique has typically used frequencies in the range ~10–40 kHz radiated by powerful (and expensive) ground-based transmitters with very large antennas run by the local military (to communicate with their submarines near world-wide). The receivers are normally rather inexpensive with small (~1–10 m) antennas recording amplitude and phase often continuously, and sometimes are fully portable. The propagation used is in the Earth-ionosphere waveguide where the lower boundary is the ground, or preferably the sea, while the upper boundary is the *D* region of the ionosphere from which the VLF radio waves undergo partial reflection. While the propagation paths over the surface of the Earth can be several thousand km or more, shorter paths of a few hundred km are often more suitable for studying the *D* region in a desired location (e.g., low latitude, mid-latitude, etc.). In particular, at a range of ~300–400 km from the transmitter, there is often a “modal-minimum” where the received amplitude has a marked minimum compared with neighboring ranges;

this can be thought of either as waveguide modes interfering or as due to the ground wave from the transmitter interfering with (mainly) the first hop of the transmitter wave reflecting from the D region (e.g., Watt, 1967).

In this short path VLF situation, two different sub-techniques have been used. In the first of these, the local geography allowed amplitude and phase to be measured as functions of range over several tens of km near the modal minimum (Thomson et al., 2017). In the second sub-technique, the amplitude and phase were measured both near ($< \sim 100$ km) the transmitter (where the ground wave is dominant and the ionospheric reflection is near negligible) and at ~ 300 km from the transmitter (near the modal minimum), where there is good sensitivity to the ionospherically reflected signal (e.g., Thomson et al., 2014). Both these sub-techniques avoid needing to otherwise know the transmitter's radiated power and, in particular, needing to otherwise know the radiated phase of the signal at the transmitter. They are both effectively measuring the changes in phase and amplitude with distance which are then compared with the corresponding calculated changes from the US Navy modal codes, ModeFinder (Morfitt & Shellman, 1976) or Long-Wavelength Propagation Capability (LWPC; Ferguson & Snyder, 1990; see also Ferguson, 1998) using a range of possible model electron number densities versus height to find the one that best fits the experimental observations. Both phase and amplitude comparisons are normally essential for VLF ionospheric measurements to avoid ambiguity and achieve good accuracy and reliability.

In contrast, the rocket measuring technique (e.g., Friedrich et al., 2018) involves transmitting a medium frequency (MF) linearly polarized radio wave of a few MHz usually (nowadays) from the ground for reception on the rocket where its received polarization is measured as the height of the rocket changes. This is aided by the rocket spinning at a slow but known rate about its vertical axis. The transmitted linearly polarized wave can be thought of as being made up of two equal right and left circularly polarized waves. Once these enter the plasma (i.e., where there are free electrons in the ionosphere) these two travel as separate modes with different phase velocities and attenuations. This means that when they “recombine” (i.e., are measured together on the rocket) their combined (quasi-linear) polarization will be found to have rotated by an amount related to the amount of plasma the waves have passed through (Faraday rotation) while the extent of the ellipticity of the received polarization is a measure of the relative attenuation of the two polarization modes (right and left) which in turn is also a measure of the amount of plasma passed through. Thus the rotating polarization and the changing relative amplitudes measured on the rocket are continuously measuring the plasma changes as the rocket rises (or falls) and so effectively measuring the refractive index and hence electron number density at the current height of the rocket.

Both the ground-based VLF-radio technique and the rocket-based MF-radio technique require the (appropriate) electron-neutral collision frequency at each height to determine the electron number density at that height. This requires the neutral number densities at each height, particularly the nitrogen (N_2) number density since N_2 is $\sim 78\%$ of the neutral atmosphere at and below D region heights. Friedrich et al. (2018), when reporting on their rocket-based FIRI-2018 electron density model, used the NRLMSISE-00 model (Picone et al., 2002), being the latest neutral atmosphere model then available. For the VLF data used here, we use the newer NRLMSIS 2.0 model (Emmert et al., 2020). VLF-determined D region results are normally reported using the “Wait” (e.g., Wait & Spies, 1964) height and sharpness parameters H' and β (e.g., Thomson et al., 2014, 2017) to describe the variation with height of the electron number density in the D region at the time and location of the VLF measurements. This has been done because these two parameters can be determined largely independently of the neutral density and collision frequency height profiles assumed at the time. Later, as here, when the best values of electron number density at the time and place of the VLF measurements are required (Thomson et al., 2018, 2021) these values can be obtained retrospectively from the measured H' and β , together with the most recent best estimates of the neutral atmospheric density and collision frequency.

In determining the electron densities from the VLF observations here we use the same formula for the monoenergetic electron-neutral collision frequency as used in developing FIRI-2018—that is, $\nu_m = Kp$ where p is the pressure at height, h , and $K = 6.4 \times 10^5$, in SI units (Friedrich & Torkar, 1983), with $p = NkT$ being determined from the appropriate NRL atmospheric model, where N is the (total) neutral number density, k is Boltzmann's constant, and T is the neutral temperature. However, as discussed by Thomson et al. (2018), the available VLF propagation codes, such as ModeFinder and LWPC, use the Appleton-Hartree equations which assume the electron-neutral collision cross section is independent of velocity whereas the rocket-based electron number densities in FIRI-2018 (Friedrich et al., 2018) used the more recent Sen-Wyller equations where the electron-neutral collision cross section is taken as proportional to the electron velocity. This means that the above $\nu_m = Kp$ can be used directly in the Sen-Wyller equations but requires modification before use in converting the VLF-measured H' and β values

into electron densities. As also discussed by Thomson et al. (2018), Deeks (1966), using both Appleton-Hartree and Sen-Wyller formulations at VLF, calculated a series of adjustment factors between 1.5 and 2.5, starting with 1.5 low in the D region and increasing monotonically to 2.5 high in the D region by which ν_m should be increased to generate an effective ν for use in Appleton-Hartree formulations to give as near as possible the same results as the Sen-Wyller formulation. Thus, as recommended by Deeks (1966), we have here used this appropriate factor at each height as plotted in his Figure 1a.

As discussed in Thomson et al. (2018), Wait and Spies (1964) defined the parameter $\omega_r = \omega_o^2/\nu$ where ω_o is the angular (electron) plasma frequency, and ν is an appropriate effective collision frequency such as the “effective ν ” described toward the end of the previous paragraph. The electron number density is thus given by $N_e \approx \nu\omega_r/3,183$ (since $e^2/\epsilon_0 m_e \approx 3,183$), and ω_r was taken to vary with height, h , as $\omega_r = 2.5 \times 10^5 \exp(h-H')\beta$ rad/s thus defining H' as the (reference) height at which $\omega_r = 2.5 \times 10^5$ rad/s, and β as a (near) constant with height, but dependent, even in quiet times, on latitude, time of day, and solar cycle.

For our main comparisons here between electron number densities from FIRI-2018 and those from VLF measurements, we have chosen (1) daytime at a low latitude in the Hawaiian Islands, latitude $\sim 20^\circ\text{N}$, using transmitter NPM on 21.4 kHz, (2) daytime at a high mid-latitude from the 23.4 kHz transmitter DHO in north Germany (a) along the west coast of Denmark's Jutland peninsula and (b) across the North Sea to Eskdalemuir, Scotland, both at latitudes $\sim 55^\circ\text{N}$, and (3) nighttime at mid-latitudes, all at relatively quiet times. The VLF phase and amplitude measurements were made using both a (hand-held) portable loop receiver for daytime observations, and at least one (usually more) separate fixed receivers recording continuously. All phases were referenced to one-second pulses from the satellite Global Positioning System (GPS). The fixed recorders enabled any phase or amplitude changes which occurred at the transmitters to be corrected for. Details can be found in Thomson et al. (2014, 2017, 2018, 2021) and references therein. No polar comparisons are made because FIRI-2018 does not extend higher than $\sim 60^\circ$ latitude. Recently Siskind et al. (2018) compared some daytime VLF-derived electron number densities with those from rocket soundings but did not make any detailed comparisons with the FIRI-2018 model.

FIRI-2018 is available as 1,980 profiles of electron number densities between 60 and 150 km altitude, including 11 solar zenith angles $0^\circ, 30^\circ, 45^\circ, 60^\circ, 75^\circ, 80^\circ, 85^\circ, 90^\circ, 95^\circ, 100^\circ, 130^\circ$, 5 latitudes $0^\circ, 15^\circ, 30^\circ, 45^\circ, 60^\circ$, three solar activities, $F_{10.7} = 75, 130, \text{ and } 200$ sfu, and for the middle of each calendar month (i.e., a total of $11 \times 5 \times 3 \times 12 = 1,980$ profiles). These model profiles are available from <https://figshare.com/s/357cb-03b3e5bed649bbc> (Friedrich et al., 2018) or <https://figshare.com/search?q=FIRI-2018>. The profiles are actually provided from 55 to 150 km, at 1 km intervals, but Friedrich et al. (2018) advise that they consider them reliable only above 60 km (and for electron densities larger than 10^6 m^{-3}). Recently Xu et al. (2021) have reported on parameterizing (fitting) each of these 1,980 FIRI-2018 electron density profiles with Wait and Spies (H' and β) parameters. These could then potentially be used to compare with the corresponding VLF-derived Wait and Spies parameters. However, we prefer here to compare electron density height-profiles from VLF with those from FIRI-2018.

2. Daytime Low-Latitude Comparisons: VLF and FIRI-2018

The VLF results used here are those from August 2012 reported by Thomson et al. (2014).

2.1. Comparisons Near Midday at Solar Zenith Angle $\sim 10^\circ$

The VLF measurements made in August 2012 in Hawaii (latitude 20.5°N) described in Thomson et al. (2014) showed that near noon when the solar zenith angle was $\sim 10^\circ$ that the D region was characterized by the Wait parameters $H' = 69.3$ km and $\beta = 0.49 \text{ km}^{-1}$. These parameters lead to the electron number densities shown by the (nearly) straight black line with large open “+” plot symbols in the upper panel of Figure 1. These were calculated using the formula $N_e = \nu\omega_r/3,183$ from Section 1 above where the (Deeks-adjusted) collision frequency, ν , was derived, as also explained in Section 1, from the NRLMSIS 2.0 neutral densities for the time and location of the measurements, thus representing the best estimate of the electron densities from these VLF observations (i.e., using $H' = 69.3$ km and $\beta = 0.49 \text{ km}^{-1}$).

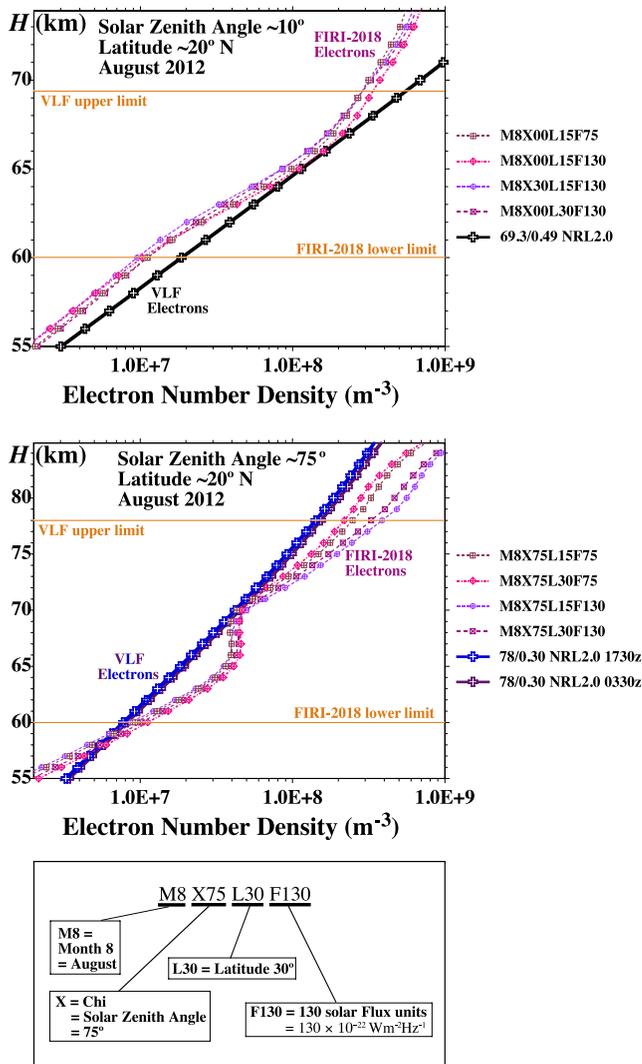


Figure 1. Low latitude ($\sim 20^\circ\text{N}$) comparisons of VLF-derived (using 21.4 kHz) and Faraday-International Reference Ionosphere (FIRI)-2018 electron number densities as functions of height. (Top panel) At solar zenith angles $\sim 10^\circ$. (Middle panel) At solar zenith angles $\sim 75^\circ$. (Bottom panel) Illustrates FIRI-2018 plot symbol labels such as “M8X75L30F130.” VLF plot symbol labels such as “69.3/0.49 NRL2.0” mean $H' = 69.3$ km and $\beta = 0.49$ km^{-1} , and NRL2.0 means the NRLMSIS 2.0 neutral atmosphere model was used to determine the collision frequencies.

The four curves to the left of the black VLF line are the appropriate FIRI-2018 (Friedrich et al., 2018) electron number density profiles for comparison, for example, in the label “M8X30L15F130”, M8 means Month 8 = August, X30 means a solar zenith angle of 30° (“Chi” in FIRI-2018), L15 means latitude 15° , F130 means $F10.7 = 130$ solar flux units, etc. So, three of these four curves are for solar zenith angles of 0° which are likely to be the nearest match to the actual 10° , and the other at 30° is for comparison. The upper orange horizontal line is at $H' = 69.3$ km, above which height the VLF measurements are essentially insensitive to the electron number density (e.g., Siskind et al., 2018). This insensitivity above height, H' , applies at least as low as $H' \approx 58$ km at the peak of an X6 flare (McRae & Thomson, 2004). The lower orange horizontal line at 60 km is the height above which Friedrich et al. (2018) state that their FIRI-2018 electron number densities are considered “reliable.” As can be seen the agreement between the VLF and FIRI-2018 electron number densities is good near the center of the height range (64–67 km) where both sets of densities are valid, but at both the lower and upper limits of the mutual height range validity (~ 60 and 69 km) the FIRI 2018 values are lower than the VLF values by a not insignificant factor of ~ 1.6 . This will be further considered in Section 2.3 below. In particular, below heights of ~ 62 km, the FIRI and VLF derived electron number densities run nearly parallel with a height difference of ~ 1 km meaning that the FIRI-2018 values are very similar to a “Wait 70/0.5” profile ($H' = 70$ km and $\beta = 0.5$ km^{-1}); this will be commented on again in the next subsection and in Section 3.

2.2. Comparisons Early/Late in the Day at Solar Zenith Angle $\sim 75^\circ$

Thomson et al. (2014) showed that, when the solar zenith angle was $\sim 75^\circ$, the low latitude (Hawaiian) D region was characterized by the Wait parameters $H' = 78.6$ km and $\beta = 0.29$ km^{-1} in the morning (~ 1730 UT) and $H' = 77.4$ km and $\beta = 0.31$ km^{-1} in the afternoon (~ 0330 UT), giving an am/pm average of $H' = 78.0$ km, and $\beta = 0.30$ km^{-1} for solar zenith angle (SZA) = 75° . These average values lead to the electron number densities shown by the two (nearly) straight lines (on the left) in the middle panel of Figure 1. These blue and red-brown lines with open “+” plot symbols show the electron number densities using the NRLMSIS 2.0 Deeks-adjusted collision frequencies at 1730 and 0330 UT respectively. In this panel, the four curves to the right are the FIRI-2018 electron number densities in August at a solar zenith angle of 75° at the two latitudes of 15° and 30° and at the two F10.7 values of 75 and 130 sfu (similar to the top panel). Also, rather similar to the top panel, the FIRI electron number densities at the lowest heights again run from $\sim 2 \times 10^6$ m^{-3} at 55 km to $\sim 2 \times 10^7$ m^{-3} at 62 km, that is, are rather similar to a “Wait 70/0.5” profile ($H' = 70$ km and $\beta = 0.5$ km^{-1}); this will be commented on again in Section 3 below.

The lower orange horizontal line remains at 60 km in the lower panel because this is the Friedrich et al. (2018) FIRI-2018 lowest height of “reliability”; the upper orange line is now at 78 km because this is $\sim H'$ at the SZA of 75° and is thus the greatest height at which the VLF measurements are sensitive to the electron number density. Overall, the agreement between VLF and FIRI-2018 at SZA = 75° is neither good nor bad; it is, however, quite passable. More discussion on this is given in the next sub-section.

2.3. VLF Phase and Amplitude Calculated With FIRI Compared With Observations

Instead of entering values of H' and β , ModeFinder allows entry of an electron density versus height profile and a collision frequency versus height profile which are then used to find the appropriate VLF modes under

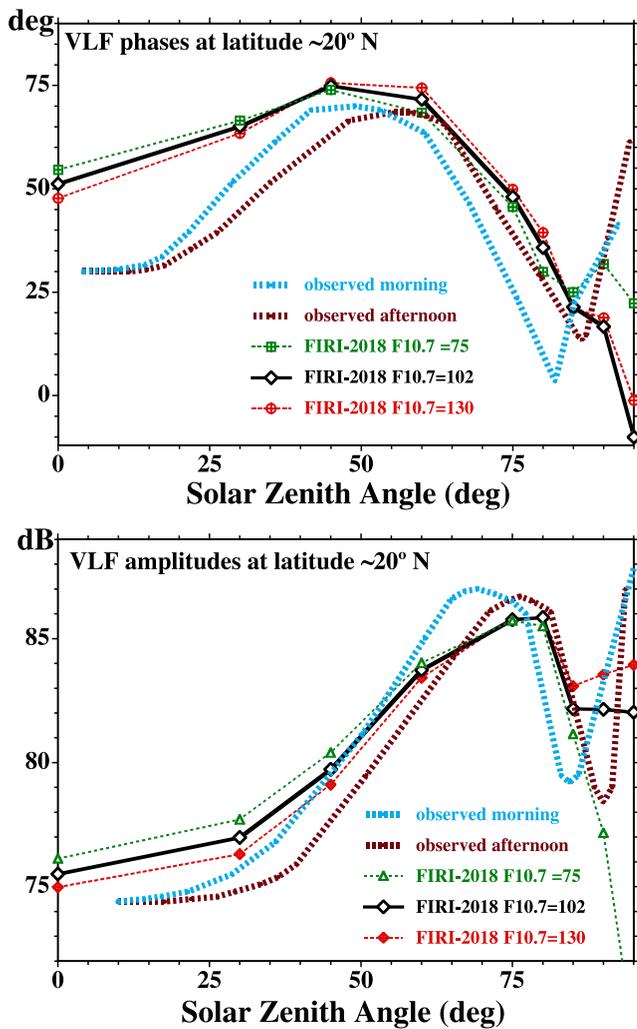


Figure 2. VLF phases and amplitudes, as functions of solar zenith angle, observed and calculated for the short (306-km), nearly all-sea, low latitude path from NPM (21.4 kHz) in Hawaii. The thick dashed lines are the observations (from Thomson et al., 2014), light blue for morning and dark red for afternoon. The green, black and red data points, joined by lines of the same colors, were calculated with ModeFinder using the corresponding Faraday-International Reference Ionosphere (FIRI)-2018 profiles (interpolated for latitude and F10.7 as appropriate—see text for details).

these supplied *D* region characteristics; from these, ModeFinder calculates the amplitude and phase of the signal at the receiver (in a very similar way to when H' and β are supplied as inputs). The FIRI-2018 electron density profiles for August for latitudes 0°, 15°, and 30° were used here. A simple quadratic interpolation was used at each height (1 km intervals) to find the “FIRI-2018” profile for latitude 20.5° (the latitude of the NPM-Hawaii path). This was done for each of the nine (FIRI-2018 tabulated) solar zenith angles between 0° and 95°, and for each of the two tabulated solar activities, F10.7 = 75 and 130 sfu. The collision frequency profile was generated from the NRLMSIS 2.0 neutral density profile, at latitude 20.5°, using the Deeks adjustment as above. The midday collision frequency profile so generated was used for all 9 solar zenith angles because, at this low latitude, the effect on the calculated phases and amplitudes of using more exact timing was found to be negligible. The results are plotted in Figure 2 with the phases shown in the upper panel and the amplitudes shown in the lower panel. The means of the two F10.7 values (75 and 130 sfu) are also shown, as F10.7 = 102 sfu which is close to the actual value of F10.7 at the time of the VLF observations (August 2012). Also shown are the best-fit curves for the VLF phases and amplitudes versus solar zenith angle actually observed for both morning and afternoon from Thomson et al. (2014). The FIRI-2018 model does not differentiate between morning and afternoon; it specifies only the solar zenith angle.

As can be seen, in both the phase and amplitude panels of Figure 2, there is clearly some agreement between the observed values as functions of solar zenith angle and those calculated from the FIRI-2018 model at least in general shape, apart from the highest solar zenith angles (90°–95°). For low solar zenith angles (~10°, near midday), there is a fairly significant difference with the FIRI-2018 calculated phase value of ~55° being higher than that observed by ~25°. From the upper panel of Figure 2, it can be seen the observed phase does not reach 55° until the solar zenith angle is 35°–40°. From Thomson et al. (2014), using their Figure 3 or 5, this implies the FIRI-2018 profile is ~0.8 km higher than the VLF profile ($H' = 69.3$ km) agreeing, at least approximately, with the apparent relative heights of the VLF and FIRI-2018 profiles shown in the upper panel of Figure 1 here.

3. Daytime High Mid-Latitude Comparisons: VLF and FIRI-2018

The VLF results used in this section are from July 2015 as reported by Thomson et al. (2017), using two closely located, high mid-latitude paths from DHO in north Germany (a) along the west coast of Jutland, Denmark and (b) across the North Sea to Eskdalemuir in Scotland.

3.1. Comparisons Near Midday at Solar Zenith Angle ~33°

The VLF measurements made in July 2015 from north Germany to Denmark (latitude ~54.5°N) of Thomson et al. (2017) showed, near noon when the solar zenith angle was ~33°, that the *D* region was characterized by the Wait parameters $H' = 72.8$ km and $\beta = 0.345$ km⁻¹. These parameters lead to the electron number densities shown by the (nearly) straight black line with large open “+” plot symbols in the upper panel of Figure 3, calculated using the formula $N_e = \nu\omega/3,183$ with the (Deeks-adjusted) collision frequency, ν , being derived from NRLMSIS 2.0 neutral densities for the time and place of the measurements. To the left, the four colored curves show for comparison the corresponding FIRI-2018 electron number density profiles for latitudes near 55°N at solar zenith angles of 30° and 45°, in July 2015 when F10.7 = 130 sfu.

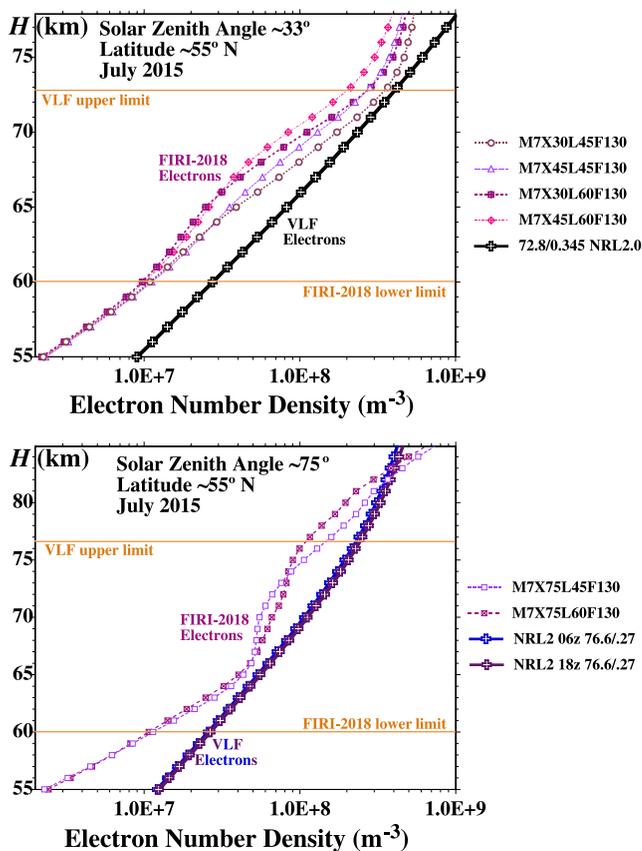


Figure 3. High mid-latitude ($\sim 55^\circ\text{N}$) comparisons of VLF-derived (using 23.4 kHz) and Faraday-International Reference Ionosphere (FIRI)-2018 electron number densities as functions of height. Plot symbol labels are explained in Figure 1 above. (Top panel) Comparisons at solar zenith angles $\sim 33^\circ$. (Bottom panel) Comparisons at solar zenith angles $\sim 75^\circ$.

At the greatest heights at which the VLF measurements are sensitive (~ 72.8 km, $\sim H'$), it can be seen, in the upper panel of Figure 3, that the agreement is quite good between the VLF-derived electron densities (at $\sim 55^\circ$ latitude) and the FIRI-2018 model at a solar zenith angle of $\sim 33^\circ$ and a latitude of 45° (FIRI data “M7X30L45F130”); however, as can also be seen, the FIRI model gives somewhat lower electron densities at the higher latitude of 60° . Although this overall agreement at heights near 72 km is quite passable, it is also clear that toward the lowest heights (~ 60 km) at which the FIRI model is considered reliable (Friedrich et al., 2018), its electron densities are significantly lower (by a factor of nearly 3) than the VLF-derived values. In particular, the FIRI electron densities at a height of 60 km are nearly the same ($\sim 1.0 \times 10^7 \text{ m}^{-3}$) at latitude 20° (as shown in Figure 1) and at latitudes $\sim 55^\circ$ shown here, whereas the ionization due to galactic cosmic rays is expected to be significantly higher at these higher latitudes (e.g., Heaps, 1978).

3.2. Comparisons Early/Late in the Day at Solar Zenith Angle $\sim 75^\circ$

The same set of VLF measurements reported by Thomson et al. (2017), from July 2015 at latitude $\sim 54.5^\circ\text{N}$, in Section 3.1 above, but using the path DHO to Eskdalemuir instead of DHO to Denmark, showed that at solar zenith angles of $\sim 75^\circ$ the D region was characterized by $H' = 76.6$ km and $\beta = 0.27 \text{ km}^{-1}$; these parameters lead to the electron number densities shown in the lower panel of Figure 3 by the two fairly similar nearly straight lines (to the right), with open “+” plot symbols, again using the formula $N_e = \nu\omega/3,183$ as above. Both these lines use NRLMSIS 2.0 but the blue line is for 6 UT (morning) while the red-brown line is for 18 UT (afternoon) to illustrate the small but noticeable difference. The FIRI model does not discriminate between morning and afternoon so such differences are not pursued further here. Again it needs to be noted that all the FIRI profiles in all four panels of Figures 1 and 3 are essentially the same at ~ 60 km and below, independent of latitude.

As can be seen in this lower panel of Figure 3, at the greatest heights at which the VLF measurements are sensitive ($H' \sim 76$ km) the FIRI-2018 electron densities, for a solar zenith angle of 75° are low (plot symbol squares, $\sim 1.1 \times 10^8 \text{ m}^{-3}$) by a factor of ~ 2 compared with the VLF results. However, if FIRI electron densities at a solar zenith angle of 60° (not shown) were used instead the agreement would be much better, being then just marginally higher ($\sim 2.5 \times 10^8 \text{ m}^{-3}$ at 76 km) than the VLF results. This might possibly be due to a scarcity of rocket profiles (in solar zenith angle and latitude).

Again, as in Section 3.1, it is clear that toward the lowest heights (~ 60 km) at which the FIRI model is considered reliable its electron densities are lower (by a factor ~ 2.5) than the VLF-derived values. Also, the FIRI electron densities at a height of 60 km are nearly the same ($\sim 1.0 \times 10^7 \text{ m}^{-3}$) at latitude 20° (as shown in Figure 1) and at latitudes $\sim 55^\circ$ shown here, whereas the ionization due to galactic cosmic rays is expected to be significantly higher at these higher latitudes (e.g., Heaps, 1978). Siskind et al. (2018) have also pointed out that sounding rocket profiles have tended to fail to show the increasing electron densities with increasing latitude (at low altitudes) expected from the well-known, corresponding increases in galactic cosmic rays.

3.3. VLF Phase and Amplitude Calculated With FIRI Compared With Observations

Section 2.3 above used ModeFinder to calculate VLF phases and amplitudes on our low latitude path as functions of solar zenith angle using appropriate low latitude FIRI-2018 electron density profiles, and compared these with the corresponding VLF observations in Figure 2. Here, in Figure 4, similar comparisons are made but for the high mid-latitude path DHO to Eskdalemuir; as can be seen, the agreement between the VLF observations (from Thomson et al., 2017) and the calculations using the appropriate FIRI profiles for both the phases (upper panel) and the amplitudes (lower panel) is relatively poor. A significant contributor to this poor level of agreement is

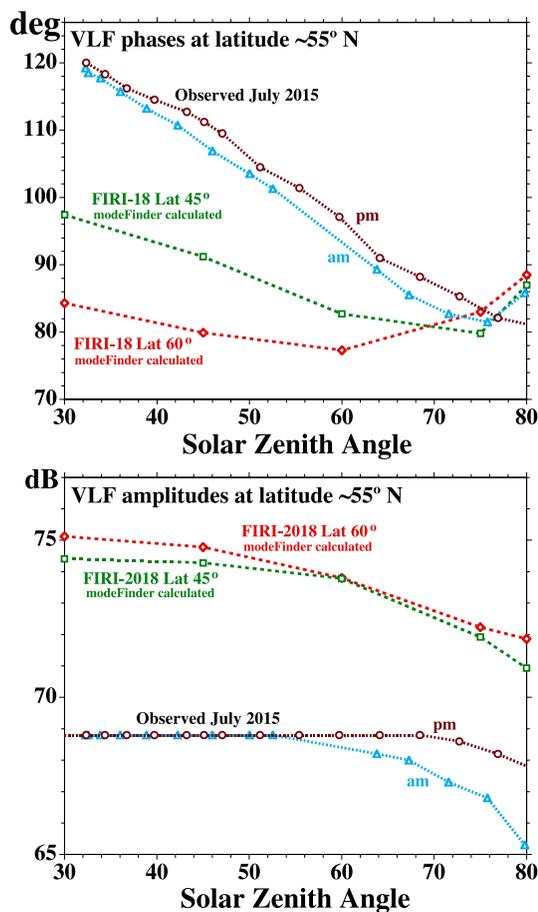


Figure 4. VLF phases and amplitudes, as functions of solar zenith, observed and calculated for the 748-km nearly all-sea, high mid-latitude path, from DHO (23.4 kHz) to Eskdalemuir. The observations (from Thomson et al., 2017) are shown in light blue for morning and dark red for afternoon. The green and bright red data points, joined by lines of the same colors, were calculated with ModeFinder using the corresponding Faraday-International Reference Ionosphere (FIRI)-2018 profiles.

likely to be related to the FIRI-2018 model not fully recognizing a greater number of galactic cosmic ray generated electrons below heights of ~ 70 km at these relatively high latitudes. While these cosmic ray generated electrons will not themselves vary much with solar zenith angle (during daytime), the electron density above ~ 70 km is solar generated (mainly Lyman- α) and so solar zenith angle dependent. Whether this effect actually shows up in the VLF phases or amplitudes depends on the exact nature of the path, including its length. Here, as the VLF observations show, the VLF phase is quite strongly solar zenith angle dependent while the amplitude is not.

4. High Mid-Latitude VLF Electron Densities Below ~ 60 km Height

Comparing the two panels of Figure 3, it can be seen that at heights below ~ 60 km the VLF-derived electron number densities near noon are lower than early or late in the day; for example, at 55 km specifically, the VLF-derived electron number density at a solar zenith angle of 75° shows $\sim 1.2 \times 10^7 \text{ m}^{-3}$ in the lower panel while at a solar zenith angle of $\sim 33^\circ$, in the upper panel, it shows as $\sim 0.9 \times 10^7 \text{ m}^{-3}$, that is, $\sim 30\%$ lower. This seems unlikely because galactic cosmic rays, which are the principal ionizing source at these heights, are not solar zenith angle dependent. Another possibility could be the photo-detachment, by visible light at sunrise, of electrons (e.g., Kazil et al., 2003; Ogawa & Shimazaki, 1975) from negative ions generated during the night from galactic cosmic rays, but this effect is likely to have faded when the sun has risen to a solar zenith of 75° (and also does not occur at dusk).

To assess whether this difference is significant in terms of likely errors, the noon profile of $H' = 72.8$ km and $\beta = 0.345 \text{ km}^{-1}$ as determined by Thomson et al. (2017) was compared with a similar profile with two parts, $H' = 76.6$ km and $\beta = 0.27 \text{ km}^{-1}$ (as for a solar zenith angle of 75°) below 60 km and $H' = 72.6$ km and $\beta = 0.355 \text{ km}^{-1}$ above 60 km as shown in the upper panel of Figure 5. The lower two panels of Figure 5 here are taken from Thomson et al. (2017) showing the measured (data points) and calculated (lines) phases and amplitudes versus distance, with the new (ModeFinder) calculated (thick green) lines superposed for the new two-part profile with $H' = 72.6$ km and $\beta = 0.35 \text{ km}^{-1}$ (for phase) and $\beta = 0.36 \text{ km}^{-1}$ (for amplitude) above 60 km and $H' = 76.6$ km and $\beta = 0.27 \text{ km}^{-1}$ below 60 km. As

can be seen the (previous) best fit black lines ($H' = 72.8$ km and $\beta = 0.345 \text{ km}^{-1}$) fit the data points just as well as the new green lines. The change of just 0.2 km in H' and 0.01 km^{-1} in β is within the originally estimated experimental error and so is not of much significance.

Of course, instead of adjusting the noon profile below 60 km to the 75° profile, the 75° profile could have been adjusted instead (e.g., by similarly making into two parts) so that, below 60 km it matched the original, $H' = 72.8$ km and $\beta = 0.345 \text{ km}^{-1}$, noon profile, or indeed both noon and 75° profiles could have each been even more slightly adjusted (in two parts) so that they matched below 60 km. This would involve <0.2 km in H' and $<0.01 \text{ km}^{-1}$ in β and so be of near negligible significance. Also new in Figure 5 are orange “+” symbols, identical to the new green profile but with it cut-off at a height just below 57 km, showing that the VLF technique is here sensitive down to heights at least as low as 57 km.

5. Nighttime D Region Comparisons: VLF and FIRI-2018

The nighttime VLF-derived electron number densities used in this section come from long, nearly all-sea paths over a range of latitudes, but excluding both the polar regions ($>\sim 60^\circ$ latitude, as does FIRI-2018) and the equatorial regions (say $<\sim 15^\circ$ geomagnetic latitude), mainly from Thomson et al. (2007) but also supported by Thomson

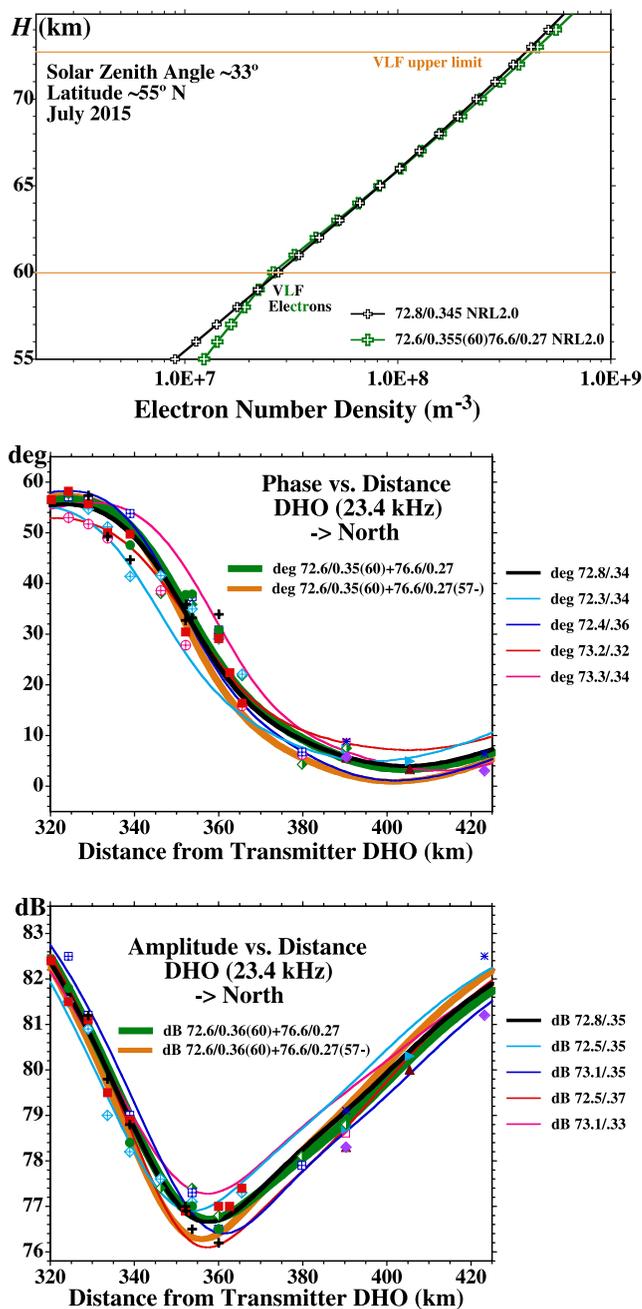


Figure 5. Assessing small shape changes in the noon, high mid-latitude, VLF-derived electron density profile. (Top panel) The original (black) $H' = 72.8$ km and $\beta = 0.345$ km⁻¹ profile from Thomson et al. (2017) together with the new (green) two-part profile as labeled. (Lower two panels) ModeFinder calculations (green lines) for the new profile overlaid on the phase and amplitude observations for 5–17 July 2015 (data points), and calculations (lines) first presented in Thomson et al. (2017) for the high mid-latitude path at noon.

and McRae (2009). Thomson et al. (2007) found $H' = 85.1 \pm 0.4$ km and $\beta = 0.63 \pm 0.04$ km⁻¹ on average, predominantly in summer and under conditions nearer solar minimum than solar maximum. The FIRI-2018 profiles compared in this section are those with solar zenith angles of 130° (the highest solar zenith angle reported in FIRI-2018) which correspond with being full night. The next highest solar zenith angle available in FIRI-2018 was 100°; it was decided not to use these 100° values for comparisons because they would be only just into full darkness and the VLF data reported in Thomson et al. (2007) was wholly or mainly with solar zenith angles greater than 100°, that is, with the ionosphere well settled into full night.

Figure 6 compares the VLF-derived nighttime electron number densities with those from FIRI-2018. The thick black solid (nearly) straight line with the open “+” plot symbols shows the VLF-derived electron number densities for $H' = 85.0$ km and $\beta = 0.65$ km⁻¹ using collision frequencies derived (as above) from the NRLMSIS 2.0 model (nominally for July at 45° latitude, 0° longitude, and 0 UT, though this was not critical). The solid blue line is similarly calculated for $H' = 85.0$ km and $\beta = 0.60$ km⁻¹. The FIRI-2018 electron number densities are shown with the dot/dashed curves with the plot-symbol labeling similar to that used above except that the “X130” (indicating a solar zenith angle of 130°) is not shown in the labels because they are all for 130°. So, for example, “M7L45F130” means Month 7 (i.e., July), Latitude 45°, solar Flux 130 sfu. The VLF data correspond to solar activity roughly midway between F10.7 solar fluxes of 75–130 sfu. Friedrich et al. (2018) have indicated that FIRI-2018 electron number densities below 1×10^6 m⁻³ are not considered reliable; so these have been included only for completeness. At night (when $H' \approx 85$ km), ModeFinder calculations show that the VLF technique is sensitive down to heights ~ 75 km, and up to heights of ~ 90 km, that is, somewhat above H' (the daytime sensitivity upper height limit mentioned in Section 2.1).

In Figure 6, it can be seen that the agreement between the night VLF-derived D region electron number densities and those from the (rocket-based) FIRI-2018 model is remarkably good, especially at latitudes near 45°. It must be emphasized that these are all averaged values; the night D region is rather variable (e.g., Thomson et al., 2007) in time and space (compared with the daytime D region) whether observed by ground-based VLF or by rocket-borne MF radio. In particular, as pointed out by Friedrich et al. (2018), the actual individual nighttime electron density profiles as measured by rockets at a particular place and time often show “ledges” where the electron density increases much more rapidly with height than shown here in Figure 6, but these ledges occur (randomly) at different heights, typically between ~ 80 and 90 km, and so, when averaged, result in the profiles in Figure 6. At night, galactic cosmic rays generate free electrons at similar rates as by day; these electrons are then similarly rapidly removed by attachment to neutral molecular oxygen molecules: $e^- + O_2 \rightarrow O_2^-$. By day, photons of (visible) sunlight immediately release these electrons again from the negative ions, resulting in cosmic ray generated free electrons dominating below heights of 65–70 km as noted above. In darkness, this last step cannot occur, resulting in very few free electrons from cosmic rays at night (Thomson et al., 2021; see also Banks & Kockarts, 1973).

Cummer et al. (1998) and Cheng et al. (2006) used VLF propagation from natural lightning (sferics) to determine nighttime D region characteristics over the USA at latitudes $\sim 35^\circ$ – 37° with both also finding $H' \approx 85.0$ km but both finding significantly lower values of β with Cummer et al. finding $\beta \approx 0.50$ km⁻¹ and Cheng et al. $\beta \approx 0.45$ km⁻¹.

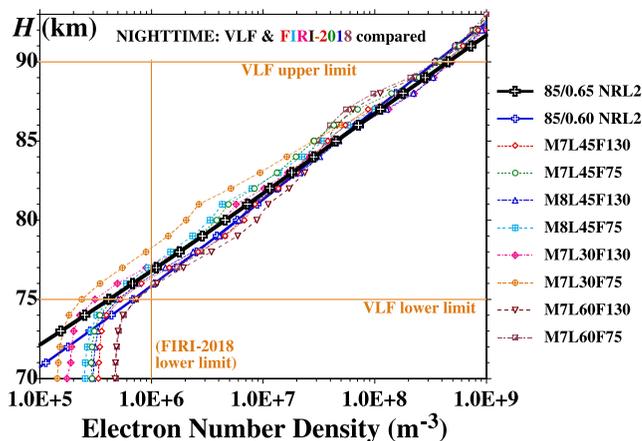


Figure 6. Nighttime, broadly mid-latitude comparisons of VLF-derived and Faraday-International Reference Ionosphere (FIRI)-2018 electron number densities as functions of height. Plot symbol labels are explained in Figure 1. The VLF technique is not sensitive here above heights of ~ 90 km. The FIRI electron densities are not considered reliable below densities of 10^6 m^{-3} .

6. Neutral Atmosphere Collision Effects on VLF Electron Densities

The VLF-derived electron number densities presented in the previous sections here were calculated from their measured H' and β values using the NRLMSIS 2.0 neutral density model (Emmert et al., 2020) as explained in Section 1 above. These VLF-derived electron number densities are now compared with (a) those similarly derived from the NRLMSISE-00 model (Picone et al., 2002) as used by Friedrich et al. (2018) for FIRI-2018, and (b) those derived using, as also discussed in Section 1, the “Wait” collision frequency (with collision cross section independent of velocity, and so with no Deeks adjustments) as incorporated in ModeFinder, that is, $\nu = 5.0 \exp(-0.15(h-70))$ MHz. In Figure 7, the top and middle panels make these comparisons for the low latitude case ($\sim 20^\circ\text{N}$) and the high mid-latitude case ($\sim 55^\circ\text{N}$) respectively, at the two appropriate daytime solar zenith angles (mid-day, 10° and 33° , respectively, and 75°), while the bottom panel is for night (at mid-latitudes).

It can be seen that the electron densities using NRLMSISE-00 are normally only very slightly (and so negligibly) larger than those using NRLMSIS 2.0. Only at both higher latitude ($\sim 55^\circ\text{N}$) and higher solar zenith angle (SAZ = 75°), and at daytime heights > 70 km, do the NRLMSISE-00 electron densities become appreciably, though still marginally, larger (by $\sim 10\%$) than those using NRLMSIS 2.0.

In contrast, while the “Wait”-derived electron densities nearly always deviate from the NRLMSIS 2.0 values more than those from NRLMSISE-00, none-the-less they do not deviate very much for both daytime low latitudes ($\sim 20^\circ\text{N}$) and nighttime mid-latitudes. However, as can be seen in the middle panel, the “Wait”-derived electron densities at the high mid-latitude of $\sim 55^\circ$ are typically significantly lower by up to factors of ~ 1.6 , compared with those derived from the more modern collision frequencies and atmospheric models.

7. Comparisons: IRI and VLF

The D region part of the FIRI-2018 model is considered a significant advance on the D region part of the current (2016) International Reference Ionosphere (IRI-2016) model (Bilitza et al., 2017). Here we compare the VLF electron number densities with the equivalent IRI-2016 profiles. In the D region, IRI-2016 uses IRI-95 which is based on an older and much smaller selection of typical rocket profiles such as those in Mechty et al. (1972) who also used Langmuir probes calibrated by Faraday and differential-absorption data. The IRI-2016 data used here came from the NASA web page https://ccmc.gsfc.nasa.gov/modelweb/models/iri2016_vitmo.php which also includes a second D region model “FT-2001” which is a much earlier version of FIRI (Friedrich & Torkar, 2001). Figure 8 shows comparisons between (a) the VLF-derived electron number densities presented here in Sections 2, 3, and 5, and (b) the corresponding IRI data. The top and middle panels are for mid-day at latitudes of $\sim 20^\circ$ and $\sim 55^\circ$, respectively, while the bottom panel is for night at mid-latitudes. As can be seen, a principal disadvantage of the IRI-2016 D region models is that their electron density profiles do not extend down as low as either the FIRI-2018 or the VLF profiles. By day (upper two panels) the IRI-2016 profiles extend down to ~ 65 km compared with down to ~ 60 km for FIRI-2018. This means they are essentially not registering electrons generated by galactic cosmic rays; it also largely precludes their use for VLF modeling such as for the phase and amplitude calculations performed using FIRI-2018 here in Sections 2.3 and 3.3. Similarly at night (bottom panel) it can be seen that the IRI-95 and FT-2001 profiles extend down only to 80 and 83 km respectively. Also, although the agreement between the VLF and FT-2001 profiles is good (down to just ~ 83 km), there is little agreement with IRI-95 profile. In the top panel, even where the red, IRI-95 profile crosses the black VLF line, the IRI-95 line has a slope corresponding to $\beta = \sim 0.40 \text{ km}^{-1}$ while the VLF observations give $\beta = \sim 0.49 \text{ km}^{-1}$ which, as can be seen in Figure 3 of Thomson et al. (2014), is quite a large difference (corresponding to 4–5 standard deviations or 4–5 dB in amplitude).

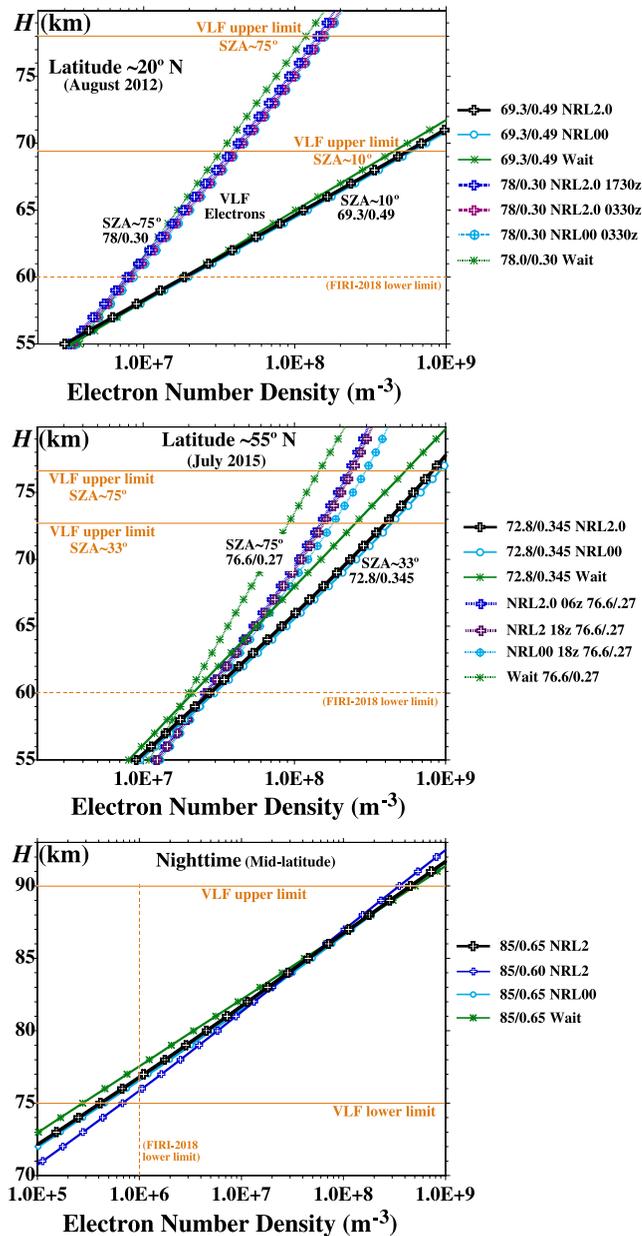


Figure 7. VLF-derived electron number densities compared depending on the neutral atmospheric model and so the electron-neutral collision frequencies. SZA = Solar Zenith Angle. (top panel) daytime, latitude $\sim 20^\circ$ N (Hawaii, August 2012). (middle panel) daytime, latitude $\sim 55^\circ$ N (North Sea, July 2015). (bottom panel) night, mid-latitude.

8. Summary and Conclusions

The FRI-2018 model uses rocket observations and modeling to provide an impressive range of ionospheric electron number densities for heights, in 1-km intervals, from 150 km down to nominally 55 km though only those above 60 km, and $>10^6 \text{ m}^{-3}$, are considered reliable (Friedrich et al., 2018). Such profiles are provided at 5 latitudes (0° – 60°), 11 solar zenith angles, 3 solar activities and for each of 12 months, making a total of 1,980 profiles. While a significant amount of VLF data exists, this has not yet been sorted comprehensively. We therefore chose to select three representative conditions where extensive quality VLF data were available: nighttime at mid-latitudes, daytime at low latitude ($\sim 20^\circ$) and daytime toward the higher end of the FRI-2018 latitudes ($\sim 55^\circ$). While VLF data can give very good electron number densities, the technique works only in the lowest parts of the Earth's ionosphere. During quiet times this means ~ 75 – 90 km altitude by night, ~ 55 – 70 km near mid-day at low latitudes, and ~ 60 – 75 km at higher solar zenith angles or mid-latitudes. Both the VLF technique and this lowest ionospheric region are none-the-less very important in geophysics. Both man-made and natural (e.g., lightning) VLF radio signals propagate in the Earth-ionosphere waveguide and so are sensitive to the lower D region, by day, night, dawn and dusk, both in quiet conditions and when disturbed, such as by energetic electron (or proton) precipitation or by solar flares or by (extra-terrestrial) gamma rays etc. Production and loss modeling in the D region is very important in attempting to get a good quantitative understanding of the key mechanisms there. Recently Siskind et al. (2018) found that, by day, their theoretical photochemical model agreed better with VLF-derived electron densities below 68–70 km than the corresponding electron densities from rockets. There appear to be multiple complex processes involved in such D region modeling which are likely to result in further developments but need the support of measured electron number densities.

The agreement found here between VLF-derived and FRI-2018 electron densities is, perhaps surprisingly, best at night. This may be partly because, at night, both techniques have good sensitivity in the height range 75–90 km. The MF rocket sensitivity has a threshold of typically between 10^9 m^{-3} at 60 km and 10^7 m^{-3} at 80 km (Friedrich et al., 2018; Jacobsen & Friedrich, 1979) and so in the relevant height range of 75–90 km the more sensitive but much less certain Langmuir probe data will have been only minimally required, if at all. Also, the lower limit for electron densities of 10^6 m^{-3} (at any height) will have been encountered only below ~ 75 km. In addition, both techniques will have averaged out any of the “ledges” mentioned here in Section 5: VLF by averaging over long paths and by averaging over several days, and FRI-2018 by averaging over a number of rocket profiles.

For daytime low latitudes ($\sim 20^\circ$), the agreement found here between VLF-derived and FRI-2018 electron densities is quite reasonable, specifically for the electron densities at the solar zenith angles of $\sim 10^\circ$ and $\sim 75^\circ$ and more generally in terms of the agreement between observed and

ModeFinder + FRI-2018 predicted VLF phases and amplitudes as functions of solar zenith angle. It might well be fair to say that probably from the very general, global perspective of the FRI-2018 model, the agreement is quite good, while from the highly focused VLF perspective, the agreement is more modest.

For daytime high mid-latitudes ($\sim 55^\circ$), the degree of agreement found here between VLF-derived and FRI-2018 electron densities is at best only quite modest, with the FRI-2018 densities being consistently lower by up to a factor of ~ 3 than the VLF-derived values. Again, it might well be fair to say that probably from the very general, global perspective of the FRI-2018 model, the agreement is generally passable while, from the highly focused

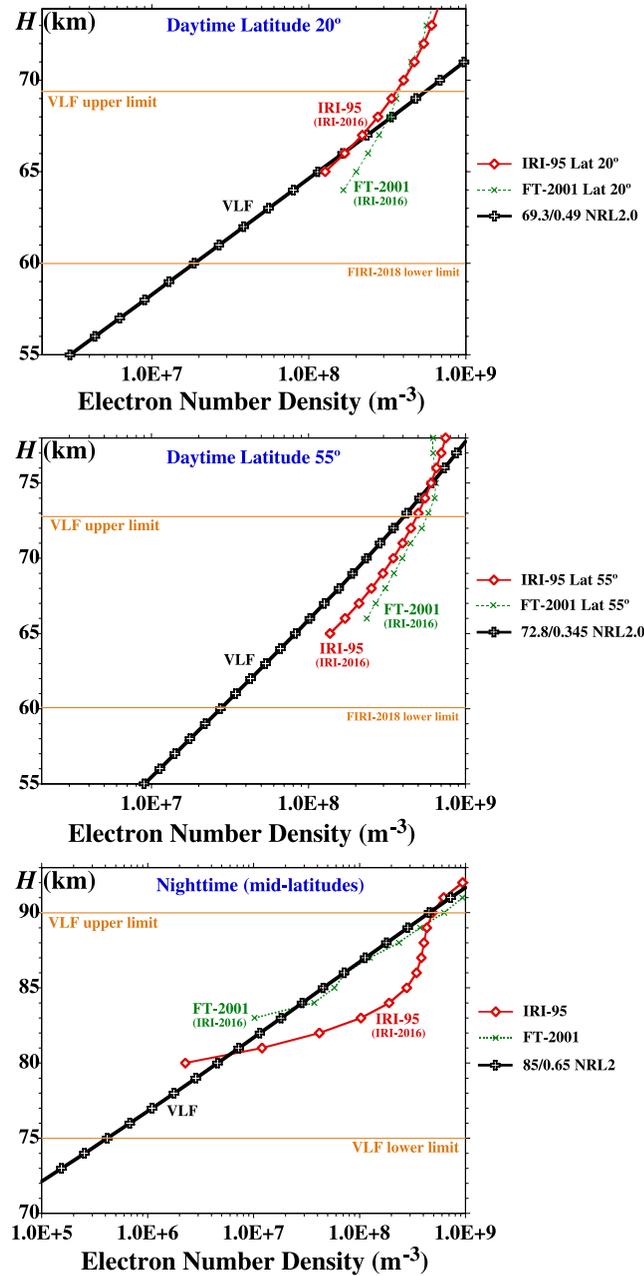


Figure 8. VLF-derived electron number densities compared with those from the International Reference Ionosphere, IRI-2016. Comparisons are with each of the two D region components of IRI-2016: IRI-95 and FT-2001 discussed in the text. (Top panel) Daytime at low latitudes ($\sim 20^\circ N$). (Middle panel) Daytime at high mid-latitudes ($\sim 55^\circ N$). (Bottom panel) Nighttime at mid-latitudes.

VLF perspective, the agreement is fairly marginal. A significant factor, at these higher latitudes in daytime, may well be that FRI-2018, at the lowest altitudes, in compromising between the rocket (probe, rather than MF) observations and modeling, has perhaps underestimated the galactic cosmic ray generated electrons. In contrast, the FRI-2018 profiles at low latitude (20°) and high solar zenith angle (75°), in the middle panel of Figure 1, show an electron density “bulge” at ~ 65 km, which appears consistent with cosmic ray generated electrons, below a “dip” at 70 km down to which height Lyman- α is likely no longer penetrating at this high solar zenith angle.

Although, as mentioned above, the FIRI-2018 model does not extend above latitudes $\sim 60^\circ$, comparisons of VLF-derived electron densities in polar regions with those from sounding rocket profiles and an MF radar have shown them to be quite comparable both by day and by night (Thomson et al., 2018, 2021).

Overall, our comparisons here with specific VLF-derived electron number densities indicate that the FIRI-2018 model, though under-estimating at higher latitudes by day, is a major advance at least compared with IRI-2016. We conclude that FIRI-2018 reasonably represents the *D* region electron number density profiles on a near global basis (i.e., at latitudes below 60°), at least in the low height ranges where it can be compared with VLF measurements.

Data Availability Statement

The NRLMSIS 2.0 neutral atmospheric data were obtained using the FORTRAN code at <https://map.nrl.navy.mil/map/pub/nrl/NRLMSIS/NRLMSIS2.0/> but should also be available from <https://kauai.ccmc.gsfc.nasa.gov/instantrun/msis>. The NRLMSISE-00 neutral atmospheric data were obtained from the NRLMSISE-00 web model at <https://ccmc.gsfc.nasa.gov/modelweb/models/nrlmsise00.php> but should also be available from <https://kauai.ccmc.gsfc.nasa.gov/instantrun/msis>. The FIRI-2018 model profiles are available from <https://figshare.com/s/357cb03b3e5bed649bbc> (Friedrich et al., 2018) or <https://figshare.com/search?q=FIRI-2018>. The VLF data used here came from Thomson et al. (2007, 2014, 2017). US Navy code LWPC is available at <https://github.com/mlhutchins/LWPC>. The US Navy computer program referred to here as ModeFinder is a slightly modified version of MODEFNDR (e.g., Nunn & Strangeways, 2000; Thomson, 1993) and MODESRCH described and listed in Morfitt and Shellman (1976). Solar zenith angles were determined at <https://gml.noaa.gov/grad/solcalc/>. The IRI data used here came from the NASA web page: https://ccmc.gsfc.nasa.gov/modelweb/models/iri2016_vitmo.php.

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