

1 **Stable AMOC off state in an eddy-permitting**

2 **Coupled Climate Model**

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7 **Abstract** Shifts between on and off states of the Atlantic Meridional Over-
8 turning Circulation (AMOC) have been associated with past abrupt climate
9 change, supported by the bistability of the AMOC found in many older, coarser
10 resolution, ocean and climate models. However, as coupled climate models
11 evolved in complexity a stable AMOC off state no longer seemed supported.
12 Here we show that a current-generation, eddy-permitting climate model has
13 an AMOC off state that remains stable for the 450-year duration of the model
14 integration. Ocean eddies modify the overall freshwater balance, allowing for
15 stronger northward salt transport by the AMOC compared with previous, non
16 eddy-permitting models. As a result, the salinification of the subtropical North
17 Atlantic, due to a southward shift of the intertropical rain belt, is counteracted
18 by the reduced salt transport of the collapsed AMOC. The reduced salinifi-
19 cation of the subtropical North Atlantic allows for an anomalous northward
20 freshwater transport into the subpolar North Atlantic dominated by the gyre
21 component. Combining the anomalous northward freshwater transport with

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22 the freshening due to reduced evaporation in this region helps stabilise the
23 AMOC off state.

24 **Keywords** AMOC · AMOC collapse · abrupt climate change · hosing
25 experiment · CGCM · eddy-permitting

26 1 Introduction

27 The Atlantic Meridional Overturning Circulation (AMOC) describes the merid-
28 ional volume transport in the Atlantic Ocean (Wunsch, 2002). The AMOC
29 brings warm waters to the high latitude North Atlantic, warming the climate
30 of Northern and Western Europe. A collapse of the AMOC would lead to
31 drastic changes in surface air temperatures over much of the Northern Hemi-
32 sphere, in particular in the Northeast Atlantic where temperatures can drop
33 by 9°C (Manabe and Stouffer, 1988; Vellinga and Wood, 2002; Jackson et al,
34 2015). As a consequence of anthropogenic climate change, warming of the high
35 latitude North Atlantic and the addition of freshwater through enhanced pre-
36 cipitation, increased melting of sea-ice and icebergs, as well as more runoff
37 from the Greenland ice sheet can cause the sinking branch of the AMOC to
38 weaken and potentially shut down. Hereafter, we refer to a collapsed AMOC
39 as an AMOC off state while, the AMOC circulation, as it is known today, is
40 referred to as an AMOC on state.

41 Climate model projections indicate a likely weakening of the AMOC, but
42 a complete collapse was deemed unlikely in the latest IPCC report by Collins
43 et al (2013). However, models have difficulty correctly simulating past abrupt

44 climate changes, including an AMOC collapse, affecting the likelihood of sim-
45 ulating future abrupt climate change (Valdes, 2011; Drijfhout et al, 2011).
46 Paleo-proxy data have shown evidence for wide spread abrupt climate change
47 events in the times before the Holocene from ice-core records (Dansgaard et al,
48 1993; Blunier and Brook, 2001) and sediment cores (de Abreu et al, 2003). A
49 possible interpretation of these events is that they are associated with switches
50 between AMOC on and off states in the past (Broecker et al, 1990), although
51 the spatial extent of these abrupt changes in climate can still be questioned
52 (Wunsch, 2006). Such switches can be theoretically understood from simple
53 box model studies showing that under the same forcing conditions it is possible
54 to have both a stable AMOC on and off state, or only a mono-stable regime
55 depending on the forcing (Stommel, 1961; Marotzke, 1990; Rahmstorf, 1996).
56 The existence of bistability in these box models depends on the freshwater
57 forcing. Similarly, some coupled climate models have found a bistable AMOC
58 dependent on freshwater forcing when freshwater hosing was applied contin-
59 uously (Hawkins et al, 2011; Hu et al, 2012; Sijp, 2012). However, in newer
60 coupled climate models after applying freshwater hosing for a set amount of
61 time the AMOC recovered after the freshwater hosing was stopped (Peltier
62 et al, 2006; Krebs and Timmermann, 2007; Jackson, 2013) while it was pos-
63 sible to maintain the AMOC off state in some older coupled climate models
64 (e.g. UVic and GFDL R30 models in Stouffer et al (2006)).

65 To identify the transition between the two regimes of mono- and bistability
66 it was proposed that the sign of the freshwater transport by the AMOC in the

67 Atlantic can be used as an indicator for its stability (referred to here as M_{ov} but
68 often also referred to as F_{ov}) (Rahmstorf, 1996; de Vries and Weber, 2005).
69 When used as an indicator for AMOC stability, M_{ov} is typically measured
70 at the southern entrance of the Atlantic near 34°S . A positive M_{ov} at 34°S
71 indicates that the AMOC imports freshwater into the Atlantic and a negative
72 M_{ov} at 34°S indicates freshwater export from the Atlantic. In an AMOC off
73 state M_{ov} is expected to tend towards zero, thereby creating an anomalous
74 salt import into the Atlantic for positive M_{ov} which leads to a destabilisation
75 of the AMOC off state. On the other hand, when M_{ov} is negative an AMOC
76 collapse will result in an anomalous freshwater import into the Atlantic helping
77 stabilise the AMOC off state. Therefore, a positive M_{ov} can be associated with
78 a mono-stable AMOC while a negative M_{ov} can be associated with a bistable
79 AMOC (Huisman et al, 2010). Observational estimates of M_{ov} at the southern
80 boundary of the Atlantic based on ship data or estimated from ARGO float
81 data support a negative M_{ov} , suggesting that the present day AMOC resides
82 in the bistable regime (Bryden et al, 2011; Garzoli et al, 2013). It has been
83 recommended that the divergence of the freshwater transport into the Atlantic
84 by the AMOC, $\Delta M_{ov} = M_{ovS} - M_{ovN}$, where S (N) is the southern (northern)
85 boundary of the Atlantic is a better indicator of bistability (Huisman et al,
86 2010; Liu and Liu, 2013).

87 When the AMOC weakens and even collapses, the reduction in northward
88 heat transport causes a wide spread cooling of the northern hemisphere surface
89 air temperatures (Manabe and Stouffer, 1988; Vellinga and Wood, 2002; Jack-

90 son et al, 2015). The cooling leads to a southward/equatorward shift of the
91 latitude of maximum heating causing the dividing latitude of the northern and
92 southern hemisphere Hadley circulations to shift southward (Drijfhout, 2010),
93 displacing the Intertropical Convergence Zone (ITCZ). The southward shift of
94 the ITCZ causes a reduction of precipitation in the subtropical North Atlantic
95 region leading to a salinification of the ocean. The saltier waters in this region
96 can be transported into the high latitude regions of the North Atlantic through
97 large-scale instabilities kick starting the convection (e.g. the large-scale eddy
98 generated in GFDL CM2.1 in Yin and Stouffer (2007)). Therefore, in order for
99 the AMOC off state to remain stable this salinification needs to be balanced
100 by an equally large freshening term, due to changes in ocean circulation.

101 In the GFDL R30 model the freshening associated with ocean circulation
102 changes is large enough to counteract the salinification due to the southward
103 ITCZ shift because the overturning circulation reverses (Yin and Stouffer,
104 2007). In that case Antarctic Intermediate Water (AAIW) sinks to a depth of
105 1000 m just south of South America and is transported northward, then up-
106 wells in the North Atlantic subtropical gyre. This circulation has been named
107 the reverse thermohaline circulation (RTHC). However, the RTHC only devel-
108 ops in coarse-resolution ocean models and often is deeper than just the upper
109 1000 m (Dijkstra, 2007; Hawkins et al, 2011; Sijp, 2012). In a newer generation
110 of coupled climate models the RTHC cell does not develop (e.g. GFDL CM2.1
111 in Yin and Stouffer (2007)) and without the additional freshwater transport of
112 the RTHC the subtropical gyre becomes so salty that a fresh subpolar ocean

113 without deep sinking is no longer stable and the AMOC recovers (Yin and
114 Stouffer, 2007; Jackson, 2013). The reason for the RTHC not to develop is
115 that stronger atmospheric feedbacks promote saltier and colder thermocline
116 water in the subtropical North Atlantic, reducing the north-south pressure
117 gradient between the subtropical North Atlantic and subpolar South Atlantic
118 that is driving the RTHC (Yin and Stouffer, 2007)).

119 In the very latest coupled climate models ocean eddies and swifter bound-
120 ary currents are allowed for, changing the salt balance in the Atlantic. Ocean
121 eddies freshen the subtropical gyre by exchanging water with the tropics and
122 subpolar gyre (Tréguier et al, 2012). As a result, eddy-permitting and eddy-
123 resolving models must feature a larger mean flow salt transport divergence
124 into the subtropical gyre to maintain equilibrium counteracting freshening by
125 the eddies. The larger mean flow salt transport divergence could allow for a
126 stronger advective salt feedback associated with an AMOC collapse without
127 the need of developing an RTHC. Indeed, using a higher resolution coupled
128 climate model Spence et al (2013) achieved a stronger drop and slower recov-
129 ery of the AMOC in a high-resolution model relative to a coarser resolution
130 model in a relatively weak and short freshwater hosing experiment. Similarly,
131 Weijer et al (2012), using an ocean only model, were able to show that the
132 drop in AMOC in response to a freshwater hosing was stronger in the higher
133 resolution model. Both studies suggest that the AMOC off state in higher
134 resolution models could become stable. Here we discuss whether a larger salt
135 transport by the AMOC into the North Atlantic subtropical gyre, which is

136 typical for higher resolution ocean models, can sustain a stable off state, even
137 if the RTHC does not develop, using a 450 year long hosing experiment in an
138 eddy-permitting coupled climate model.

139 **2 Model Configuration and Experiment Setup**

140 **2.1 Model Configuration**

141 For this study the Global Climate version 2 (GC2) (Williams et al, 2015) con-
142 figuration of Hadley Centre Global Environmental Model version 3 (HadGEM3)
143 (Hewitt et al, 2011) is used. This coupled climate model consists of an ocean,
144 atmosphere, sea-ice and land-surface model coupled together with data ex-
145 changing between the atmosphere and ocean components every 3 hours. The
146 ocean model component of GC2, HadGEM3 uses the Global Ocean version
147 5 (GO5) (Megann et al, 2013) of the ORCA025 configuration of the Nucleus
148 for European Modelling of the Ocean (NEMO) (Madec, 2008) version 3.4.
149 The ORCA025 grid uses a tri-polar structure with poles over Antarctica,
150 Siberia and Canada and has a horizontal resolution of 0.25° , with the res-
151 olution decreasing when moving towards the poles so that the grid remains
152 quasi-isotropic. The ocean model contains 75 vertical levels with thicknesses
153 ranging from 1 m at the surface and increasing with depth up to 200 m in
154 the bottom layer. The sea-ice model is the global sea ice version 6 (GSI6)
155 configuration of the Los Alamos National Laboratory sea ice model (CICE)
156 version 3.4 (Rae et al, 2015) and is used at the same model grid as the ocean

157 model. The Global Atmosphere version 6 (GA6.0) of the Met Office unified
158 model is used with a horizontal resolution of N216, which has a resolution of
159 about 60 km in mid-latitudes, and has 85 levels in the vertical leading to an
160 improved resolution in the stratosphere. Global Land version 6 (GL6) config-
161 uration of the land model Joint UK Land Environment Simulator (JULES) is
162 also used in this model setup but none of its data is analysed in this study.
163 Heat, freshwater and momentum fluxes are passed between the atmosphere
164 and ocean/ice model every three hours through the OASIS coupler while the
165 ocean and sea-ice model exchange fluxes every ocean model time step (22.5
166 min) without the use of flux adjustment. The eddy permitting resolution of
167 the ocean model has lead to a reduction in the North Atlantic cold bias and
168 the atmospheric model shows improved Atlantic and European blocking events
169 (Scaife et al, 2011) and the ability to better predict the winter North Atlantic
170 Oscillation (Scaife et al, 2014), in previous versions of the HadGEM3 model
171 setup, i.e. GloSea5.

172 2.2 Experiment Setup

173 In this study two experiments from the GC2 model are considered, a 150-year
174 long present day control simulation and a 450 year long hosing experiment.
175 The hosing experiment is a continuation of the experiment analysed in Jackson
176 et al (2015) (See reference for more details). The present day control simula-
177 tion was started from a 36 year long development run of HadGEM3, which
178 was initialised with EN3 data (Ingleby and Huddleston, 2007) averaged over

179 2004-2008 and the hosing experiment is started from year 42 of the control
180 experiment. The control simulation uses CO₂ concentrations based on 1978
181 levels and held constant throughout both simulations. The main goal of the
182 hosing experiment was to collapse the AMOC, therefore, the methodology is
183 based on Vellinga and Wood (2002), which allows for a rapid collapse of the
184 AMOC but is very idealised. For the first 10 years of the hosing experiment
185 the salinity in the model is perturbed by an amount equivalent to a hosing
186 of 10 Sv, making a total of 100 Sv-years additional freshwater. This is done
187 through reducing the salinity in the Atlantic Ocean north of 20°N and in the
188 Arctic by 0.64 psu in the upper 350 m and then tapering to zero over the
189 next 186 m (Fig. 1). This is done instantaneously every December 1 and, as is
190 common practice in hosing experiments, is compensated by adding 0.008 psu
191 everywhere else in the ocean allowing for the total salinity to be conserved
192 (Fig. 1). After the 10 years of hosing is completed and the model is allowed to
193 continue without changes for another 440 years.

194 **3 Results**

195 In the 450 year long hosing experiment the AMOC is able to collapse and
196 remain very weak for the entire duration of the model integration (Fig. 2).
197 During and after the 10 year hosing period the ocean begins to adjust, with
198 salinity anomalies slowly spreading southward from the hosing region towards
199 the equator and also spreading downward in the water column. Since we want
200 to discuss the evolution of the ocean fields in 100 year time-slices, we will take

201 the period 311-410 (301 to 400 years after the hosing stopped) as representa-
202 tive for the final state of the model. The mean salinity is 0.86 psu fresher in the
203 hosing region towards the end of the hosing simulation (years 311-410) relative
204 to the control simulation. The sea surface salinity (SSS) anomaly with respect
205 to the control run features a comma shaped pattern in the North Atlantic sub-
206 tropical gyre (Fig. 3a), as typical with most fresh water hosing experiments
207 (Krebs and Timmermann, 2007; Yin and Stouffer, 2007). The sea surface tem-
208 peratures (SSTs) also drop due to the reduction of northward heat transport
209 from the AMOC off state (Fig. 3b and Jackson et al (2015)). The decrease in
210 SSTs allow for the seasonal sea-ice to extend further southward reaching as
211 far south in winter as the Grand Banks, as well as covering a large portion of
212 the Norwegian and Baltic Seas (Fig. 3b). The reductions in SSS and SST fall
213 within the range of what has been seen in previous modelling studies with a
214 similar magnitude of freshwater hosing (Yin and Stouffer, 2007).

215 3.1 AMOC Streamfunction

216 The control simulation features a realistic AMOC with a maximum strength
217 of 17.4 Sv at 27°N and at a depth of 773 m in the mean (Fig. 4a). The depth
218 reached by the North Atlantic Deep Water cell is slightly shallower than that in
219 observations (3000 m as opposed to 4000 m in (Kanzow et al, 2010; Smeed et al,
220 2014)), a common problem in ocean models (Danabasoglu et al, 2014). The
221 Faroe Bank Channel overflow (defined as waters denser than $\sigma_\theta = 27.8 \text{ kg/m}^3$)
222 is slightly weaker in this model than in observations (1.8 Sv as opposed to 1.9

223 Sv (Hansen and Østerhus, 2007)). This overflow is mainly missing the weak
224 cold waters below 0 °C which account for the majority of the overflow waters
225 in the observations, making the model overflow less dense. For the Denmark
226 Strait the overflows are considerably weaker when considering waters denser
227 than $\sigma_\theta = 27.8 \text{ kg/m}^3$ (1.4 Sv as opposed to 3.4 Sv (Jochumsen et al, 2012)),
228 which again is missing the very cold water masses. However, for the Denmark
229 Strait choosing the density cut off to be $\sigma_\theta = 27.8 \text{ kg/m}^3$ misses a lot of the
230 overflow waters. By choosing the density class cut off of to be $\sigma_\theta = 27.6 \text{ kg/m}^3$,
231 matching the depth of density cutoff in Jochumsen et al (2012), the overflow
232 increases to 2.9 Sv. These differences in the overflows between the model and
233 observations could potentially lead to the shallower North Atlantic Deep Water
234 cell. The main convection sites are in the Labrador Sea, Greenland Sea and
235 South of Iceland (Fig. 4b) as expected from observations (de Boyer Montégut
236 et al, 2004). However, the too buoyant overflows could potentially account
237 for the slightly weaker and shallower AMOC as compared to observations at
238 26.5°N (Fig. 2a, 15.7 Sv as opposed to 17.5 Sv (Smeed et al, 2014)) but this
239 is not investigated in more detail.

240 Based on an AMOC index at 26.5°N and between 500-2000 m the AMOC
241 collapses very rapidly during the hosing, leading to a minimum in AMOC at
242 year 4 (Fig. 2a). After the hosing has stopped the AMOC recovers slightly,
243 achieving a maximum at year 21, before dropping in strength again and re-
244 maining in a very weak state for the duration of the model integration. How-
245 ever, there is a noticeable weak trend in the AMOC index at 26.5°N which by

246 the end of the model integration causes the AMOC to increase in strength to
247 just over 5 Sv (Fig. 2a). This increase in AMOC strength is slow and occurs
248 later in the model integration than seen in previous climate model studies
249 (Vellinga and Wood, 2002; Stouffer et al, 2006; Jackson, 2013). Also, it only
250 applies to a shallow, wind-driven, AMOC that does not extend further north
251 than the subtropics. Considering an AMOC index further to the north (maxi-
252 mum between 50°N - 65°N and 500-2000 m depth) the AMOC collapse shows
253 no hint of recovering (Fig. 2b). There is no sign of increasing mixed layer depth
254 in the subpolar North Atlantic due to the onset of deep convection (Fig. 4d).
255 Both subtropical and subpolar wind-driven cells are enhanced near the sur-
256 face related to the positive North Atlantic Oscillation (NAO) that develops
257 in response to the AMOC collapse (Jackson et al, 2015). The AMOC stream-
258 function does not develop a stable RTHC after the AMOC collapses. Despite
259 this, the AMOC off state appears stable, at least for 450 years. In year 311-410
260 there appears to be no convection present in the high latitude regions (Fig. 4d)
261 and similarly the overflows in the Denmark Strait and Faroe Banks Channel
262 have completely collapsed to 0 with no signs of recovery.

263 3.2 Atmospheric Response

264 The southward shift of the ITCZ is reflected in the net precipitation (pre-
265 cipitation - evaporation + runoff, PER) and causes a reduction in the surface
266 freshwater flux into the ocean just north of the equator and an increase south of
267 the equator (Fig. 5a,b). These changes in PER reduce the amount of freshwater

268 added to the subtropical North Atlantic with the majority of the reduction in
269 precipitation occurring in the subtropical North Atlantic which loses 0.047 Sv
270 in years 311-410 (Table 1, Fig. 5c). This reduction in PER is an atmospheric
271 feedback to the AMOC collapse that acts to destabilise the AMOC off state
272 by salinifying the North Atlantic.

273 Over the subpolar North Atlantic evaporation is reduced due to the increase
274 in sea-ice cover blocking latent heat exchange and the decrease in atmospheric
275 temperatures reducing the amount of atmospheric water vapour content (Ta-
276 ble 1, Fig. 5b) (Drijfhout, 2014). Despite the reduction in evaporation being
277 small relative to the precipitation changes in the subtropical regions, it is large
278 enough to outweigh the reduction in precipitation over the subpolar Atlantic.
279 The subsequent increase in PER causes an anomalous freshening of the sinking
280 regions (Fig. 5d) with a magnitude of 0.042 Sv in the years 311-410 (Table 1).
281 The rate at which the precipitation and evaporation anomalies change reduces
282 as the model integration continues, especially for the evaporation. This subpo-
283 lar freshening is an atmospheric feedback that stabilises the AMOC off state
284 through freshening the North Atlantic. The salinification over the subtropical
285 North Atlantic is marginally stronger than the freshening over the subpolar
286 North Atlantic (Table 1). The salinification of the subtropical North Atlantic
287 could eventually lead to more saline waters being transported in the subpolar
288 North Atlantic as seen in the GFDL CM2.1 model (Yin and Stouffer, 2007).
289 Nevertheless, the off state remains stable here, while it quickly destabilises
290 in the GFDL CM2.1 model. It should be noted that the initial atmospheric

291 response of precipitation and evaporation in HadGEM3 is similar to that in
 292 GFDL CM2.1. However, in the GFDL CM2.1 model the precipitation anoma-
 293 lies associated with a southward shift of the ITCZ are not maintained as the
 294 AMOC recovers, while here the anomaly continues to show the characteristic
 295 dipole pattern over the equator although the amplitude is slowly decreasing
 296 (Fig. 5c). This brings up the question which additional feedbacks are present
 297 in HadGEM3, stabilising the AMOC off state? To answer this question we
 298 analyse in detail the freshwater budget in the subtropical and subpolar North
 299 Atlantic.

300 3.3 Freshwater Budget

301 The freshwater budget analysis is based on an extension to the calculations
 302 detailed in Drijfhout et al (2011) (see appendix for details). The freshwater
 303 budget can be summarised as follows:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + PER + M_{mix}, \quad (1)$$

304 where M_{trend} is the freshwater trend in the region of interest, $\Delta M_{ov/az/eddy}$
 305 represents the divergence of the freshwater transport for the specific region,
 306 in our case the southern boundary minus the northern boundary, for the vari-
 307 ous components of the transport, PER is the precipitation minus evaporation
 308 plus runoffs over the specific region of interest and finally M_{mix} is the residual
 309 term of the budget, mainly comprised of mixing along the boundaries. In eqn.
 310 1 the decomposition of mean flow transport divergence into an overturning

(M_{ov}) and gyre (M_{az}) component was motivated by the much stronger coupling between M_{ov} and AMOC than between M_{az} and AMOC at the southern boundary of the Atlantic, when they budget is applied to the Atlantic as a whole. Especially in the North Atlantic subpolar gyre this decomposition can be questioned. However, this framework can still be used to link area integrated changes in freshwater budget to changes in the AMOC, especially in the North Atlantic subtropics. It appears that changes in M_{ov} are first order in eqn. 1 and can be understood from the AMOC collapse as they are dominated from changes in the zonal mean velocity field. In addition it allows for comparison with observations where freshwater transports have been diagnosed using the same framework (McDonagh et al, 2010; Bryden et al, 2011; Garzoli et al, 2013; McDonagh et al, 2015-in press). When the model is in an equilibrium state the changes in PER are approximately balanced by changes in freshwater transport by overturning circulation (M_{ov}), azonal circulation (M_{az}) and eddies (M_{eddy}). We apply the freshwater budget analysis to the subtropical North Atlantic, defined as 10°N to 45°N , and to the subpolar North Atlantic, defined to be 45°N to 70°N . These boundaries were chosen to coincide with the boundaries of the subtropical and subpolar gyres, with the subpolar gyre region containing the main sinking regions of the North Atlantic. The region specific freshwater budget analyses are summarised in Table 1 and graphically in Fig. 6. The atmospheric contributions to the freshwater budget have already been discussed; below we discuss the freshwater transport terms as well as the freshwater budget as a whole.

334 3.3.1 Freshwater Transports

335 The AMOC off state is associated with changes in the freshwater transport
336 terms that must be able to balance the changes in PER, especially in the
337 subtropical North Atlantic, to prevent a salinification of the North Atlantic
338 and hence a return to the AMOC on state. In the control simulation the
339 freshwater transport due to the overturning, M_{ov} , is negative throughout the
340 entire Atlantic Ocean, indicating that the AMOC is transporting freshwater
341 southward/salt northward (Fig. 7). The negative M_{ov} at 34°S is consistent with
342 observations (Bryden et al (2011), Garzoli et al (2013), McDonagh Personal
343 Communications based on McDonagh and King (2005)), despite being slightly
344 weaker, and is a possible indication for a bistable AMOC (Fig. 7). After the
345 AMOC collapses the magnitude of M_{ov} , as expected, decreases and over time
346 adjusts to a new equilibrium (Fig. 8a). The reduction in magnitude of M_{ov} can
347 be attributed to the reduction in AMOC transport, with changes in salinity
348 only having a small effect (Fig. 8b). These changes lead to an anomalous
349 northward transport of freshwater south of 45°N in the Atlantic Ocean (Fig.
350 8b). Even more important, however, is the sign of the divergence of M_{ov} instead
351 of the sign of M_{ov} itself, since it is the divergence that determines whether
352 or not there will be a freshening or salinification in the region of interest.
353 The subtropical North Atlantic has an increase of 0.132 Sv of freshwater due
354 to the changes in the divergence of M_{ov} (Fig. 6c, Table 1). The associated
355 increase in freshwater is twice the amount of freshwater required to balance
356 the anomalous salinification caused by changes in PER (Table 1). Changes in

357 the divergence of M_{az} and M_{eddy} need to enhance the salinification caused by
358 PER and thereby balance the changes in divergence of M_{ov} .

359 The salinity decrease after the AMOC collapse is largest at the eastern side
360 of the basin, which does not only hold for the surface (Fig. 3a) but, also at
361 depth (not shown). This decrease in salinity is strongest over the southward
362 branch of the subtropical gyre and northward branch of the subpolar gyre,
363 near the eastern boundary, leading to changes in M_{az} . This results in a de-
364 crease in M_{az} in the subtropical gyre and an increase in M_{az} in the subpolar
365 gyre, while changes at the gyre boundaries are small (e.g 10°N and 45°N)
366 (Fig. 8c,d). Relative to the climate models in Yin and Stouffer (2007), which
367 have a coarser resolution, HadGEM3 has larger amplitude in M_{az} divergence,
368 also leading to larger changes in its divergence after the collapse. This is due
369 to the increase in model resolution leading to stronger gyres (Tréguier et al,
370 2005; Spence et al, 2013) and less east-west difference in salinity bias (Yin and
371 Stouffer (2007) their Fig. 1), likely due to the Gulf Stream separation being
372 too far north in lower resolution models. The change in divergence of M_{az} for
373 both the subtropical and subpolar North Atlantic reduces the amount of fresh-
374 water being transported into these regions (Fig. 6, Table 1). This anomaly in
375 freshwater transport partially balances the additional freshwater being added
376 to the subtropical North Atlantic by changes in M_{ov} (Table 1, Fig. 6c) and
377 changes in the subpolar gyre PER and mixing (Fig. 6b).

378 The resolution of the model used in this study allows for the analysis of the
379 effect of eddies on the freshwater budget from the equator to mid latitudes.

380 Here, freshwater transport due to eddies is defined as the difference between
381 total freshwater transport and freshwater transport calculated by using the
382 seasonal fields only (see appendix for more details). The main effect of the
383 eddies is to exchange water between the subtropical and subpolar gyres, fresh-
384 ening the former and salinifying the latter (Fig. 8e and Table 1). Immediately
385 after the AMOC collapse the salinity gradient at the edge of the hosing re-
386 gion becomes very large leading to a large increase in the southward freshwater
387 transport by the eddies at 20°N. Within a few decades after the freshwater hos-
388 ing M_{eddy} becomes relatively small again compared to M_{ov} and M_{az} with val-
389 ues similar to the control integration (Fig. 8e). In the eddy-permitting model
390 the freshening of the subtropical North Atlantic by M_{eddy} and the increased
391 freshening by a larger M_{az} play a similar role to the flux adjustment in coarser
392 resolution climate models in the control integration (e.g. GFDL R30 model in
393 Yin and Stouffer (2007)). This helps to stabilise the freshwater budget by al-
394 lowing for a larger negative M_{ov} in the control integration and subsequently
395 a larger change in M_{ov} after the AMOC collapses. The change in M_{ov} is now
396 large enough to balance all other terms in the freshwater budget without leav-
397 ing a strong positive salinity trend in the subtropical North Atlantic. As model
398 resolution increases further towards eddy-resolving the magnitude of M_{eddy} is
399 expected to become even larger (Tréguier et al, 2012), further adding to the
400 stabilising effect of the eddies.

401 3.3.2 Total Freshwater Budget

402 The total freshwater trend in the subtropical North Atlantic still shows a small
403 salinification over the 311-410 year period, slightly stronger than the salini-
404 fication in the control run (Table 1, Fig. 6). Despite the salinification of the
405 subtropical North Atlantic, the subpolar North Atlantic shows a freshening
406 trend, enhancing the salinity gradient between the two (Table 1, Fig. 6). The
407 anomalous freshening trend of the subpolar North Atlantic can be attributed
408 to the combination of decreased evaporation in this region, the anomalous
409 northward freshwater transport at the gyre boundaries and an increased mix-
410 ing term (i.e. M_{mix}). The gradient in salinity across the North Atlantic, despite
411 being stronger than in the control integration, does not lead to large-scale in-
412 stabilities that suddenly give rise to very strong salinity transports as seen
413 in Yin and Stouffer (2007). The eddies are likely helping to keep the gradient
414 small enough to avoid a sudden large-scale instability to develop and to restart
415 the convection in the high latitude sinking regions.

416 4 Discussion

417 The AMOC response to freshwater perturbations has been previously investi-
418 gated in a large CMIP/PMIP coordinated experiment (Stouffer et al, 2006).
419 A freshwater hosing of 0.1 Sv and 1 Sv was applied for 100 years, versus a
420 hosing of 10 Sv over 10 years in the present experiment. Of the nine models
421 involved in the 100 Sv-year hosing experiment, seven models had started the

422 transition from the off state back to the on state before 100 years after the
423 completion of the hosing. Two models remained in the off state; one model of
424 intermediate complexity, Uvic, and one older GFDL model, GFDL R30. The
425 different behaviour between GFDL R30 and a newer version, GFDL CM2.1,
426 was afterwards analysed (Yin and Stouffer, 2007) and it was argued that the
427 stable off state in GFDL R30 was maintained by flux adjustment and weak
428 atmospheric feedbacks allowing the RTHC to develop. This result led to the
429 paradigm that newer generation climate models that no longer use flux adjust-
430 ment and feature more realistic atmospheric dynamics are not able to maintain
431 a stable AMOC off state (Yin and Stouffer, 2007; Liu et al, 2014). Here we
432 show that an eddy-permitting coupled climate model is able to maintain a
433 stable AMOC off state for 440 years after the hosing is completed, which is
434 more than twice as long as the runs performed in the CMIP/PMIP experiment.
435 The increase in freshwater transport into the subtropical North Atlantic due to
436 higher-resolution eddies and increased boundary currents allow the AMOC to
437 transport more salt northwards across the entire Atlantic basin. This stronger
438 advective salt feedback is key for the model to be able to counteract the strong
439 atmospheric response over the tropical/subtropical North Atlantic basin that
440 features in complex climate models when the AMOC collapses. In a sense, ed-
441 dies and swifter boundary currents play a similar role in the freshwater budget
442 to the flux adjustment used in older generation climate models.

443 Some coupled climate models of lower complexity have been integrated for
444 even longer durations with some of them having the AMOC off state become

445 unstable after many centuries (Krebs and Timmermann, 2007). We cannot ex-
446 clude that such a transition will eventually occur in HadGEM3, but at present
447 there is no deep water formation site returning to the high latitude North
448 Atlantic (Fig. 4d) and the freshwater budget shows no signs of a potential
449 recovery. While the subtropical North Atlantic is continuing to increase its
450 salinity, albeit with a very small trend, the subpolar North Atlantic is getting
451 relatively fresher, hampering the restart of deep convection. Also when taking
452 the subpolar North Atlantic and the Arctic into account there is an overall
453 freshening trend suggesting that having a return of deep convection in the high
454 latitude North Atlantic in the near future is very unlikely.

455 When taking the salinity of the entire Atlantic into account, as was done
456 in Sijp (2012), we do not see a difference in salinity between the hosing and
457 control simulations. In Sijp (2012) the two states in Atlantic mean salinity are
458 associated with the AMOC on and off states. However in Sijp (2012) an RTHC
459 develops, which is responsible for the low salinity state, while in HadGEM3
460 the AMOC off state still has a shallow wind-driven cell that extends into the
461 Northern Hemisphere, preventing a low salinity state. However if we focus
462 on the region north of 35°N only, the hosing integration is 0.7 psu fresher
463 in the upper 3000 m than the control integration, indicating that low and
464 high salinity states in the subpolar gyre can be associated with the AMOC on
465 and off state in this model. This suggests that a bifurcation in basin average
466 salinity no longer exists in HadGEM3 but bistability in subpolar gyre salinity
467 is still existent.

468 The increase in northward salt transport by the AMOC in HadGEM3,
469 relative to the coarser resolution climate models (Yin and Stouffer, 2007) is
470 associated with a reduction in vertical gradient of salinity bias in the Atlantic.
471 The model using flux adjustment in Yin and Stouffer (2007), GFDL R30,
472 showed little bias, but the climate model that did not use flux adjustment,
473 GFDL CM2.1, featured larger biases. In particular, the salinity bias in the
474 GFDL CM2.1 model contained a pronounced vertical gradient with a negative
475 salinity bias near the surface and a positive bias at deeper levels throughout
476 most of the Atlantic. Combined with an AMOC that transports surface water
477 northward and deep water southward this salinity bias leads to M_{ov} being
478 strongly biased towards positive values. With a positive M_{ov} , when the AMOC
479 collapses, more saline water will be transported into the Atlantic, aiding the
480 recovery of the AMOC, as is clearly the case with GFDL CM2.1 in Yin and
481 Stouffer (2007). These results are supported by the analysis of Liu et al (2014),
482 where they see a larger negative salinity bias in the surface for the un-flux
483 adjusted models relative to flux adjusted models. This led to a less negative
484 M_{ov} at 34°S, reducing the likelihood of bistability. For the model used in this
485 study, HadGEM3, the salinity bias has a weak negative vertical gradient in the
486 Southern Atlantic in the depths corresponding with the North Atlantic Deep
487 Water (NADW) cell of the AMOC and a mostly positive bias in the upper
488 1000 m throughout the rest of the Atlantic (Fig. 9). This weaker salinity bias
489 is likely due to the fact that the model is eddy permitting and has swifter
490 and narrower boundary currents. In GFDL CM2.1 the positive salinity bias

491 peaks near 20°N (Yin and Stouffer, 2007), while in HadGEM3 the model bias
492 is smaller there (Fig. 9) since 20°N coincides with a convergence in freshwater
493 transport due to the eddies (Fig. 8e). The vertical structure of the salinity
494 bias in HadGEM3 is too small to affect the sign of M_{ov} : it only has a minor
495 effect on M_{ov} south of the equator and an even weaker effect between the
496 equator and 30°N (Fig. 7). However, a further reduction of the salinity bias
497 would move the model values of M_{ov} even closer to the estimates based on
498 observations of M_{ov} throughout the Atlantic (Fig. 7).

499 At 26°N M_{ov} is -0.601 Sv in the control integration of HadGEM3 (about
500 -0.6 Sv GFDL CM2.1 Yin and Stouffer (2007)) and -0.78 Sv in observations
501 (McDonagh et al, 2015-in press). A larger difference between HadGEM3 and
502 the models analysed in Yin and Stouffer (2007) occurs at the southern bound-
503 ary of the subtropical gyre (10°N). In HadGEM3 M_{ov} is largely negative at
504 those latitudes, -0.361 Sv, while in GFDL CM2.1 M_{ov} has about half the
505 amplitude, approximately -0.2 Sv. Both models agree on M_{ov} being slightly
506 negative at the subtropical-subpolar boundary, around -0.2 Sv. Thus the dif-
507 ferent values at the southern boundary of the subtropical gyre in the models
508 determines the sign of the divergence of M_{ov} over the subtropical gyre and
509 the sign of the advective salt feedback in this area when the AMOC weak-
510 ens or collapses. Unfortunately there are no estimates of M_{ov} near 10°N , but
511 the reduced salinity bias in HadGEM3 suggests that a negative M_{ov} at those
512 latitudes is the more likely.

513 Of some concern is the absence of an RTHC in the AMOC streamfunction
514 after hosing is applied. Stability analysis of coarse-resolution ocean-only mod-
515 els suggests that the collapsed AMOC is an unstable steady state, dividing
516 the attractor space between a stable on state and a stable RTHC reaching
517 to the bottom of the Atlantic (Dijkstra, 2007). Furthermore, the studies of
518 Saenko et al (2003) and Sijp et al (2012) point out that it is the density dif-
519 ference between the NADW and the Antarctic Intermediate Water (AAIW)
520 formation regions which are important for the existence of an RTHC. In this
521 study the density of the NADW formation region is not reduced enough after
522 the initial hosing to become lighter than the water in the AAIW formation
523 region as RTHC is not maintained. This study and the results of Yin and
524 Stouffer (2007) suggest that the development of the RTHC is suppressed by
525 atmospheric feedbacks. However, there is at present insufficient analysis to
526 conclude whether atmospheric feedbacks really prevent a stable RTHC to de-
527 velop, or whether there are other reasons for why it is absent in HadGEM3.
528 For HadGEM3, we believe there are two possibilities; 1) the AMOC off state,
529 despite the maintaining an AMOC off state for much longer than the models
530 used in the PMIP experiment of Stouffer et al (2006), will eventually return to
531 an AMOC on state, or 2) the AMOC off state is a stable solution of coupled
532 climate models at eddy-permitting or higher resolution.

533 In HadGEM3 the presence of eddies and swifter boundary currents (stronger
534 gyres) allows for stronger northward salt transport of the AMOC, stabilising
535 the off state (Fig. 8). An even higher-resolution (1/12 degree), eddy-resolving

536 ocean model features even larger northward salt transport by M_{eddy} than the
537 eddy-permitting version (Tréguier et al, 2012), implying an AMOC off state
538 could potentially be favoured by even stronger advective salt feedbacks. On
539 the other hand, the latitudinal structure of M_{ov} in HadGEM3 seems broadly
540 consistent with the few estimates we have at different latitudes (Fig. 7) and
541 we anticipate only a small improvement in this respect when going to higher
542 resolution in the ocean component of climate models.

543 **5 Conclusions**

544 The goal of the model run analysed in this study was to rapidly collapse the
545 AMOC and study the stability of the AMOC off state. Several other studies
546 have been done choosing a freshwater hosing setup that more realistically rep-
547 represents what could happen in the climate system (Weijer et al, 2012; Spence
548 et al, 2013; Swingedouw et al, 2013). These studies have all shown that it is
549 possible to weaken the AMOC using a more realistic hosing setup. On top of
550 that Weijer et al (2012) and Spence et al (2013) have shown that when using
551 higher resolution the amount by which the AMOC weakens is larger relative
552 to their coarse resolution models used in those studies. However, these studies
553 often only have been run for 50 years in the high resolution setting. These re-
554 sults plus the results presented in this study support the possibility of coupled
555 models being more likely to model abrupt climate changes as model resolutions
556 continue to improve. At higher resolution a stronger advective salt feedback
557 associated with the AMOC, leading to a freshening of the subtropical North

558 Atlantic, overcomes the damping feedback that salinifies this region, associated
 559 with the atmospheric response to an AMOC collapse. This changed balance
 560 between the different feedbacks makes the transition to a stable AMOC off
 561 state possible, when the freshwater transports at high latitudes in the North
 562 Atlantic increases. This is illustrated by the eddy-permitting climate model,
 563 HadGEM3, being able to maintain an AMOC off state for 440 years.

564 Appendix: Freshwater Budget Calculation

565 The freshwater budget calculation used in this study is based on the method presented
 566 in Drijfhout et al (2011) with modifications to include the effects of a northern and southern
 567 boundary, as well as specifics to the version of NEMO used (GO5, version 3.4 of NEMO)
 568 (Megann et al, 2013). Mean flow transports are based on 3 month means, while total trans-
 569 ports (i.e. vS) are calculated online and are updated after each ocean model time step, which
 570 are later averaged over the years of interest removing the effects of the seasonal cycle on the
 571 budget. Following Drijfhout et al (2011), the equation for the volume budget is as follows:

$$V_t = T_S - T_N - T_{Med} + PER - Res_V, \quad (2)$$

572 where V_t is the rate of change of the volume, $T_{(N/S)}$ are volume transports through the
 573 northern and southern boundaries, T_{Med} is the volume transport through the Strait of
 574 Gibraltar, PER is the precipitation minus evaporation plus runoffs and Res_V is the error
 575 generated by the choice of differencing scheme and temporal resolution of the data. The
 576 value of Res_V is computed as a residual to close the budget. Since the model has a free
 577 surface V_t is equivalent to the changes in the sea surface height using backwards differencing.
 578 The main differences between eqn. 2 and eqn. 4 in Drijfhout et al (2011) are that we have left
 579 the choice of the northern and southern boundaries as arbitrary as opposed to choosing $34^\circ S$
 580 and the Bering Strait and we have included a term, T_{med} for the volume transports through
 581 the Strait of Gibraltar. In this configuration of NEMO the transports are computed without
 582 taking the changes in sea surface height into account. For the regions of interest used in this
 583 study the values of Res_V are relatively small resulting in $O(10^{-4} \text{ Sv})$ for the North Atlantic

584 subtropical gyre and $O(10^{-5}$ Sv) for the North Atlantic subpolar gyre, which in both cases
 585 is the smallest term in the budget with the remaining terms ranging from $O(10^{-3}$ Sv) to
 586 $O(1$ Sv). Choosing instantaneous values of sea surface height from the model restart files in
 587 the computation of V_t leads to Res_V having the same order as the precision in which the
 588 data is stored but, not all model restart files were available.

589 Similarly the salinity budget in terms of freshwater becomes the following:

$$M_{trend} - V_t = M_S - M_N - M_{Med} + M_{Mix} - Res_V + H, \quad (3)$$

590 where M_{trend} is the rate of change of freshwater in the region of interest, $M_{(N/S)}$ are the
 591 northward/southward freshwater transports, M_{med} is the freshwater transport through the
 592 Strait of Gibraltar, H represents the freshwater hosing and M_{mix} , computed as a residual,
 593 closes the budget capturing mixing and errors introduced by the temporal resolution of the
 594 data, as well as, the choice of reference salinity, S_o . The conversion between salinity based
 595 terms to the freshwater based terms in eqn. 3 is done through multiplying all the terms in
 596 the equation by $-1/S_o$. Note that we have dropped the negative sign before M_{trend} in eqn.
 597 3, contrary to Drijfhout et al (2011) so that positive values indicate an increase in freshwater
 598 not salinity. In this case the hosing is included in the salinity budget and not the volume
 599 budget since it is computed as a redistribution of salinity in this model study. Combining
 600 eqns. 2 and 3 gives the following expression for the fresh water budget:

$$M_{trend} = (M_S + T_S) - (M_N + T_N) - (M_{Med} + T_{Med}) + M_{Mix} + PER + H. \quad (4)$$

601 The $M_{(N/S)}$ terms can be divided into eddy and mean flow components since the ocean
 602 model output includes vS computed at every model time step. The eddy contribution to
 603 the freshwater transport is defined as follows:

$$M_{(eddy(N/S))} = \frac{-1}{S_o} \int_{N/S} (\overline{vS} - \overline{v}\overline{S})dA = M_{(N/S)} - M_{(mean(N/S))}, \quad (5)$$

$$\rightarrow M_{(N/S)} = M_{(mean(N/S))} + M_{(eddy(N/S))}, \quad (6)$$

604 where the integral is taken over each zonal section of the Atlantic basin, \overline{vS} is the total
 605 seasonal mean transport, \bar{v} and \bar{S} are the seasonal mean meridional velocity and salinity and
 606 $M_{(mean(S/N))} = -1/S_o \int_{N/S} \bar{v}\bar{S}dA$ represents the non-eddy transports, with the overbar
 607 denoting a mean computed over 3 months. A map of the eddy kinetic energy (Fig. 10)
 608 shows that the eddy field in HadGEM3 is very similar to other models of similar resolution
 609 (Delworth et al, 2012), perhaps even slightly closer to what is expected from observations.
 610 The eddy contribution is computed in a very similar way to Tréguier et al (2012), in which it
 611 was also shown that the eddy contribution will be even stronger at higher model resolutions.
 612 Since the current model resolution is eddy-permitting it is not possible to completely resolve
 613 eddies at all latitudes, therefore caution must be taken in interpreting the role of the eddies
 614 in the high latitudes. Similar to what is done in Drijfhout et al (2011), $M_{(mean(S/N))}$ can
 615 be divided into an overturning $M_{(ov(S/N))}$, azonal $M_{(az(S/N))}$ and the volume transport
 616 $T_{(S/N)}$ terms as follows:

$$M_{mean(N/S)} = M_{ov(N/S)} + M_{az(N/S)} - T_{(N/S)}, \quad (7)$$

$$M_{ov(N/S)} = \frac{-1}{S_o} \int_{N/S} v^* \langle S \rangle dA, \quad (8)$$

$$M_{az(N/S)} = \frac{-1}{S_o} \int_{N/S} v' S' dA, \quad (9)$$

617 where $\langle f \rangle = \int f dx / \int x$ is the zonal mean, $f' = f - \langle f \rangle$ is the difference from the zonal
 618 mean, $\hat{f} = \int f dA / \int dA$ is the zonal section mean or barotropic component and $f^* = \langle f \rangle - \hat{f}$
 619 is the zonal mean baroclinic component for $f = \bar{v}$ or $f = \bar{S}$. Substituting eqns. 6 and 7 into
 620 eqn. 4 gives the final form for the zonal freshwater budget equation:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + \Delta M_{Med} + M_{Mix} + PER + H, \quad (10)$$

621 where $\Delta M_{ov} = M_{(ov(S))} - M_{(ov(N))}$, $\Delta M_{az} = M_{(az(S))} - M_{(az(N))}$, $\Delta M_{eddy} = M_{(eddy(S))} -$
 622 $M_{(eddy(N))}$ and $\Delta M_{Med} = -M_{Med} - T_{Med}$.

623 There are several possible valid choices of the reference salinity; the mean salinity over the
 624 entire volume of the region used in the budget calculation, the mean salinity over the section

625 used as the northern (southern) boundary or the mean salinity from the Strait of Gibraltar.
626 For this study it was chosen to use the mean salinity at the boundary between the North
627 Atlantic subtropical and subpolar gyres for S_o , the reference salinity. Choosing one of the
628 other salinities as a reference salinity creates a maximum difference of $O(10^{-4} \text{ Sv})$, which is
629 less than 10% of the smallest value represented in our budget analysis. To further simplify
630 the budget analysis only times when there is no hosing being applied are considered and
631 the freshwater transport through the Strait of Gibraltar is combined with the mixing term,
632 resulting in the following final equation for the budget analysis:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + M_{mix} + PER. \quad (11)$$

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Table 1 Summary of freshwater budget for subtropical (10°N-45°N) and subpolar North Atlantic (45°N-70°N). All values are given in Sv with positive values indicating an addition of freshwater into the region. The bottom row of each section is the anomalous change in freshwater (i.e. Hosing (311-410) - Control).

		Overturning ΔM_{ov}	Azonal ΔM_{az}	Eddy ΔM_{eddy}	P-E+R PER	Mixing M_{mix}	Total M_{trend}
Subpolar N. Atlantic	Control	-0.170	-0.090	-0.033	0.241	0.041	-0.012
	Hosing (311-410)	-0.165	-0.187	-0.031	0.283	0.093	-0.007
	Anomaly	0.006	-0.097	0.001	0.042	0.052	0.005
Subtropical N. Atlantic	Control	-0.164	0.524	0.102	-0.504	-0.009	-0.051
	Hosing (311-410)	-0.032	0.490	0.061	-0.551	-0.024	-0.055
	Anomaly	0.132	-0.033	-0.041	-0.047	-0.015	-0.004

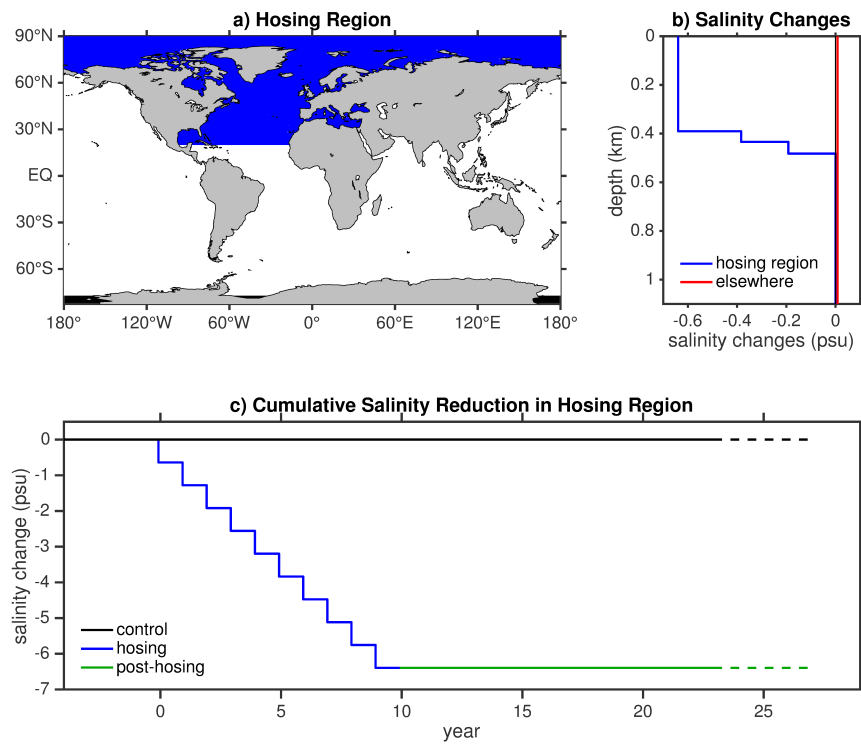


Fig. 1 (a) The region where the freshwater hosing is applied. (b) The redistribution of salinity in the hosing region (blue) and everywhere else (red). (c) The cumulative salinity reduction in the hosing region (upper 350 m) in the model experiments for the control (black), hosing (blue) post-hosing (green).

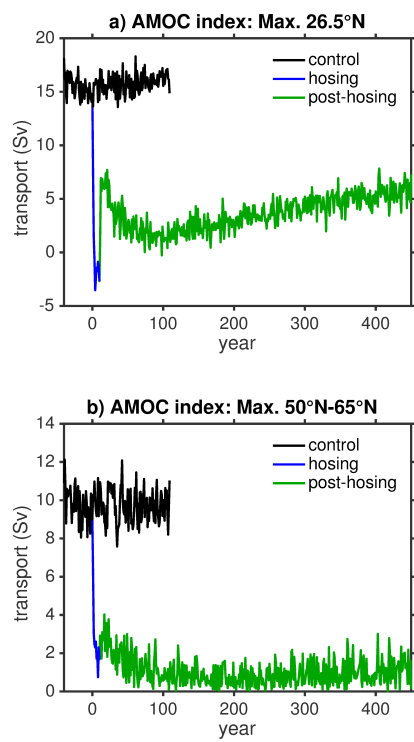


Fig. 2 (a) The AMOC index computed as the maximum AMOC streamfunction at 26.5°N below a depth of 500 m and above 2000 m for the control experiment (black), hosing period (blue) and post-hosing period (green). (b) same as a expect computed between 50°N and 65°N .

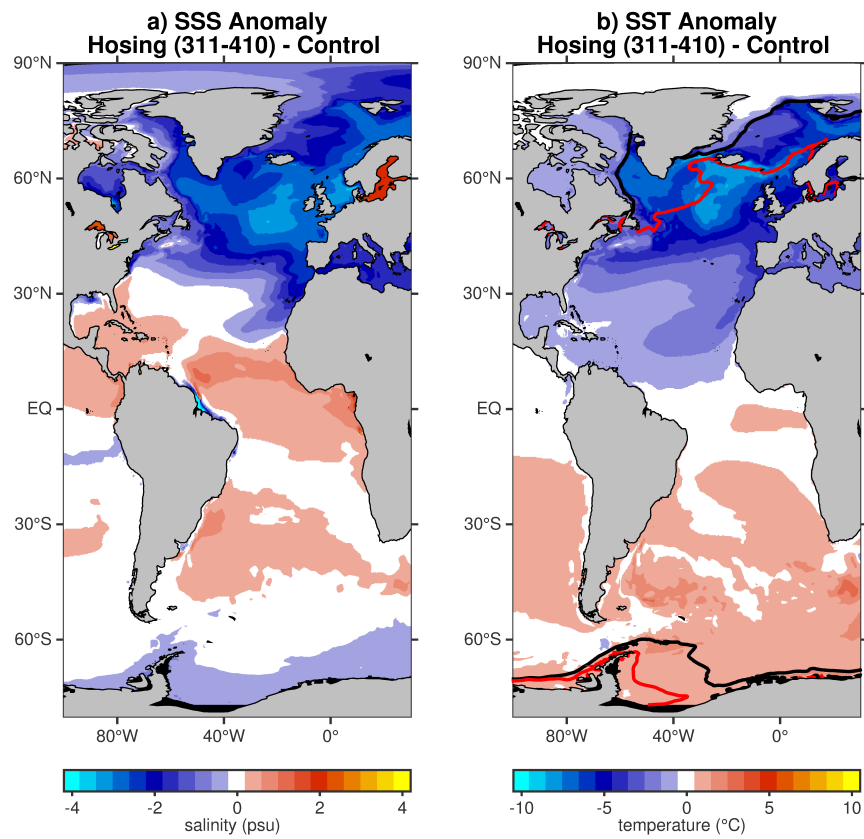


Fig. 3 (a) Mean SSS from years 311-410 of the hosing simulation minus the mean SSS from the control simulation. (b) same as in (a) but for SST with the black contour indicating the annual maximum sea-ice extent in the control simulation and the red contour from years 311-410 of the hosing simulation.

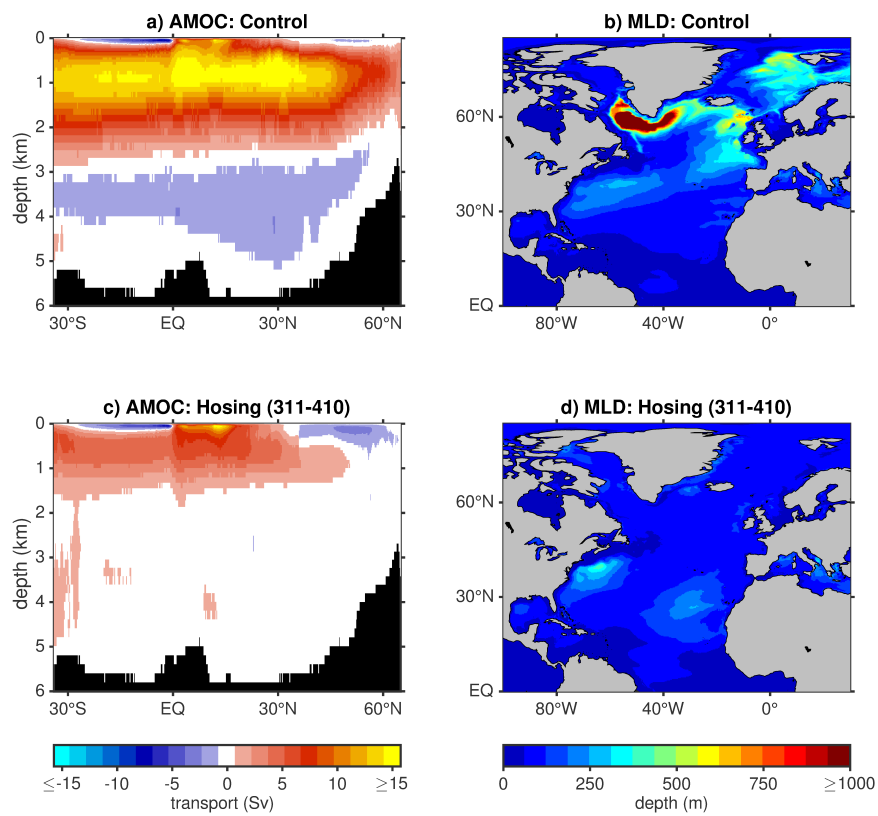


Fig. 4 (a) The mean AMOC streamfunction and (b) the mean annual maximum mixed layer depth from the control simulation. (c) and (d) same as a and b but for years 311-410 of the hosing simulation.

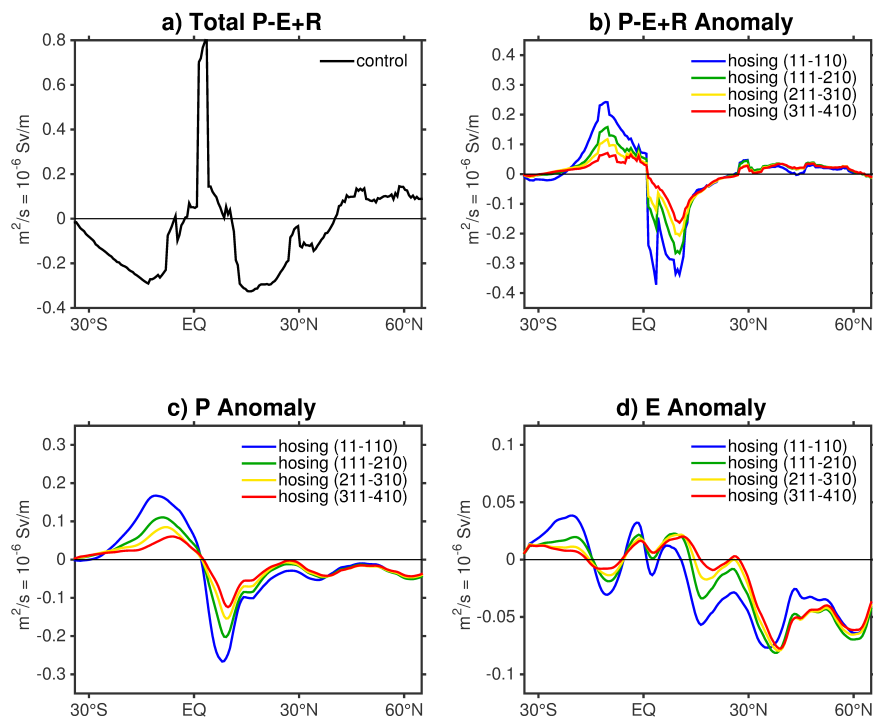


Fig. 5 (a) The zonally integrated P-E+R from the control simulation normalized to Sv per meter in latitude. (b) the anomalous P-E+R from various 100 year means in the hosing simulation, (c) same as b but for precipitation only, (d) same as b but for evaporation only with blue years 11-110, green years 111-210 yellow years 211-310 and red years 311-410. All data is smoothed using a 2° latitude window to reduce the spikes from the river runoffs.

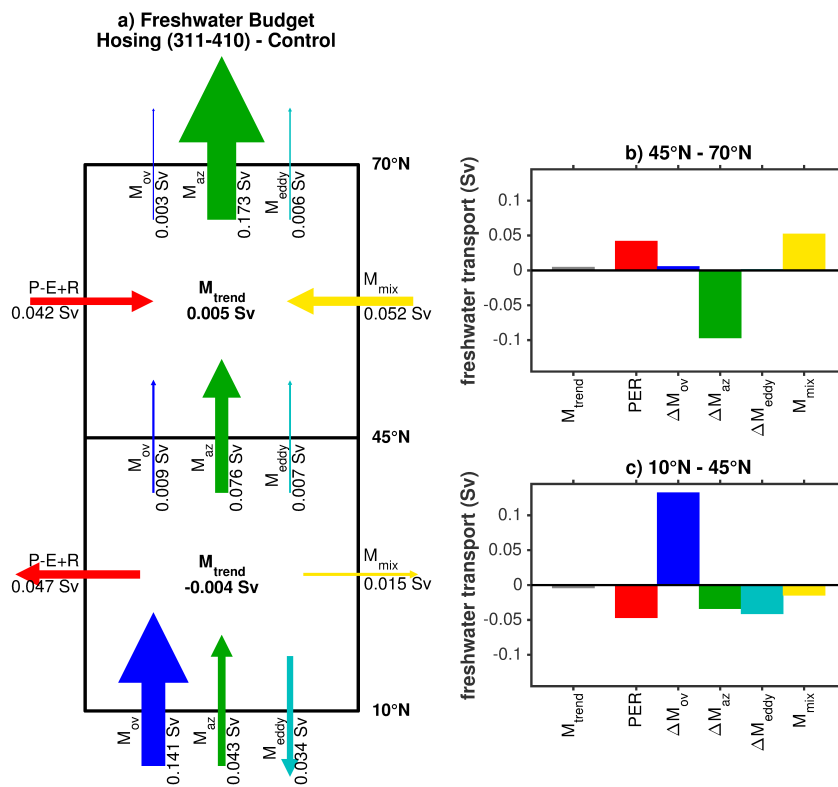


Fig. 6 (a) Anomalous freshwater budget boxes for the subtropical (10°N-45°N) and subpolar (45°N-70°N) North Atlantic. The width of the arrows and arrow heads have been scaled according to the strength of the freshwater transport anomalies. (b) Summary of the anomalous freshwater budget for the subpolar North Atlantic. (c) Same as (b) but for the subtropical North Atlantic.

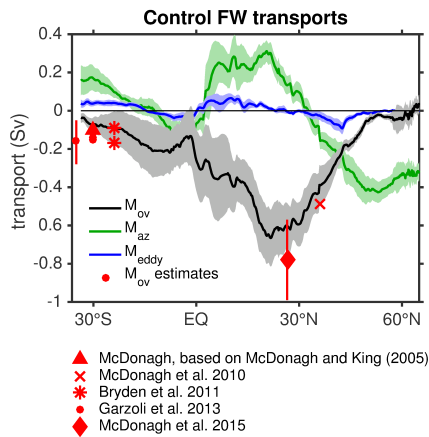


Fig. 7 Mean M_{ov} from control simulation with \pm one standard deviation of seasonal data (black/grey shading), mean M_{az} from control simulation with \pm one standard deviation of seasonal data (green/green shading) and mean M_{eddy} from control simulation with \pm one standard deviation of seasonal data (blue/blue shading). Estimates of M_{ov} from observations (red): triangle based on McDonagh and King (2005); cross McDonagh et al (2010); stars Bryden et al (2011); circles Garzoli et al (2013) with vertical line representing the range in estimates; and diamond McDonagh et al (2015-in press) with the vertical line indicating the standard deviation of 10 day timeseries. Note that the standard deviations/range are computed using data available on different timescales.

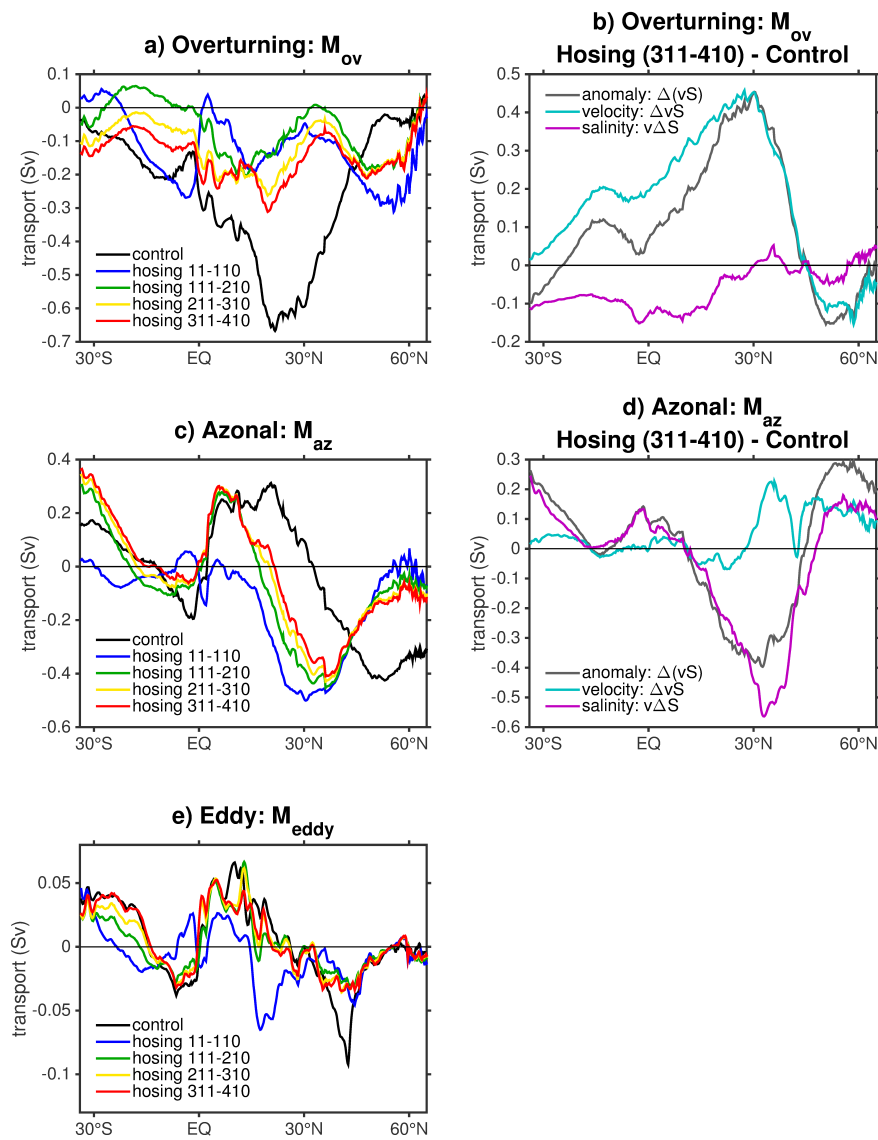


Fig. 8 The freshwater transports along latitude bands in the Atlantic. (a) Freshwater transport due overturning M_{ov} . The different colours represent different means over various years; control (black), hosing 11-110 (blue), hosing 111-210 (green), hosing 211-310 (yellow) and hosing 311-410 (red). (b) Decomposition of M_{ov} anomalies (hosing years 311-410 minus control) into contributions from velocity (cyan) and salinity (magenta) compared to total anomaly (dark gray). (c and d) same as (a and b) but for M_{az} and (e) same as (a) but for M_{eddy} . Note the different scales on panels a-e.

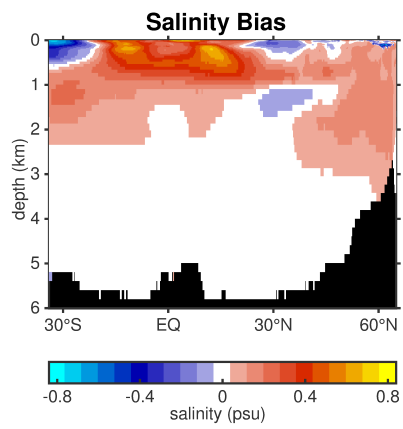


Fig. 9 Zonal mean salinity bias of the control experiment relative to EN3 data (Ingleby and Huddleston, 2007).

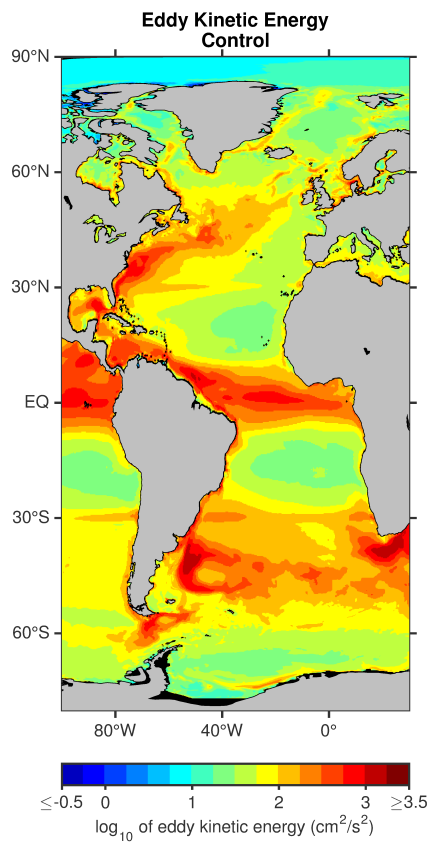


Fig. 10 Logarithm of the surface eddy kinetic energy in the control simulation. The eddy kinetic energy was computed from the model's surface velocity fields using the difference between the instantaneous velocities and seasonal mean velocities before averaging over all years of the simulation.