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1	Multi-scale sensible heat fluxes in the suburban environment from
2	large aperture scintillometry and eddy covariance
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## 10 Abstract

Sensible heat fluxes  $(Q_H)$  are determined using scintillometry and eddy covariance over a 11 12 suburban area. Two large aperture scintillometers provide spatially integrated fluxes across path 13 lengths of 2.8 km and 5.5 km over Swindon, UK. The shorter scintillometer path spans newly built residential areas and has an approximate source area of 2-4 km<sup>2</sup>, whilst the long path extends from 14 the rural outskirts to the town centre and has a source area of around 5-10 km<sup>2</sup>. These large-scale 15 16 heat fluxes are compared with local-scale eddy covariance measurements. Clear seasonal trends are revealed by the long duration of this dataset and variability in monthly  $Q_H$  is related to the 17 18 meteorological conditions. At shorter time scales the response of  $Q_H$  to solar radiation often gives 19 rise to close agreement between the measurements, but during times of rapidly changing cloud 20 cover spatial differences in the net radiation ( $Q^*$ ) coincide with greater differences between heat fluxes. For clear days  $Q_H$  lags  $Q^*$ , thus the ratio of  $Q_H$  to  $Q^*$  increases throughout the day. In 21 22 summer the observed energy partitioning is related to the vegetation fraction through use of a 23 footprint model. The results demonstrate the value of scintillometry for integrating surface 24 heterogeneity and offer improved understanding of the influence of anthropogenic materials on 25 surface-atmosphere interactions.

## 26 Keywords

27 Energy balance; Large aperture scintillometer; Seasonality; Sensible heat flux; Urban

## 28 1. Introduction

Understanding the interactions between the land surface and the atmosphere is central to developing our predictive power in terms of weather forecasting, air quality events, thermal comfort, flood risk and tools for urban design. The surface energy balance has been closely linked to land use and land cover, from studies both within and between cities. This has been achieved largely 33 through eddy covariance (EC) measurements at multiple sites in a city (e.g. in Los Angeles (Grimmond et al. 1996), Basel (Christen and Vogt 2004), Łódź (Offerle et al. 2006), Melbourne 34 (Coutts et al. 2007), Montreal (Bergeron and Strachan 2010), Dublin (Keogh et al. 2012), Essen 35 36 (Weber and Kordowski 2010), Oberhausen (Goldbach and Kuttler 2013) and Helsinki (Nordbo et al. 2013)); and through comparison of measurements from different cities (e.g. Grimmond and Oke 37 38 (1995; 2002)). Also, studies at individual sites have combined footprint models and land cover maps 39 to capture differences in the surface cover sampled as the source area of the EC measurement 40 changes with atmospheric conditions (Vesala et al. 2008; Järvi et al. 2012; Goldbach and Kuttler 41 2013). The derived relations between surface cover and fluxes offer valuable indications of the 42 underlying processes and form a basis for modelling turbulent fluxes (Grimmond and Oke 2002; Järvi et al. 2011; Loridan and Grimmond 2012). 43

44 Scintillometry provides a means to estimate fluxes at a much larger scale than eddy covariance, typically of the order of several km<sup>2</sup> (Hoedjes et al. 2007; Guyot et al. 2009; Kleissl et al. 2009a). 45 46 Path-averaging along the electromagnetic beam of the scintillometer means that measurements are 47 inherently spatially integrated, offering a particular advantage over heterogeneous surfaces (Beyrich et al. 2002; Meijninger et al. 2002b; Evans 2009; Samain et al. 2011a). Despite the complexity of the 48 49 urban surface, patches of impervious land cover (roads, car parks, paved areas) adjacent to green 50 spaces (parks, gardens) are not very different to the juxtaposition of fields containing differently 51 ripening and senescing crops in mixed agricultural landscapes. In such studies, measuring sufficiently 52 high above the surface ensures the influences of surface heterogeneity are well-blended at the height of the measurement and reliable fluxes can be obtained (Meijninger et al. 2002b; Ezzahar et 53 54 al. 2007). In addition to the increased spatial representativeness of such large-area measurements, 55 their increased scale facilitates comparison with satellite remote sensing products or land-surface models. A key conclusion of the work by Chehbouni et al. (2000a) was that model development 56 57 should focus on establishing relations to replicate observations made at the large-scale. Studies comparing model output with scintillometry data include the work of Cheinet et al. (2011), Samain 58 59 et al. (2011b), Steeneveld et al. (2011) and Maronga et al. (2013).

The use of scintillometers in urban environments can be divided into two groups: (a) studies involving small aperture instruments, usually deployed within or near the top of the roughness sublayer on path lengths of the order of 100 m, e.g. in Tokyo (Kanda et al. 2002), Basel (Roth et al. 2006) and London (Pauscher 2010); and (b) studies over much longer path lengths (500 m - 10 km) using large aperture scintillometers. This study uses large aperture scintillometry. The much larger sampling volume enables robust retrieval of turbulence statistics and the measurement sensitivity is 66 greatest near the centre of the path, away from the instruments and their mounting structures, such 67 that the influence of locally-produced turbulence around these structures has minimal impact on the 68 measurements. Furthermore, direct access to the measurement area is not required – an 69 electromagnetic beam is simply transmitted high above the surface – unlike point measurements 70 requiring *in situ* mounting. This remote sensing capability makes the scintillometry technique 71 particularly valuable in the urban environment.

There are an increasing number of studies using large aperture scintillometry in urban areas. Lagouarde et al. (2006) present results from a three week trial in summer 2001 in which two large aperture scintillometers were used over Marseille. Other studies include work in London (Gouvea and Grimmond 2010), Nantes (Mestayer et al. 2011), Łódź (Zieliński et al. 2012) and Helsinki (Wood and Järvi 2012).

77 To date, the use of large aperture scintillometry in urban areas has been to derive the sensible 78 heat flux. In order to directly estimate the latent heat flux, a second scintillometer of longer 79 wavelength is required (Hill et al. 1988; Andreas 1989), however these are not yet commercially 80 available. Hence there are only a handful of two-wavelength studies documented (Kohsiek and 81 Herben 1983; Green et al. 2001; Meijninger et al. 2002a; Meijninger et al. 2006; Evans 2009). A 82 millimetre-wave scintillometer was installed alongside one of the infrared scintillometers used in this 83 study. These first two-wavelength results from the urban environment are presented in two 84 companion papers (Ward et al. in preparation a,b).

The goal of the work presented here is to investigate the influence of the surface on sensible 85 86 heat fluxes across different spatial and temporal scales in the suburban environment. Eddy 87 covariance measurements are analysed together with results from large aperture scintillometers 88 installed on 2.8 km and 5.5 km paths in Swindon, UK. The two year study period enables analysis of 89 seasonal and inter-annual patterns (Sect. 3.1) as well as the short-term response of heat fluxes to 90 solar radiation (Sect. 3.2). Consideration is given to the experimental uncertainties, including the 91 representativeness of point measurements required as inputs for scintillometry algorithms (Sect. 92 2.3) and the observational challenges associated with urban environments. The contributions of 93 different land cover classes are related to the observations at each scale by applying a footprint 94 model (Sect. 3.3).

## 95 2. Methodology

### 96 2.1. Derivation of the sensible heat flux from single-wavelength scintillometry

97 Scintillometers measure the intensity of an electromagnetic beam after propagation through the 98 turbulent atmosphere. Changes in beam intensity are related to the strength of turbulence and can 99 be converted to the structure parameter of the refractive index of air  $(C_n^2)$ . First, these refractive 100 index fluctuations must be related to temperature fluctuations,

101 
$$C_T^2 \approx \frac{T^2}{A_T^2} C_n^2$$
, (1)

102 where  $C_T^{\ 2}$  is the temperature structure parameter and  $A_T$  is the structure parameter coefficient 103 which depends on temperature (*T*), pressure (*p*) and weakly on humidity (Hill et al. 1980; Andreas 104 1988; Ward et al. 2013b). The Bowen ratio ( $\beta$ ) correction given by Wesely (1976) (see Moene (2003) 105 for full details and a discussion) is not implemented in Equation 1 because  $\beta$  is not known *a priori*. 106 The effect of this is discussed in Sect. 2.3.

107 The conversion from structure parameters to fluxes entails iteration of similarity functions,  $f_{MO}(\zeta)$ , 108 using Monin-Obukhov similarity theory (MOST) and the effective height of the scintillometer,  $z_{ef}$ , the 109 wind speed, U (which is measured at height  $z_U$ ), displacement height,  $z_d$ , and aerodynamic 110 roughness length,  $z_0$ . Commonly used forms of the similarity functions are

111 
$$f_{MO}(\zeta) = c_{T1}(1 - c_{T2}\zeta)^{-2/3}$$
 (2)

112 for unstable conditions and

113 
$$f_{MO}(\zeta) = c_{T1}(1 - c_{T3}\zeta^{2/3})$$
 (3)

114 for stable conditions. The stability variable  $\zeta$  is given by  $(z_m - z_d)/L_{Ob}$ , where  $L_{Ob}$  is the Obukhov length and  $z_m$  the measurement height, or  $z_{ef}/L_{Ob}$  when  $z_{ef}$  has been calculated incorporating the 115 displacement height (Hartogensis et al. 2003). Different values of the empirically derived constants 116 117  $c_{TI-3}$  are used in the literature (Andreas 1988; Hill et al. 1992), and alternative functional forms appear, e.g. Thiermann and Grassl (1992). The Wyngaard (1973) values were adjusted by Andreas 118 119 (1988) to give  $c_{TI}$  = 4.9,  $c_{T2}$  = 6.1 and  $c_{T3}$  = 2.2 (hereafter An88). De Bruin et al. (1993) found  $c_{TI}$  = 4.9, 120  $c_{T2}$  = 9 and  $c_{T3}$  = 0 (hereafter DB93). These similarity functions relate the temperature structure 121 parameter to the temperature scaling variable  $(T_*)$ ,

122 
$$T_* = \left\{ \frac{C_T^2 (z_m - z_d)^{2/3}}{f_{MO}(\zeta)} \right\}^{1/2}.$$
 (4)

123 The friction velocity,  $u_*$ , is estimated from measured wind speed assuming a logarithmic wind profile 124 adjusted for stability and a value for the roughness length (e.g. Stull (1988)). The wind profile 125 equation is solved iteratively with

126 
$$L_{Ob} = \frac{u_*^2 T}{g \kappa_v T_*}$$
, (5)

where g is the acceleration due to gravity and  $\kappa_{\nu}$  is the von Kármán constant (0.4). Note that this equation for  $L_{Ob}$  neglects the buoyancy correction. Finally, the sensible heat flux is obtained using

129 
$$Q_H = -\rho c_p u_* T_*$$
, (6)

130 where  $\rho$  is the density of air and  $c_p$  the specific heat capacity at constant pressure. In the following, 131  $Q_{H\_BLS}$  and  $Q_{H\_LAS}$  are used to denote the sensible heat flux for the 5.5 km path and 2.8 km path, 132 respectively (Sect. 2.2). The sensible heat flux from the EC station is denoted  $Q_{H\_EC}$ .

One drawback of scintillometry is that the sign of the heat flux is unknown and must be assigned based on other information. Following a comparison of possible algorithms, Samain et al. (2012) recommended using an algorithm based on the minima in the diurnal cycle of  $C_n^2$  to indicate a transition of stability. We follow a similar methodology here, which has the key advantage that it is based on path-averaged information.

## 138 2.2. Site description and experimental details

139 This study took place in Swindon (population 175 000), situated 120 km west of London (top 140 right, Fig. 1). Typical of the UK suburban landscape, Swindon consists mainly of residential areas with houses of varying ages extending outwards from the town centre, interspersed with greenspace, 141 142 small parades of shops and institutional buildings. Larger industrial and commercial zones are mostly 143 situated towards the edges of the development. The town centre comprises commercial areas, with some pedestrianized streets, offices, public buildings and transport hubs. Building density in the 144 145 town centre is greater than in the surrounding suburbs and buildings are taller, larger and more 146 variable in height. Outside of the urban core the buildings are more uniform, houses are mostly 1-3 147 storeys, semi-detached or terraced and usually have at least a small garden. There are a few small blocks of flats (4-5 storeys) and larger warehouses in industrial areas. Trees are of a similar height to 148 149 the buildings and found mostly in undeveloped green corridors between residential areas, along roadsides and in gardens. The area is relatively well vegetated (cover fraction 53%), largely due to the prevalence of grassed areas: parks, playing fields, green corridors, gardens, verges and a large nature reserve near the centre of the study area (Fig. 1).





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**Fig. 1** Land cover surrounding the two scintillometer paths (BLS and LAS), eddy covariance station (EC) and two meteorological stations (MET<sub>sub</sub> and MET<sub>roof</sub>). Example footprints for typical atmospheric conditions (wind direction = 225°,  $L_{Ob}$  = -200 m,  $u_*$  = 0.5 m s<sup>-1</sup> and  $\sigma_v$  = 0.9 m s<sup>-1</sup>) are indicated by the cumulative source area: the region within the solid (dashed) line contributes 80% (95%) to the measured flux. The location of Swindon within the British Isles is shown (top right). Details of the land cover classification are given in the text. Where land cover data were unavailable areas are left unclassified (white).

161

162 Observations at multiple scales are achieved through a combination of the eddy covariance 163 technique and two scintillometer paths. The largest measurement scale reaches between the town 164 centre and the rural fringe at the northern edge of the settlement: an infrared scintillometer, the

165 BLS900 (Scintec, Rottenburg, Germany), was installed on a 5.5 km path orientated approximately 166 north-south. A second infrared scintillometer, a LAS 150 (Kipp and Zonen, Delft, The Netherlands), was aligned on a shorter path of length 2.8 km. This path is located over relatively recently 167 developed suburbs (in the last 20 years or so) 3-5 km north of the town centre. Both are large 168 169 aperture scintillometers operating at a wavelength of 880 nm. Although LAS is an abbreviation for 170 large aperture scintillometer, in this study BLS is used to denote the scintillometer on the long path and LAS the scintillometer on the short path. The EC system was installed approximately 3 km north 171 172 of the town centre, close to the middle of the long path.

173 Footprint models can be used to aid the interpretation of observed fluxes by relating them to the 174 probable area of the surface that influenced the measurements. Although some of the assumptions may be challenged by complex environments, footprint models have been used successfully in urban 175 176 areas (Schmid et al. 1991; Järvi et al. 2009; Hiller et al. 2011), providing measurements are made at 177 sufficient height that the influences of individual obstacles or heterogeneities are averaged out. 178 Meijninger et al. (2002b) extended footprint theory to scintillometers by combining source areas 179 calculated for a single point measurement with the scintillometer path weighting function. This has 180 since been adopted by other studies (Meijninger et al. 2006; Hoedjes et al. 2007; Samain et al. 181 2011a; Evans et al. 2012; Liu et al. 2013). A range of footprint models exist; here we use the 182 analytical model of Hsieh et al. (2000) and assume the lateral dispersion is Gaussian (Schmid 1994; 183 Detto et al. 2006).

Results of the footprint model for each of the three systems are shown in Fig. 1 for typical atmospheric conditions (wind direction = 225°,  $L_{Ob}$  = -200 m,  $u_*$  = 0.5 m s<sup>-1</sup> and standard deviation of lateral wind  $\sigma_v = 0.9$  m s<sup>-1</sup>). Source areas vary depending on atmospheric conditions and wind direction, as well as measurement height and surface roughness. The difference in measurement scales is apparent. The sizes of the areas contributing 80% (95%) of the observed fluxes are approximately 0.06, 1.0 and 3.0 km<sup>2</sup> (0.5, 3.0 and 7.5 km<sup>2</sup>) for the EC, LAS and BLS, respectively. The size of the footprints increases with stability.

Beam heights, land cover and building and tree height were obtained using a spatial database incorporating surface cover information (OS MasterMap 2010 ©Crown Copyright), a digital terrain model and digital surface model from lidar (2007, ©Infoterra Ltd) and aerial photography (2009, ©GeoPerspectives). For this study a spatial resolution of 5 m was used, further details are given in Ward et al. (2013a). Some of the residential area at the far north-west of the study area has very recently been completed, with some ongoing development of the rural outskirts during the study period. The overall effect here may be a small overestimation of the vegetated land cover fraction for the LAS path, when winds are from the west or north-west and in stable conditions, as this recentdevelopment has not yet been incorporated in the spatial database.

The roughness length for momentum was estimated based on the mean height of the roughness elements ( $z_H$ ) within the area influencing the measurements using the approximation  $z_0 = 0.1 z_H$ (Garratt 1992). The resulting values (Table 1) are reasonable based on comparison with the literature (Grimmond and Oke 1999), however there is appreciable uncertainty associated with this (and other) methods (Sect. 3.1). The zero plane displacement height,  $z_d$ , is estimated at  $0.7 z_H$  and incorporated into the effective height calculation for the scintillometers, after Hartogensis et al. (2003) (their equation 15).

207 Both infrared scintillometer transmitters were installed on a telecommunications mast at the 208 northern edge of the suburbs at 27.9 m (BLS) and 26.6 m (LAS) above ground level. The LAS receiver 209 was bolted to a 1.7 m high post at the same property as the EC mast, whilst the BLS receiver was 210 mounted on a building in the town centre at 26.2 m. The combination of the topography (Fig. 2), 211 mounting on existing structures and the path weighting provides sufficient beam height for the 212 scintillometers to be above the blending height (which is estimated to be between 15 and 30 m 213 (Pasquill 1974; Garratt 1978)). The blending height will be larger above landscapes with larger scale 214 heterogeneity (e.g. Wood and Mason (1991)), or may not exist at all once the scale of heterogeneity 215 exceeds the boundary layer height (Maronga and Raasch 2013). With these limitations in mind, the 216 Swindon sites were selected such that the patches of different land cover within the footprints are, 217 for the most part, reasonably small (100 - 300 m or less).

Instrumentation	Dates	Location	<i>z<sub>m</sub></i> [m]	Zef	Path	Bearing	Z0	Zd
				[m]	length [m]	[°]	[m]	[m]
BLS	12 Jan 2011 –	51°36'33.9" N	44.3	45.0	5492	170	0.7	4.9
	31 Dec 2012	1°47'38.6" W (Tx)						
		51°33'38.1" N						
		1°46'55.3" W (Rx)						
LAS	22 Jun 2011 –	51°36'33.9" N	32.4	35.9	2761	184	0.6	4.5
	31 Dec 2012	1°47'38.6" W (Tx)						
		51°35'4.9" N						
		1°47'53.0" W (Rx)						
EC	09 May 2011 –	51°35'4.6" N	12.5	-	-	-	0.5	3.5
	31 Dec 2012	1°47'53.2" W						
	09 May 2011 –	51°35'4.6" N	10.6 (WXT)	-	-	-	0.5	3.5
	31 Dec 2012	1°47'53.2" W	10.1 (NR01)					
MET <sub>roof</sub> *	01 Jan 2011 –	51°34'0.3" N	2.0 (WXT)	-	-	-	-	-
	31 Dec 2012	1°47'5.3" W	1.1 (NR01)					

218 Table 1 Details of the instrumental setup. Tx denotes transmitter, Rx receiver. For the scintillometers the mean heights of

219 the beams above the surface  $(z_m)$  and the effective measurement heights  $(z_{ef})$  are given. The date range refers to the data

used here. \*For MET<sub>roof</sub> the heights above the roof surface are given;  $z_0$  and  $z_d$  were not calculated for this site.



Fig. 2 Cross section of the topography (solid black line) and mean obstacle height (dotted line; buildings and trees within a
 radius of 100 m) along the (a) BLS and (b) LAS paths (coloured lines).

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A CR5000 datalogger (Campbell Scientific Ltd., Loughborough, UK) sampled the intensity of the received LAS beam at 500 Hz and measured the  $C_n^2$  signal (calculated onboard the instrument and stored as a logarithm) every second which was then output at 1 min intervals and these were averaged to 10 min. For the BLS, the mean and standard deviation of the beam intensity of each disk (the BLS900 is a dual disk instrument with two transmitting apertures) were obtained from the supplied signal processing unit at 30 s intervals, then converted to log-amplitude variances and  $C_n^2$ and averaged up to 10 min. Data from a single disk are presented here.

The EC system consists of a sonic anemometer (R3, Gill Instruments Ltd., Lymington, UK) and an open-path infrared gas analyser (LI-7500, LI-COR Biosciences, Lincoln, USA) at a height of 12.5 m. As this is 2-3 times the height of the surrounding buildings and trees, it is therefore sufficiently high to deliver fluxes representative of the local-scale. Data were processed using EddyPro (LI-COR) following conventional procedures. Further details of the EC measurements can be found in Ward etal. (2013a).

238 Meteorological instruments were installed on the same mast as the EC equipment (denoted MET<sub>sub</sub>). A second set of meteorological data were collected at a rooftop site close to the town 239 240 centre (MET<sub>roof</sub>). Both stations included a four-component radiometer (NR01, Hukseflux, Delft, The 241 Netherlands), automatic weather station (WXT510, Vaisala, Helsinki, Finland) and tipping bucket rain 242 gauge (0.2 mm tip, Casella CEL, Bedford, UK). At MET<sub>sub</sub>, the radiometer was installed at a height of 243 10.1 m so that the downward-facing field of view comprises a mixture of surfaces: grass lawns and 244 verges, road, pavement, hedges and small trees, bare soil, gravel, roofs of garages, small sheds and 245 single-storey extensions and walls (brick and painted). At MET<sub>roof</sub>, the radiometer was installed at 1.1 m above the roof surface made of grey synthetic material and black rubber matting. Additionally at 246 247 MET<sub>roof</sub>, a heat flux plate (HFP01, Hukseflux) was installed between the roof surface and rubber 248 sheet, providing an approximation of the change in storage through the roof. At both sites the 249 meteorological data were logged at 1 min intervals (CR1000, Campbell Scientific Ltd.) and 250 subsequently averaged to obtain 10 min resolution for calculation of the scintillometer fluxes or 30 251 min for comparison with EC fluxes. The 10 min scintillometer fluxes were also averaged to 30 min for 252 comparison with EC results. Details of the observational setup are summarized in Table 1.

To provide nearly continuous auxiliary data required for scintillometry processing, results from the two meteorological stations were combined. When available, data from MET<sub>sub</sub> are used as the siting of this station is more appropriate. Based on the regression of concurrent data (9 May 2011 -31 Dec 2012), temperature, relative humidity (RH), pressure and wind speed at MET<sub>roof</sub> were adjusted to gap-fill the combined dataset, including the period prior to installation of MET<sub>sub</sub> on 9th May 2011. This is considered further in Sect. 2.3.

All data were subject to quality control routines. Data were removed at times of known instrument malfunction. Meteorological data were excluded when they (or their standard deviations) fell outside physically reasonable thresholds. Quality control of the scintillometry data included rejecting times when the received signal intensity dropped below half of the value in clear conditions, which usually indicated rain or fog. Data points neighbouring those that failed the intensity check were also removed. Out of the total data collected, 84% of BLS and 82% of LAS data (10 min) remained for analysis.

266 Both scintillometers were corrected for the effects of saturation using the modulation transfer 267 function of Clifford et al. (1974). Using the threshold value suggested by Kleissl et al. (2010), 16% of the BLS data and 0.2% of the LAS data might be expected to suffer from saturation. Overall, the correction increased  $C_n^2$  by 4% and 1% for the BLS and LAS respectively (naturally the corrections are larger with increased scintillation and rise to 8% and 2% for the midday periods).

271 Recent studies have indicated sometimes severe discrepancies between certain scintillometers, 272 in particular the LAS 150 model (Kleissl et al. 2008), whereas the BLS900 model tends to give more reproducible results (Kleissl et al. 2009b). Prior to deployment in Swindon, the LAS and BLS were run 273 274 alongside each other at a fairly homogenous grass test site at Chilbolton Observatory, Hampshire, UK (17 April 2010 – 25 May 2010). Observed  $C_n^2$  ranged between  $10^{-16}$  m<sup>-2/3</sup> and  $10^{-12}$  m<sup>-2/3</sup>, which 275 spans the range of values observed for the Swindon paths. Results suggested the response of the 276 LAS is reasonable but compared to the BLS  $C_n^2$  is overestimated by 9.8%. This adjustment has been 277 applied to the LAS  $C_n^2$  for the Swindon data. As a result of these comparisons (e.g. Kleissl et al. 278 (2008; 2009b), Van Kesteren and Hartogensis (2011)), Kipp and Zonen have updated their original 279 280 LAS 150 instrument to a LAS MkII model (Mustchin et al. 2013).

#### 281 2.3. Assessment of the input meteorological data

282 First, the suitability of the combined meteorological input data used to process the scintillometry 283 fluxes is considered. To calculate  $Q_H$  from single-wavelength scintillometry, air temperature, pressure and humidity are required to first obtain the structure parameter of temperature,  $C_T^2$ . Both 284 285 T and RH are similar between the MET<sub>sub</sub> and unadjusted MET<sub>roof</sub> sites. The regression slopes are within 3% and there is high correlation ( $r^2 > 0.98$ ). Sensitivity of  $Q_H$  to these input meteorological 286 287 variables is small (Hartogensis et al. 2003) and indeed these very small differences have minimal 288 impact on the fluxes. The average difference in  $Q_H$  is < 0.5% when calculated using T, RH and p from 289 each site. Use of this combined dataset is therefore judged unproblematic and to be a sufficiently 290 accurate representation of T, RH and p across the study area.

291 An initial estimate of the Bowen ratio is recommended to account for the contribution of humidity and combined temperature-humidity fluctuations to optical  $C_n^2$  (Wesely 1976). Usually the 292 293 value of  $\beta$  is arrived at iteratively through incorporation of the available energy (e.g. Meijninger et al. 294 (2002b)). However, estimating the available energy is challenging in urban areas as the net storage 295 heat flux ( $\Delta Q_s$ ) plays a more significant role in the energy balance than for most rural areas (e.g. 296 grassed or agricultural land), yet it is very difficult to measure directly (Offerle et al. 2005; Roberts et 297 al. 2006). Other (rural) studies have used  $\beta$  measured at a nearby station (Hoedjes et al. 2002; 298 Samain et al. 2011a) or have calculated  $Q_H$  using a series of values of  $\beta$  (Meijninger and De Bruin 299 2000). When  $\beta$  is expected to be large (e.g. > 0.6 for Chehbouni et al. (2000b); > 1 for Moene (2003)) the correction may be neglected. Given the uncertainty in estimating the available energy and the lack of representative EC data across the whole study area, the Bowen ratio correction has not been applied for the results presented here. The potential impact is an average overestimation in  $Q_H$  from the scintillometers of less than 5% for  $\beta > 1$ , and less than 10% for  $\beta = 0.5$ . For the BLS, the  $C_T^2$  values here were found to be within around 6% of the  $C_T^2$  values calculated incorporating data from the millimetre-wave scintillometer (Sect. 1), which do not require a Bowen ratio correction (see Ward et al. (in preparation b) for details).

To process scintillometry data, the friction velocity is usually estimated from wind speed measured at a single point and adjusted to beam height using the logarithmic profile accounting for stability. As with the other meteorological inputs, wind speed from MET<sub>roof</sub> was adjusted to produce the combined dataset with optimum availability of input data. Concurrent  $Q_H$  values calculated using the MET<sub>sub</sub> wind speed or the adjusted MET<sub>roof</sub> wind speed differ by less than 3%,  $r^2$  is high (0.98) and there is little scatter (root mean squared error, RMSE < 10 W m<sup>-2</sup>).

The dual-disk design of the BLS900 enables estimation of the path-averaged crosswind, i.e. the 313 314 component of the wind speed perpendicular to the scintillometer path. To check that the point 315 measurements of wind speed were a realistic proxy for the wind field over the scintillometer source 316 area, a comparison was made between the BLS crosswind speed and the equivalent crosswind speed 317 calculated using wind speed and direction from MET<sub>sub</sub> and scaled to the effective height of the BLS using stability from the EC station. Overall the crosswind estimates displayed similar trends across a 318 range of wind speeds and directions. The high correlation obtained ( $r^2 = 0.922$ ) implies that these 319 320 point measurements generally capture the variability of the wind field at the larger scale and gives 321 confidence in their use in processing the scintillometry data.

## 322 **3.** Analysis of sensible heat fluxes

### 323 **3.1. Assessment of seasonal cycles and annual variations**

324 Large-area sensible heat fluxes from the 5.5 km scintillometer path are presented for two years 325 (2011-12), alongside 18 months of data from the shorter 2.8 km scintillometer path and almost 20 months of eddy covariance data (Fig. 3). The annual cycle is evident, with mean daily (24 h)  $Q_H$ 326 327 reaching a maximum in early summer (May-June in 2011, May in 2012) and minimum in December. In December, average  $Q_H$  is negative even during daytime (defined as times when incoming 328 shortwave radiation  $K_{\perp}$  > 5 W m<sup>-2</sup>, see Fig. 3b) as the typical diurnal course of  $Q_H$  becomes positive 329 330 only for a few hours around midday (Fig. 4). This behaviour is observed consistently across the three datasets and in both years but contrasts with the majority of urban studies in more built-up areas, 331

where greater heat release from storage and anthropogenic activity can maintain a positive sensible heat flux all year round (Offerle et al. 2005; Kotthaus and Grimmond in press-a). For two sites in Oberhausen, Germany, Goldbach and Kuttler (2013) found  $Q_H$  to be positive for most of the daytime throughout the year at their urban site, whereas their suburban site exhibits similar behaviour to Swindon. Besides the smaller storage and anthropogenic heat flux in suburban areas, more of the available energy is partitioned into evaporation, owing to increased moisture availability from soil surfaces and greater total evapotranspiration from a larger vegetation fraction.

339 Within the trends of the expected annual cycle, there are notable differences between the two 340 years studied. Peak  $Q_H$  is larger in summer 2011 compared to 2012 and month-to-month variation is 341 smaller in 2011. Broadly speaking, much of 2011 was under threat of drought, with dry soil moisture 342 conditions and depleted ground water supplies. Despite frequent rain and very few clear sky days 343 the annual rainfall total was 530 mm compared to the average 780 mm for southern England<sup>1</sup>. Dry 344 conditions continued through early spring 2012, until early April. Very wet weather followed and remained throughout 2012 (total rainfall 1020 mm), with brief drier and warmer spells in late July 345 346 and early September.

June 2012 was particularly wet and cloudy (sunshine hours were only 70% of normal<sup>1</sup>; mean 347 daytime  $K_{\perp}$  was 174 W m<sup>-2</sup> in 2012 compared to 212 W m<sup>-2</sup> in 2011). Monthly mean daily  $Q_H$  was 348 19%, 34% and 31% lower than in June 2011 for the BLS, LAS and EC, respectively. August 2012 also 349 350 had a notably lower  $Q_H$  during daytime compared to 2011 (also shown in Fig. 4), despite similar 351 radiative energy inputs in both years. A long dry spell and generally sunny weather in September 352 2012 allowed surfaces to dry out and  $Q_H$  to increase, resulting in a larger average value than for the 353 previous month. Large negative nocturnal  $Q_{H BLS}$  in February 2011 means the daily (24 h) average is 354 negative, whereas high  $Q^*$  and high daytime  $Q_{H BLS}$  in 2012 contribute to a positive 24h average in 355 2012 (Fig. 3).

<sup>&</sup>lt;sup>1</sup> Met Office climate statistics (1971-2000), <u>http://www.metoffice.gov.uk/climate</u>, last accessed 29 March 2013





**Fig. 3** Monthly mean sensible heat flux observations from scintillometry (BLS and LAS) and eddy covariance (EC) for all available data (a) over 24 h and (b) separated into day ( $K_{\downarrow} > 5$  W m<sup>-2</sup>) and night times. Partial months in relation to the installation dates (Table 1) are January 2011 (BLS), May 2011 (EC) and June 2011 (LAS, note only 4 days of data due to an instrument fault). Error bars indicate the impact on the scintillometer fluxes of altering the input roughness length by ± 0.2 m (a) or using the similarity functions of De Bruin et al. (1993) (b). The net radiation is indicated by shading (b, right-hand axis).



Fig. 4 Temporal variation of monthly mean diurnal cycles of sensible heat fluxes from (a) eddy covariance, (b) the LAS and(c) the BLS.

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370 Overall there is remarkably good agreement across the three datasets, which capture seasonal 371 similarities and inter-annual variability. The different source areas of each instrument, and that the 372 BLS measures across a large proportion of northern Swindon, suggest that these trends are local-to-373 city-scale responses to regional weather variability. Furthermore, this agreement implies that any 374 bias in the monthly averages due to the effect of the wind direction distribution on the 375 measurement footprints is outweighed by the changes in surface conditions, prevailing weather and 376 the resulting surface-atmosphere interactions. Given the much smaller source area of the EC 377 technique compared to scintillometry, it is reasonable to expect that heterogeneity of the surface 378 has a larger influence on the EC observations than the scintillometry observations (Sect. 3.3).

The LAS tends to give the largest  $Q_{H}$ , particularly during daytime, compared to both EC and the BLS. In summer 2012 the EC and BLS average values do not reach above 150 W m<sup>-2</sup>, in contrast to  $Q_{H}_{LAS}$  (Fig. 4). During winter months (November 2011-January 2012) and at night the BLS gives the largest fluxes. Daily average  $Q_{H\_EC}$  often lies between the two scintillometer averages but during winter (November-December 2011, December 2012) and at night the scintillometers tend to give larger magnitude  $Q_{H}$ . This can also be seen in Fig. 4: the absolute size of  $Q_{H}$  from the scintillometers is larger (e.g. around transition times in December), whether positive or negative, whereas EC values are much closer to zero. Larger scintillometer fluxes in neutral-to-stable conditions may reflect the performance of the similarity functions (Sect. 3.2).

388 The widely implemented similarity functions of Andreas (1988) were used here. Using the De 389 Bruin et al. (1993) similarity functions instead increases  $Q_H$  by about 13-14% (bars in Fig. 3b). This is 390 similar to results in Marseille (Lagouarde et al. 2006) and within the 10-15% range given by Beyrich 391 et al. (2012). The large uncertainty introduced by the choice of similarity function is a major 392 limitation of the scintillometry technique across all environments; it is not confined to urban sites 393 although there is the added question of whether functions developed over homogeneous terrain 394 should be applied to more heterogeneous locations. Kanda et al. (2002) and Roth et al. (2006) both 395 derived 'urban forms' of the similarity functions for their small aperture scintillometer studies, 396 however their paths were closer to, or within, the roughness sub-layer. Other large aperture studies 397 in urban environments have used the more common functions (Lagouarde et al. 2006; Zieliński et al. 2012). 398

399 Typically, the uncertainty in  $z_0$  is large as  $z_0$  can vary spatially, with time of day and stability 400 (Grimmond et al. 1998; Hoedjes et al. 2007; Zilitinkevich et al. 2008), and with shape, density and 401 arrangement of surface structure (Grimmond and Oke 1999). For the study area, the true value is 402 expected to be within the range 0.4 to 1.0 m based on values in the literature. The impact on the 403 scintillometer estimation of  $Q_H$  of changing the prescribed values of  $z_0$  by ±0.2 m is ±7% (error bars 404 in Fig. 3a). Although the flux is fairly sensitive to the value of  $z_0$  used, the overall trends do not 405 change significantly. No adjustment was made to account for seasonal variation in  $z_0$  (or  $z_d$ ), though 406 these values may be 10-20% smaller in winter than in summer (Grimmond et al. 1998). Using a 407 smaller value of  $z_0$  during leaf-off periods decreases the wintertime fluxes slightly (the error bars in 408 Fig. 3a represent a change in  $z_0$  of about ±30%).

Allowing a  $\pm 5\%$  uncertainty in  $z_{ef}$  ( $\pm 2.25$  m) affects the fluxes by  $\pm 3\%$ . This uncertainty in  $z_{ef}$ includes measurement accuracy and variation of the effective height with stability as well as accounting for spatial differences in obstacle height (hence  $z_d$ ) and topography. The large beam height and relatively small displacement height help to keep the sensitivity to  $z_{ef}$  (and  $z_d$ ) small.

#### 413 **3.2. Short-term variability**

Direct comparison of 30 min sensible heat fluxes obtained from scintillometry and EC (Fig. 5) 414 415 indicates reasonably good agreement between the measurement techniques and across the scales with strong correlation ( $r^2 \approx 0.87$ ). The slope of the regression between  $Q_{H_{LAS}}$  and  $Q_{H_{EC}}$  is close to 416 1, with a small positive offset, whereas the BLS tends to give lower  $Q_H$  than EC particularly towards 417 large values of  $Q_{H}$ . Whilst the linear fit between  $Q_{H_{LAS}}$  and  $Q_{H_{EC}}$  indicates a good match between 418 419 these data, the BLS data distribution appears more curved at high  $Q_{H\_EC}$ . The source area of the EC 420 mast and BLS are quite different, which may partly explain why the highest EC fluxes are not 421 matched by the BLS. Specifically, the area to the south and south-west of the EC mast has a 422 particularly high proportion of built and impervious surfaces and little vegetation, whilst the BLS 423 source area always includes open green spaces. Thus when the EC footprint is over the least 424 vegetated sector (180-240°), the measured  $Q_{H EC}$  tends to be larger compared to other wind sectors around the flux mast as well as to the scintillometer results. This effect would be amplified when 425 426 surface water is scarce. During summertime, when the wind is from the south, both  $Q_{H_{LAS}}$  and  $Q_{H_{BLS}}$  are lower than  $Q_{H_{EC}}$  (Fig. 6). For more westerly winds (240-270°) the EC source area contains 427 more vegetation and there is closer agreement between  $Q_{H_{EC}}$  and  $Q_{H}$  from the scintillometers. 428

429 However, the remaining curvature in Fig. 6a, but not seen in Fig. 6b, most likely indicates 430 saturation affecting the BLS (5.5 km path) but not the LAS on the shorter path. Despite having corrected the scintillometers for saturation (Sect. 2.2), a comparison of the distribution of  $C_T^2$  values 431 from the BLS and LAS suggests that the highest BLS fluxes are still affected: whilst the LAS provides 432  $C_T^2$  values up to about 0.03 K<sup>2</sup> m<sup>-2/3</sup>, the BLS distribution drops off sharply at around 0.009 K<sup>2</sup> m<sup>-2/3</sup>. 433 Recently, Wood et al. (2013) found an upper  $C_T^2$  threshold of 0.02 K<sup>2</sup> m<sup>-2/3</sup> for their shorter path 434 length of 4.2 km. Other studies have also suggested that the effects of saturation may still be 435 observed above commonly-used thresholds (Kohsiek et al. 2006). 436

During night and transition times, the agreement between the datasets is poorer ( $r^2$  decreases to 437 around 0.4 for  $K_{\downarrow} \leq 5$  W m<sup>-2</sup>). This is to be expected for several reasons. Firstly, the limitations in 438 439 instrument performance are reached when fluxes are small, for both EC and scintillometry. 440 Secondly, the time of stability transition may vary with location, even along the scintillometer paths, so that the three values of  $Q_H$  obtained for a given time period may not have the same sign. Data 441 points in the second and fourth quadrants indicate when scintillometer and EC derived  $Q_H$  have 442 443 opposite signs. The stability may also change more than twice per day which would mean the scintillometer data are processed assuming the wrong stability regime. Thirdly, the corrections for 444 the influence of humidity fluctuations on  $C_T^2$  and  $L_{Ob}$  are generally larger at these times (when  $\beta$  is 445

small). The Bowen ratio correction to  $C_T^2$  introduces the larger error of these two approximations; 446 447 neglecting the buoyancy correction to the Obukhov length (e.g. Green et al. (2001)) is thought to lead to a slight underestimation in  $Q_H$  of  $\approx 0.5$  W m<sup>-2</sup>. Finally, near-neutral to stable atmospheric 448 449 conditions do not always satisfy the assumptions required for the measurement theory (e.g. weak 450 turbulence, non-stationarity, poorer performance of similarity functions). Removing the night time 451 data causes the regression slopes in Fig. 5 to decrease slightly to 0.77 (BLS) and 0.94 (LAS), and the intercepts to increase to 13 W m<sup>-2</sup> (BLS) and 9 W m<sup>-2</sup> (LAS). For night time data only, the intercepts 452 are similar in size but of opposite sign. These intercepts are thought to result from the 453 454 overestimation of small fluxes by the similarity functions. Considering all data together (Fig. 5) the 455 lack of small  $Q_H$  values from the scintillometers can be identified around zero. Using functions of a conventional form (such as Equations 2 and 3) appears to under represent  $Q_H$  values close to zero 456 457 and overestimates  $Q_H$  in neutral conditions ( $f_{MO}$  is too small so the  $T_*$  obtained is too large). Investigation into the scaling of  $C_T^2$  with stability is presented in more detail elsewhere (Ward et al. 458 in preparation a) and Lagouarde et al. (2006) also noted an overestimation (15 W m<sup>-2</sup>) of small night 459 time  $Q_H$  values using An88 and DB93 (unstable forms). Although this effect is undesirable, the small 460 461 size of the fluxes at these times means that absolute errors are small.

462



464 Fig. 5 Comparison of 30 min sensible heat fluxes derived from the scintillometers (BLS, LAS) and eddy covariance (EC) for all

465 available data.



467 Fig. 6 As for Fig. 5 but for summertime (May-Sep 2011-12) data only and for wind directions 180-270° (colours).

466

469 The diurnal course of  $Q_H$  obtained from the three systems follow each other closely: example 470 days from July 2012 are shown in Fig. 7. No rainfall was observed during these mostly clear-sky days 471 although the influence of cloud cover can be seen on the morning of 22 July and afternoon of 25 472 July. On 22 July the fluxes respond consistently to changes in the net radiation and the peaks and 473 troughs are closely matched between EC, LAS and BLS observations. Data from MET<sub>roof</sub>, 474 approximately 3 km southwards (Fig. 1), closely matches the variation in  $Q^*$  measured at the EC site. 475 Time-lapse photography reveals fairly uniform, almost full cloud cover at sunrise which clears 476 throughout the morning. On the afternoon of 25 July, however, the situation is quite different. 477 Rapidly changing patchy cloud cover creates spatial variability in the radiation balance components. The responses of the two radiometers are less well correlated (compare  $Q^*$  and  $Q^*_{roof}$  in Fig. 7a) and 478 479  $Q_H$  is seen to respond differently across the different measurement scales. Not surprisingly,  $Q_{H EC}$ 480 most closely matches  $Q^*$  as both are measured at the same location and have more similarly sized 481 and at least partially coincident source areas. In general, the scintillometers yield a more smoothly 482 varying diurnal course than EC, often attributed to the greater spatial averaging by scintillometers 483 (e.g. Lagouarde et al. (2006), Guyot et al. (2009)). The BLS appears to vary more smoothly than the 484 LAS (e.g. 24 July) which is consistent with the size of their source areas.

For clear days, the phase of  $Q_H$  is lagged relative to  $Q^*$ . At the three scales  $Q_H$  peaks after  $Q^*$  and remains positive later into the evening than  $Q^*$ . One component of the urban net heat storage flux is approximated by a heat flux plate installed under the roof covering at MET<sub>roof</sub> ( $\Delta Q_{S_{roof}}$  in Fig. 7a). This flux increases earlier in the day and becomes negative long before  $Q^*$ . In this way, release of 489 stored energy enables  $Q_H$  to remain positive even when  $Q^*$  is negative (Oke and Cleugh 1987; 490 Lemonsu et al. 2004). Normalising these fluxes by the net radiation clearly demonstrates the 491 opposing hysteresis patterns of  $Q_H$  compared to  $\Delta Q_{S roof}$  (Fig. 7b). The proportion of  $Q^*$  directed into 492 sensible heat increases throughout the day whereas the proportion of energy used to heat the 493 surface decreases. Strong hysteresis is evident on clear days but it tends to be less obvious on 494 cloudier days. Similar patterns have been observed at other urban sites at the local-scale (Grimmond 495 and Cleugh 1994; Grimmond and Oke 2002; Grimmond et al. 2004). Here we demonstrate that the phase lag between  $Q_H$  and  $Q^*$  is observed right across the urban environment, from the local-scale 496 497 up to the city-scale. The shift in peak  $Q_H$  around midday and into early afternoon can also be seen to 498 some extent in the average monthly values (Fig. 4), particularly in spring and early summer 2011.

499 Other than under conditions of rapidly changing  $Q^*$ , and its associated high spatial variability, the 500 diurnal patterns in  $Q_H$  derived from EC and the scintillometers match those of  $Q^*$  measured at the 501 EC site (Fig. 8). On 21, 22 and 27 July 2011 the sudden drop in  $Q^*$  during the middle of the day is also 502 seen in  $Q_{H}$ . The day-to-day variation in these two quantities is also very similar. For example  $Q^*$  and  $Q_H$  steadily increase to reach over 600 W m<sup>-2</sup> and 200-300 W m<sup>-2</sup>, respectively, on 25 July when peak 503  $Q_{H BLS}$  is about 2/3 of  $Q_{H LAS}$ . Both  $Q^*$  and  $Q_{H}$  are lower during the following few days until 29 July 504 when the net radiation remained very small throughout the day (< 100 W  $m^{-2}$ ) and conditions were 505 506 mostly near-neutral. On this day the agreement in the shape of the diurnal cycle between the scintillometers and EC is poorer, although the fluxes show some agreement in responding to the dip 507 508 in  $Q^*$  in the afternoon. Under these near-neutral conditions it is likely that the stability transitions 509 occur more often than twice daily as prescribed by the algorithm used to determine the sign of 510 scintillometric  $Q_{H}$ . Indeed,  $Q_{H_{EC}}$  is seen to change sign several times during the afternoon and 511 evening.



513

**Fig. 7** Diurnal variation in sensible heat fluxes  $(Q_H)$  and net all-wave radiation  $(Q^*)$  for four days in July 2012. Data from a heat flux plate installed on a rooftop, representing one component of the storage heat flux  $(\Delta Q_{S_roof})$  and a second radiometer located on the rooftop  $(Q^*_{roof})$  are also shown. In (b) the fluxes have been normalized by the net all-wave radiation measured at the EC site  $(Q^*)$ .



**Fig. 8** Sensible heat fluxes from EC and the scintillometers alongside net all-wave radiation from the EC site ( $Q^*$ ), rainfall and wind direction (also measured at the EC site) for two weeks in July-August 2011.

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The sign of the scintillometer sensible heat flux must be assigned during processing. Here, the 523 stability was assumed to change from stable to unstable at the first minimum in  $C_n^2$  on each day, and 524 from unstable to stable at the second minimum, providing these transitions occurred within the 525 526 likely time frames for sunrise and sunset. Additionally, the net radiation can be used to check 527 whether the minima identified are likely to indicate stability transitions rather than sudden increases in cloud cover, for example. For each 24 h period the algorithm always results in some stable and 528 some unstable data and the proportion of each depends on the observed behaviour of  $C_n^2$ 529 530 (effectively on the time between the morning and evening minima). As is evident from the data, this method generally performs well in Swindon, where EC data suggests  $Q_H$  tends to be positive for 531 some duration around midday and negative at night (Ward et al. 2013a). However, there are some 532 533 days when the stability transition does not occur and either unstable conditions prevail throughout 534 the night or stable conditions throughout the day. In these cases the sign of the fluxes from the scintillometers may be incorrect but these occasions are observed infrequently and the size of thefluxes tends to be small so the likely impact is minimal.

537 The day-to-day (night-to-night) changes in amplitude are usually captured (e.g. decreasing magnitude of nocturnal  $Q_H$  24-27 July 2011 in Fig. 8b) and for some days the evolution of  $Q_H$ 538 539 throughout the night is similar (e.g. decreasing 20-21, increasing 25-26 and 26-27 July 2011, Fig. 8b). 540 This clear relation between the scintillometer and EC fluxes gives confidence that the measurement 541 heights are suitable; in particular that the scintillometers are not measuring above the surface layer. 542 In the winter months, occasionally there are periods of a few hours to days when the shallow surface 543 layer means the scintillometer measurements cannot be related to surface fluxes via MOST (Braam 544 et al. 2012). The EC data further supports these findings with very few cases of strongly stable 545 stratification observed ( $\zeta_{EC}$  < 0.1 for 89% of data).

#### 546 **3.3. Influence of the surface**

547 Comparing the relative sizes of the fluxes can offer insight into key controls on suburban energy partitioning. Towards the end of the case study in Fig. 8 (30 July-01 August 2011),  $Q_{H EC}$  peaks at 548 larger values than either of the scintillometers, whilst  $Q_{H LAS}$  is generally largest near the beginning 549 550 of the period (21-25 July). The wind direction (Fig. 8c) provides a partial explanation due to the 551 variation in source areas. For westerly to northerly winds,  $Q_{H LAS}$  tends to be largest. All three fluxes 552 become similar during northerly winds, when there is a greater vegetation fraction within the source 553 area of each instrument. For the scintillometers the footprint will extend to include some of the rural farmland beyond the edge of the suburbs; at the EC site the increased vegetation is due to more 554 555 gardens to the north of the mast (Ward et al. 2013a).

The period shown in Fig. 9 (21 May – 31 July 2012) coincides with sudden vegetation growth in response to warm, sunny conditions at the end of May, completing the leaf-out period to reach maturity. Vegetation is then fully active throughout June and July. In this period a range of synoptic conditions (cloudy, mixed and clear days), frequent rainfall and a wide distribution of wind directions (although south-westerly was still dominant) occurred.

Footprint calculations for each 30 min period reveal an overall ranking of the vegetation fraction for each instrument that is in accordance with broad expectations given their respective sitings (EC < LAS < BLS). The mean vegetation fractions ( $\pm$  standard deviations) are 44.1 ( $\pm$ 5.0) %, 53.9 ( $\pm$ 2.9) % and 56.9 ( $\pm$ 4.5) % for EC, LAS and BLS, respectively, for the data shown in Fig. 9. The standard deviation is largest for the EC site, as might be expected (a) given the far smaller size of the source area and (b) the differences in surface cover with wind sector around the mast. The vegetation 567 fraction ranges between 32.6% and 56.8% according to the EC footprint estimation for this period. 568 The LAS source area characteristics are much less variable (minimum 47.7%, maximum 60.2%). The 569 retail park to the west of the path (Fig. 1) constitutes a small proportion of the total source area and 570 for westerly wind directions there is only a small increase in the built and impervious fractions. 571 Despite having the largest area, the BLS footprint shows appreciable variability (48.3% - 65.7%). 572 mostly associated with southerly or northerly winds when the town centre and nearby industrial 573 areas (Fig. 1) or rural surroundings are included in the footprint. For small changes in wind direction 574 the BLS source area composition hardly changes, whereas the EC source area composition can vary 575 considerably (particularly for the 180-270° sector). In addition to the directional aspect of the 576 surface heterogeneity, the total area included in the scintillometer footprint is smaller when the 577 wind direction is parallel, as opposed to perpendicular, to the scintillometer path (Meijninger et al. 578 2002b). In this case, the spatial integration occurs over a smaller area so the footprint composition, and observed fluxes, may be expected to be more variable. 579

580



**Fig. 9** Ratio of observed sensible heat flux to net all-wave radiation versus the proportion of vegetation within the flux footprint of the EC station, LAS and BLS in Swindon. Points are 30 min values around midday (1100-1500 UTC) for the period 21 May – 31 July 2012. Data are excluded for times during and  $\leq$  2 h after rainfall and when  $K_{\downarrow} \leq$  200 W m<sup>-2</sup>. Black symbols with error bars represent the mean ± standard deviation of the respective observed values binned in 5% intervals of the vegetated cover fraction (bins with > 10 data points are plotted). Those data collected more than 2 days since rainfall are outlined in red. Average summertime values from various sites in the literature are shown for comparison (see references for details).

590 The ratio of  $Q_H$  to  $Q^*$  decreases as the proportion of vegetation within each instrument's source 591 area increases (Fig. 9). Normalising the turbulent fluxes by an indicator of the energy available 592 largely removes the otherwise often dominant dependence on insolation. Additionally, to moderate 593 the influence of the diurnal hysteresis pattern (Fig. 7), only data around midday (1100-1500 UTC) 594 have been included in Fig. 9. The observed relation between vegetation cover and partitioning of 595 energy into  $Q_H$  is in agreement with other published studies, including summertime data from 596 Kansas City (Balogun et al. 2009), seven sites in Basel (Christen and Vogt 2004), two sites (high and 597 medium density) in Melbourne (Coutts et al. 2007) and four sites in Łódź (Offerle et al. 2006). Use of 598 the scintillometers in Swindon enables this comparison to be extended to larger scales.

599 Relations between land cover and energy partitioning have mostly been developed for summer 600 months, when the majority of field campaigns have taken place and do not account for differences 601 in surface or synoptic conditions. Whilst there is generally good agreement between summertime 602 datasets across a range of sites, those studies extending to winter demonstrate very different 603 behaviour of  $Q_H/Q^*$ . In dense urban areas, the anthropogenic heat flux and much larger storage flux 604 can sustain a positive sensible heat flux all year round (Goldbach and Kuttler 2013; Kotthaus and 605 Grimmond in press-a). In these locations, building density may be a more appropriate variable to use than vegetation fraction and the effect of the anthropogenic heat flux can result in  $Q_H$  that is 606 607 significantly greater than  $Q^*$ . The few campaigns spanning multiple seasons indicate temporal evolution of daytime  $Q_{\rm H}/Q^*$ , e.g. between about 0.30 (winter) and 0.55 (summer) in Melbourne 608 609 (Coutts et al. 2007), and between 0.29 (December) and 0.49 (July) in Tokyo (Moriwaki and Kanda 2004). The data presented here reveal daytime  $Q_{H}/Q^{*}$  peaks in spring between 0.4 and 0.5 and 610 drops to about 0.2 in winter for Swindon. These seasonal changes incorporate multiple effects. The 611 612 anthropogenic influences already mentioned, vegetative activity and the amount of incoming 613 radiation are major factors, but do not account for inter-annual variability in meteorological 614 conditions or rainfall. In February 2012 the limited moisture availability likely contributed to an 615 atypically high  $Q_H/Q^*$  of around 0.4.

At shorter time scales, the meteorological conditions and local stability both have an influence. Reduced availability of moisture constrains the latent heat flux and allows the sensible heat to rise. Following rainfall, the surface dries out and the ratio  $Q_H/Q^*$  tends to increase (outlined points in Fig. 9 represent data collected following more than 2 days without rainfall). Inter-annual variations in rainfall can lead to differences in the size of the fluxes from year to year that cannot only be attributed to variations in  $Q^*$  (Fig. 3). Although normalising by  $Q^*$  removes much of the dependence 622 on the radiative energy, whether conditions are clear or cloudy can affect the response of the surface. Some studies have stratified results by cloud cover conditions (Grimmond and Oke 1995; 623 Balogun et al. 2009) although the effect on  $Q_H/Q^*$  is small. In Fig. 9, data are excluded for  $K_{\perp} \leq 200$ 624 W m<sup>-2</sup> and most of the remaining points greater than 0.6 occur under low insolation. For large  $K_{\downarrow}$ 625 values the scatter is further reduced; this likely to be a result of differing conditions within the 626 627 instruments' source areas under variable cloud cover. The sensible heat flux is dependent on the amount of energy stored and released, which itself depends on the season (Offerle et al. 2005), 628 629 surface wetness (Kawai and Kanda 2010) and cloud cover (Grimmond and Oke 1995). The ability of 630 the surface to store or dissipate heat depends primarily on the physical properties of the constituent 631 materials, but may also be affected by changes in surface conditions, for example a wet surface may have a lower albedo than when dry (e.g. in Cairo (Frey et al. 2011)) and soil moisture affects its 632 633 conductivity. Different materials respond differently to direct and diffuse radiation (Kotthaus and 634 Grimmond in press-b). In combination with surface morphology and changing solar elevation with 635 latitude and time of year, this determines the impact of shadowing. To account for shading patterns in energy flux parameterization schemes Loridan and Grimmond (2012) propose an 'active' built 636 637 index. The latent heat flux also depends on these, and other, factors. To further develop 638 understanding of such processes and interactions it will be necessary to focus more attention on the interdependencies between energy fluxes and how these are affected by surface conditions in urban 639 640 areas.

Finally, although the Bowen ratio correction has not been applied to the data here, the biggest impact of the correction would be at small  $\beta$ . For  $\beta = 0.5$ , scintillometric  $Q_H$  is overestimated by 9% which would result in mean  $Q_H/Q^*$  being overestimated by 0.04. Implementing the  $\beta$  correction would act to further decrease  $Q_H/Q^*$  with vegetation fraction. As  $\beta$  itself has been shown to depend on the vegetation fraction, smaller  $\beta$  at larger vegetation fraction again acts to amplify rather than reduce the trend.

## 647 **4. Conclusions**

This work demonstrates the applicability of large aperture scintillometry for making spatially integrated observations over urban areas. With selection of a suitable path, adequately sited auxiliary meteorological measurements and knowledge of the land surface, sensible heat flux estimates are obtained that are representative of several neighbourhoods or across the settlement. Whilst EC measurements are representative of the local-scale (0.5 km<sup>2</sup>), the scintillometer data in this study have much larger source areas: 3.0 and 7.5 km<sup>2</sup> (95% contribution) for the LAS and BLS, respectively. 655 Remarkable temporal agreement is observed across the three different areal extents for both 656 short-term variability (e.g. the response to radiation patterns over a few hours to days) and seasonal trends. Differences in magnitudes of the fluxes between sites are attributed primarily to the role of 657 658 vegetation and reveal the influence of anthropogenic materials on surface-atmosphere interactions. Empirical relations between land cover and fluxes often underpin urban energy models and are 659 660 valuable for gauging the likely partitioning of energy, and hence the environmental conditions 661 (including thermal comfort and moisture availability), in cities where measurements have not been 662 made.

663 Comparison of the EC dataset with large-area fluxes at the city-scale provides some context to 664 the results and confirms that the EC site selection was appropriate. The scintillometer fluxes tend to 665 be smoother as a result of the greater spatial averaging. The large-scale flux measurements are also 666 much less sensitive to source area variability, for example due to changing wind direction over 667 heterogeneous surfaces. As they encompass a larger proportion of the study area, these large-area 668 fluxes are more representative and suffer less from sampling bias, whereas EC measurements are 669 easily influenced by spatially variable land cover or surface characteristics around the mast. The effect can be decreased by measuring at a greater height, but in general the land cover must be 670 671 carefully examined for each wind sector before drawing conclusions on the representativeness of 672 data from a single EC site.

673 For many purposes we are interested in fluxes at large scales, whether the application is input 674 data for, or evaluation of, land-surface models or numerical weather prediction, assessment of 675 satellite remote sensing products or representative observational datasets to characterize a 676 particular environment. Scintillometry offers a promising way forward, but there are still limitations. 677 A major source of uncertainty arises from the MOST functions. This is an area that would benefit 678 from further attention for all land cover types and has implications beyond improving the accuracy 679 of fluxes from scintillometry. Single-wavelength scintillometry may be best suited to urban areas 680 with little vegetation as the higher the Bowen ratio the smaller the error due to neglecting the  $\beta$ -681 correction (Moene 2003). Given the potential for saturation, particularly if the sensible heat flux is 682 large, it is recommended that an extra-large aperture scintillometer is considered for long paths (e.g. 683 > 4 km, for paths of similar height and fluxes of similar magnitude). Future work will likely focus on 684 the development of the scintillometry technique and the application for routine monitoring at largescales, e.g. Kleissl et al. (2009a). Such observational networks would offer valuable data for 685 686 assimilation into models that assess e.g. air quality or heat stress, both highly relevant to human 687 health and well-being.

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