

1 **Stable AMOC off state in an eddy-permitting**

2 **Coupled Climate Model**

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6 Received: date / Accepted: date

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7 **Abstract** Shifts between on and off states of the Atlantic Meridional Over-  
8 turning Circulation (AMOC) have been associated with past abrupt climate  
9 change, supported by the bistability of the AMOC found in many older, coarser  
10 resolution, ocean and climate models. However, as coupled climate models  
11 evolved in complexity a stable AMOC off state no longer seemed supported.  
12 Here we show that a current-generation, eddy-permitting climate model has  
13 an AMOC off state that remains stable for the 450-year duration of the model  
14 integration. Ocean eddies modify the overall freshwater balance, allowing for  
15 stronger northward salt transport by the AMOC compared with previous, non  
16 eddy-permitting models. As a result, the salinification of the subtropical North  
17 Atlantic, due to a southward shift of the intertropical rain belt, is counteracted  
18 by the reduced salt transport of the collapsed AMOC. The reduced salinifi-  
19 cation of the subtropical North Atlantic allows for an anomalous northward  
20 freshwater transport into the subpolar North Atlantic dominated by the gyre  
21 component. Combining the anomalous northward freshwater transport with

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22 the freshening due to reduced evaporation in this region helps stabilise the  
23 AMOC off state.

24 **Keywords** AMOC · AMOC collapse · abrupt climate change · hosing  
25 experiment · CGCM · eddy-permitting

## 26 1 Introduction

27 The Atlantic Meridional Overturning Circulation (AMOC) describes the merid-  
28 ional volume transport in the Atlantic Ocean (Wunsch, 2002). The AMOC  
29 brings warm waters to the high latitude North Atlantic, warming the climate  
30 of Northern and Western Europe. A collapse of the AMOC would lead to  
31 drastic changes in surface air temperatures over much of the Northern Hemi-  
32 sphere, in particular in the Northeast Atlantic where temperatures can drop  
33 by 9°C (Manabe and Stouffer, 1988; Vellinga and Wood, 2002; Jackson et al,  
34 2015). As a consequence of anthropogenic climate change, warming of the high  
35 latitude North Atlantic and the addition of freshwater through enhanced pre-  
36 cipitation, increased melting of sea-ice and icebergs, as well as more runoff  
37 from the Greenland ice sheet can cause the sinking branch of the AMOC to  
38 weaken and potentially shut down. Hereafter, we refer to a collapsed AMOC  
39 as an AMOC off state while, the AMOC circulation, as it is known today, is  
40 referred to as an AMOC on state.

41 Climate model projections indicate a likely weakening of the AMOC, but  
42 a complete collapse was deemed unlikely in the latest IPCC report by Collins  
43 et al (2013). However, models have difficulty correctly simulating past abrupt

44 climate changes, including an AMOC collapse, affecting the likelihood of sim-  
45 ulating future abrupt climate change (Valdes, 2011; Drijfhout et al, 2011).  
46 Paleo-proxy data have shown evidence for wide spread abrupt climate change  
47 events in the times before the Holocene from ice-core records (Dansgaard et al,  
48 1993; Blunier and Brook, 2001) and sediment cores (de Abreu et al, 2003). A  
49 possible interpretation of these events is that they are associated with switches  
50 between AMOC on and off states in the past (Broecker et al, 1990), although  
51 the spatial extent of these abrupt changes in climate can still be questioned  
52 (Wunsch, 2006). Such switches can be theoretically understood from simple  
53 box model studies showing that under the same forcing conditions it is possible  
54 to have both a stable AMOC on and off state, or only a mono-stable regime  
55 depending on the forcing (Stommel, 1961; Marotzke, 1990; Rahmstorf, 1996).  
56 The existence of bistability in these box models depends on the freshwater  
57 forcing. Similarly, some coupled climate models have found a bistable AMOC  
58 dependent on freshwater forcing when freshwater hosing was applied contin-  
59 uously (Hawkins et al, 2011; Hu et al, 2012; Sijp, 2012). However, in newer  
60 coupled climate models after applying freshwater hosing for a set amount of  
61 time the AMOC recovered after the freshwater hosing was stopped (Peltier  
62 et al, 2006; Krebs and Timmermann, 2007; Jackson, 2013) while it was pos-  
63 sible to maintain the AMOC off state in some older coupled climate models  
64 (e.g. UVic and GFDL R30 models in Stouffer et al (2006)).

65 To identify the transition between the two regimes of mono- and bistability  
66 it was proposed that the sign of the freshwater transport by the AMOC in the

67 Atlantic can be used as an indicator for its stability (referred to here as  $M_{ov}$  but  
68 often also referred to as  $F_{ov}$ ) (Rahmstorf, 1996; de Vries and Weber, 2005).  
69 When used as an indicator for AMOC stability,  $M_{ov}$  is typically measured  
70 at the southern entrance of the Atlantic near  $34^{\circ}\text{S}$ . A positive  $M_{ov}$  at  $34^{\circ}\text{S}$   
71 indicates that the AMOC imports freshwater into the Atlantic and a negative  
72  $M_{ov}$  at  $34^{\circ}\text{S}$  indicates freshwater export from the Atlantic. In an AMOC off  
73 state  $M_{ov}$  is expected to tend towards zero, thereby creating an anomalous  
74 salt import into the Atlantic for positive  $M_{ov}$  which leads to a destabilisation  
75 of the AMOC off state. On the other hand, when  $M_{ov}$  is negative an AMOC  
76 collapse will result in an anomalous freshwater import into the Atlantic helping  
77 stabilise the AMOC off state. Therefore, a positive  $M_{ov}$  can be associated with  
78 a mono-stable AMOC while a negative  $M_{ov}$  can be associated with a bistable  
79 AMOC (Huisman et al, 2010). Observational estimates of  $M_{ov}$  at the southern  
80 boundary of the Atlantic based on ship data or estimated from ARGO float  
81 data support a negative  $M_{ov}$ , suggesting that the present day AMOC resides  
82 in the bistable regime (Bryden et al, 2011; Garzoli et al, 2013). It has been  
83 recommended that the divergence of the freshwater transport into the Atlantic  
84 by the AMOC,  $\Delta M_{ov} = M_{ovS} - M_{ovN}$ , where S (N) is the southern (northern)  
85 boundary of the Atlantic is a better indicator of bistability (Huisman et al,  
86 2010; Liu and Liu, 2013).

87 When the AMOC weakens and even collapses, the reduction in northward  
88 heat transport causes a wide spread cooling of the northern hemisphere surface  
89 air temperatures (Manabe and Stouffer, 1988; Vellinga and Wood, 2002; Jack-

son et al, 2015). The cooling leads to a southward/equatorward shift of the  
latitude of maximum heating causing the dividing latitude of the northern and  
southern hemisphere Hadley circulations to shift southward (Drijfhout, 2010),  
displacing the Intertropical Convergence Zone (ITCZ). The southward shift of  
the ITCZ causes a reduction of precipitation in the subtropical North Atlantic  
region leading to a salinification of the ocean. The saltier waters in this region  
can be transported into the high latitude regions of the North Atlantic through  
large-scale instabilities kick starting the convection (e.g. the large-scale eddy  
generated in GFDL CM2.1 in Yin and Stouffer (2007)). Therefore, in order for  
the AMOC off state to remain stable this salinification needs to be balanced  
by an equally large freshening term, due to changes in ocean circulation.

In the GFDL R30 model the freshening associated with ocean circulation  
changes is large enough to counteract the salinification due to the southward  
ITCZ shift because the overturning circulation reverses (Yin and Stouffer,  
2007). In that case Antarctic Intermediate Water (AAIW) sinks to a depth of  
1000 m just south of South America and is transported northward, then up-  
wells in the North Atlantic subtropical gyre. This circulation has been named  
the reverse thermohaline circulation (RTHC). However, the RTHC only devel-  
ops in coarse-resolution ocean models and often is deeper than just the upper  
1000 m (Dijkstra, 2007; Hawkins et al, 2011; Sijp, 2012). In a newer generation  
of coupled climate models the RTHC cell does not develop (e.g. GFDL CM2.1  
in Yin and Stouffer (2007)) and without the additional freshwater transport of  
the RTHC the subtropical gyre becomes so salty that a fresh subpolar ocean

113 without deep sinking is no longer stable and the AMOC recovers (Yin and  
114 Stouffer, 2007; Jackson, 2013). The reason for the RTHC not to develop is  
115 that stronger atmospheric feedbacks promote saltier and colder thermocline  
116 water in the subtropical North Atlantic, reducing the north-south pressure  
117 gradient between the subtropical North Atlantic and subpolar South Atlantic  
118 that is driving the RTHC (Yin and Stouffer, 2007)).

119 In the very latest coupled climate models ocean eddies and swifter bound-  
120 ary currents are allowed for, changing the salt balance in the Atlantic. Ocean  
121 eddies freshen the subtropical gyre by exchanging water with the tropics and  
122 subpolar gyre (Tréguier et al, 2012). As a result, eddy-permitting and eddy-  
123 resolving models must feature a larger mean flow salt transport divergence  
124 into the subtropical gyre to maintain equilibrium counteracting freshening by  
125 the eddies. The larger mean flow salt transport divergence could allow for a  
126 stronger advective salt feedback associated with an AMOC collapse without  
127 the need of developing an RTHC. Indeed, using a higher resolution coupled  
128 climate model Spence et al (2013) achieved a stronger drop and slower recov-  
129 ery of the AMOC in a high-resolution model relative to a coarser resolution  
130 model in a relatively weak and short freshwater hosing experiment. Similarly,  
131 Weijer et al (2012), using an ocean only model, were able to show that the  
132 drop in AMOC in response to a freshwater hosing was stronger in the higher  
133 resolution model. Both studies suggest that the AMOC off state in higher  
134 resolution models could become stable. Here we discuss whether a larger salt  
135 transport by the AMOC into the North Atlantic subtropical gyre, which is

136 typical for higher resolution ocean models, can sustain a stable off state, even  
137 if the RTHC does not develop, using a 450 year long hosing experiment in an  
138 eddy-permitting coupled climate model.

## 139 **2 Model Configuration and Experiment Setup**

### 140 **2.1 Model Configuration**

141 For this study the Global Climate version 2 (GC2) (Williams et al, 2015) con-  
142 figuration of Hadley Centre Global Environmental Model version 3 (HadGEM3)  
143 (Hewitt et al, 2011) is used. This coupled climate model consists of an ocean,  
144 atmosphere, sea-ice and land-surface model coupled together with data ex-  
145 changing between the atmosphere and ocean components every 3 hours. The  
146 ocean model component of GC2, HadGEM3 uses the Global Ocean version  
147 5 (GO5) (Megann et al, 2013) of the ORCA025 configuration of the Nucleus  
148 for European Modelling of the Ocean (NEMO) (Madec, 2008) version 3.4.  
149 The ORCA025 grid uses a tri-polar structure with poles over Antarctica,  
150 Siberia and Canada and has a horizontal resolution of  $0.25^\circ$ , with the res-  
151 olution decreasing when moving towards the poles so that the grid remains  
152 quasi-isotropic. The ocean model contains 75 vertical levels with thicknesses  
153 ranging from 1 m at the surface and increasing with depth up to 200 m in  
154 the bottom layer. The sea-ice model is the global sea ice version 6 (GSI6)  
155 configuration of the Los Alamos National Laboratory sea ice model (CICE)  
156 version 3.4 (Rae et al, 2015) and is used at the same model grid as the ocean

157 model. The Global Atmosphere version 6 (GA6.0) of the Met Office unified  
158 model is used with a horizontal resolution of N216, which has a resolution of  
159 about 60 km in mid-latitudes, and has 85 levels in the vertical leading to an  
160 improved resolution in the stratosphere. Global Land version 6 (GL6) config-  
161 uration of the land model Joint UK Land Environment Simulator (JULES) is  
162 also used in this model setup but none of its data is analysed in this study.  
163 Heat, freshwater and momentum fluxes are passed between the atmosphere  
164 and ocean/ice model every three hours through the OASIS coupler while the  
165 ocean and sea-ice model exchange fluxes every ocean model time step (22.5  
166 min) without the use of flux adjustment. The eddy permitting resolution of  
167 the ocean model has lead to a reduction in the North Atlantic cold bias and  
168 the atmospheric model shows improved Atlantic and European blocking events  
169 (Scaife et al, 2011) and the ability to better predict the winter North Atlantic  
170 Oscillation (Scaife et al, 2014), in previous versions of the HadGEM3 model  
171 setup, i.e. GloSea5.

## 172 2.2 Experiment Setup

173 In this study two experiments from the GC2 model are considered, a 150-year  
174 long present day control simulation and a 450 year long hosing experiment.  
175 The hosing experiment is a continuation of the experiment analysed in Jackson  
176 et al (2015) (See reference for more details). The present day control simula-  
177 tion was started from a 36 year long development run of HadGEM3, which  
178 was initialised with EN3 data (Ingleby and Huddleston, 2007) averaged over

179 2004-2008 and the hosing experiment is started from year 42 of the control  
180 experiment. The control simulation uses CO<sub>2</sub> concentrations based on 1978  
181 levels and held constant throughout both simulations. The main goal of the  
182 hosing experiment was to collapse the AMOC, therefore, the methodology is  
183 based on Vellinga and Wood (2002), which allows for a rapid collapse of the  
184 AMOC but is very idealised. For the first 10 years of the hosing experiment  
185 the salinity in the model is perturbed by an amount equivalent to a hosing  
186 of 10 Sv, making a total of 100 Sv-years additional freshwater. This is done  
187 through reducing the salinity in the Atlantic Ocean north of 20°N and in the  
188 Arctic by 0.64 psu in the upper 350 m and then tapering to zero over the  
189 next 186 m (Fig. 1). This is done instantaneously every December 1 and, as is  
190 common practice in hosing experiments, is compensated by adding 0.008 psu  
191 everywhere else in the ocean allowing for the total salinity to be conserved  
192 (Fig. 1). After the 10 years of hosing is completed and the model is allowed to  
193 continue without changes for another 440 years.

### 194 **3 Results**

195 In the 450 year long hosing experiment the AMOC is able to collapse and  
196 remain very weak for the entire duration of the model integration (Fig. 2).  
197 During and after the 10 year hosing period the ocean begins to adjust, with  
198 salinity anomalies slowly spreading southward from the hosing region towards  
199 the equator and also spreading downward in the water column. Since we want  
200 to discuss the evolution of the ocean fields in 100 year time-slices, we will take

201 the period 311-410 (301 to 400 years after the hosing stopped) as representa-  
202 tive for the final state of the model. The mean salinity is 0.86 psu fresher in the  
203 hosing region towards the end of the hosing simulation (years 311-410) relative  
204 to the control simulation. The sea surface salinity (SSS) anomaly with respect  
205 to the control run features a comma shaped pattern in the North Atlantic sub-  
206 tropical gyre (Fig. 3a), as typical with most fresh water hosing experiments  
207 (Krebs and Timmermann, 2007; Yin and Stouffer, 2007). The sea surface tem-  
208 peratures (SSTs) also drop due to the reduction of northward heat transport  
209 from the AMOC off state (Fig. 3b and Jackson et al (2015)). The decrease in  
210 SSTs allow for the seasonal sea-ice to extend further southward reaching as  
211 far south in winter as the Grand Banks, as well as covering a large portion of  
212 the Norwegian and Baltic Seas (Fig. 3b). The reductions in SSS and SST fall  
213 within the range of what has been seen in previous modelling studies with a  
214 similar magnitude of freshwater hosing (Yin and Stouffer, 2007).

### 215 3.1 AMOC Streamfunction

216 The control simulation features a realistic AMOC with a maximum strength  
217 of 17.4 Sv at 27°N and at a depth of 773 m in the mean (Fig. 4a). The depth  
218 reached by the North Atlantic Deep Water cell is slightly shallower than that in  
219 observations (3000 m as opposed to 4000 m in (Kanzow et al, 2010; Smeed et al,  
220 2014)), a common problem in ocean models (Danabasoglu et al, 2014). The  
221 Faroe Bank Channel overflow (defined as waters denser than  $\sigma_\theta = 27.8 \text{ kg/m}^3$ )  
222 is slightly weaker in this model than in observations (1.8 Sv as opposed to 1.9

223 Sv (Hansen and Østerhus, 2007)). This overflow is mainly missing the weak  
224 cold waters below 0 °C which account for the majority of the overflow waters  
225 in the observations, making the model overflow less dense. For the Denmark  
226 Strait the overflows are considerably weaker when considering waters denser  
227 than  $\sigma_\theta = 27.8 \text{ kg/m}^3$  (1.4 Sv as opposed to 3.4 Sv (Jochumsen et al, 2012)),  
228 which again is missing the very cold water masses. However, for the Denmark  
229 Strait choosing the density cut off to be  $\sigma_\theta = 27.8 \text{ kg/m}^3$  misses a lot of the  
230 overflow waters. By choosing the density class cut off of to be  $\sigma_\theta = 27.6 \text{ kg/m}^3$ ,  
231 matching the depth of density cutoff in Jochumsen et al (2012), the overflow  
232 increases to 2.9 Sv. These differences in the overflows between the model and  
233 observations could potentially lead to the shallower North Atlantic Deep Water  
234 cell. The main convection sites are in the Labrador Sea, Greenland Sea and  
235 South of Iceland (Fig. 4b) as expected from observations (de Boyer Montégut  
236 et al, 2004). However, the too buoyant overflows could potentially account  
237 for the slightly weaker and shallower AMOC as compared to observations at  
238 26.5°N (Fig. 2a, 15.7 Sv as opposed to 17.5 Sv (Smeed et al, 2014)) but this  
239 is not investigated in more detail.

240 Based on an AMOC index at 26.5°N and between 500-2000 m the AMOC  
241 collapses very rapidly during the hosing, leading to a minimum in AMOC at  
242 year 4 (Fig. 2a). After the hosing has stopped the AMOC recovers slightly,  
243 achieving a maximum at year 21, before dropping in strength again and re-  
244 maining in a very weak state for the duration of the model integration. How-  
245 ever, there is a noticeable weak trend in the AMOC index at 26.5°N which by

246 the end of the model integration causes the AMOC to increase in strength to  
247 just over 5 Sv (Fig. 2a). This increase in AMOC strength is slow and occurs  
248 later in the model integration than seen in previous climate model studies  
249 (Vellinga and Wood, 2002; Stouffer et al, 2006; Jackson, 2013). Also, it only  
250 applies to a shallow, wind-driven, AMOC that does not extend further north  
251 than the subtropics. Considering an AMOC index further to the north (maxi-  
252 mum between 50°N - 65°N and 500-2000 m depth) the AMOC collapse shows  
253 no hint of recovering (Fig. 2b). There is no sign of increasing mixed layer depth  
254 in the subpolar North Atlantic due to the onset of deep convection (Fig. 4d).  
255 Both subtropical and subpolar wind-driven cells are enhanced near the sur-  
256 face related to the positive North Atlantic Oscillation (NAO) that develops  
257 in response to the AMOC collapse (Jackson et al, 2015). The AMOC stream-  
258 function does not develop a stable RTHC after the AMOC collapses. Despite  
259 this, the AMOC off state appears stable, at least for 450 years. In year 311-410  
260 there appears to be no convection present in the high latitude regions (Fig. 4d)  
261 and similarly the overflows in the Denmark Strait and Faroe Banks Channel  
262 have completely collapsed to 0 with no signs of recovery.

### 263 3.2 Atmospheric Response

264 The southward shift of the ITCZ is reflected in the net precipitation (pre-  
265 cipitation - evaporation + runoff, PER) and causes a reduction in the surface  
266 freshwater flux into the ocean just north of the equator and an increase south of  
267 the equator (Fig. 5a,b). These changes in PER reduce the amount of freshwater

268 added to the subtropical North Atlantic with the majority of the reduction in  
269 precipitation occurring in the subtropical North Atlantic which loses 0.047 Sv  
270 in years 311-410 (Table 1, Fig. 5c). This reduction in PER is an atmospheric  
271 feedback to the AMOC collapse that acts to destabilise the AMOC off state  
272 by salinifying the North Atlantic.

273 Over the subpolar North Atlantic evaporation is reduced due to the increase  
274 in sea-ice cover blocking latent heat exchange and the decrease in atmospheric  
275 temperatures reducing the amount of atmospheric water vapour content (Ta-  
276 ble 1, Fig. 5b) (Drijfhout, 2014). Despite the reduction in evaporation being  
277 small relative to the precipitation changes in the subtropical regions, it is large  
278 enough to outweigh the reduction in precipitation over the subpolar Atlantic.  
279 The subsequent increase in PER causes an anomalous freshening of the sinking  
280 regions (Fig. 5d) with a magnitude of 0.042 Sv in the years 311-410 (Table 1).  
281 The rate at which the precipitation and evaporation anomalies change reduces  
282 as the model integration continues, especially for the evaporation. This subpo-  
283 lar freshening is an atmospheric feedback that stabilises the AMOC off state  
284 through freshening the North Atlantic. The salinification over the subtropical  
285 North Atlantic is marginally stronger than the freshening over the subpolar  
286 North Atlantic (Table 1). The salinification of the subtropical North Atlantic  
287 could eventually lead to more saline waters being transported in the subpolar  
288 North Atlantic as seen in the GFDL CM2.1 model (Yin and Stouffer, 2007).  
289 Nevertheless, the off state remains stable here, while it quickly destabilises  
290 in the GFDL CM2.1 model. It should be noted that the initial atmospheric

291 response of precipitation and evaporation in HadGEM3 is similar to that in  
 292 GFDL CM2.1. However, in the GFDL CM2.1 model the precipitation anoma-  
 293 lies associated with a southward shift of the ITCZ are not maintained as the  
 294 AMOC recovers, while here the anomaly continues to show the characteristic  
 295 dipole pattern over the equator although the amplitude is slowly decreasing  
 296 (Fig. 5c). This brings up the question which additional feedbacks are present  
 297 in HadGEM3, stabilising the AMOC off state? To answer this question we  
 298 analyse in detail the freshwater budget in the subtropical and subpolar North  
 299 Atlantic.

### 300 3.3 Freshwater Budget

301 The freshwater budget analysis is based on an extension to the calculations  
 302 detailed in Drijfhout et al (2011) (see appendix for details). The freshwater  
 303 budget can be summarised as follows:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + PER + M_{mix}, \quad (1)$$

304 where  $M_{trend}$  is the freshwater trend in the region of interest,  $\Delta M_{ov/az/eddy}$   
 305 represents the divergence of the freshwater transport for the specific region,  
 306 in our case the southern boundary minus the northern boundary, for the vari-  
 307 ous components of the transport,  $PER$  is the precipitation minus evaporation  
 308 plus runoffs over the specific region of interest and finally  $M_{mix}$  is the residual  
 309 term of the budget, mainly comprised of mixing along the boundaries. In eqn.  
 310 1 the decomposition of mean flow transport divergence into an overturning

( $M_{ov}$ ) and gyre ( $M_{az}$ ) component was motivated by the much stronger coupling between  $M_{ov}$  and AMOC than between  $M_{az}$  and AMOC at the southern boundary of the Atlantic, when they budget is applied to the Atlantic as a whole. Especially in the North Atlantic subpolar gyre this decomposition can be questioned. However, this framework can still be used to link area integrated changes in freshwater budget to changes in the AMOC, especially in the North Atlantic subtropics. It appears that changes in  $M_{ov}$  are first order in eqn. 1 and can be understood from the AMOC collapse as they are dominated from changes in the zonal mean velocity field. In addition it allows for comparison with observations where freshwater transports have been diagnosed using the same framework (McDonagh et al, 2010; Bryden et al, 2011; Garzoli et al, 2013; McDonagh et al, 2015-in press). When the model is in an equilibrium state the changes in PER are approximately balanced by changes in freshwater transport by overturning circulation ( $M_{ov}$ ), azonal circulation ( $M_{az}$ ) and eddies ( $M_{eddy}$ ). We apply the freshwater budget analysis to the subtropical North Atlantic, defined as  $10^{\circ}\text{N}$  to  $45^{\circ}\text{N}$ , and to the subpolar North Atlantic, defined to be  $45^{\circ}\text{N}$  to  $70^{\circ}\text{N}$ . These boundaries were chosen to coincide with the boundaries of the subtropical and subpolar gyres, with the subpolar gyre region containing the main sinking regions of the North Atlantic. The region specific freshwater budget analyses are summarised in Table 1 and graphically in Fig. 6. The atmospheric contributions to the freshwater budget have already been discussed; below we discuss the freshwater transport terms as well as the freshwater budget as a whole.

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### 334 3.3.1 Freshwater Transports

335 The AMOC off state is associated with changes in the freshwater transport  
336 terms that must be able to balance the changes in PER, especially in the  
337 subtropical North Atlantic, to prevent a salinification of the North Atlantic  
338 and hence a return to the AMOC on state. In the control simulation the  
339 freshwater transport due to the overturning,  $M_{ov}$ , is negative throughout the  
340 entire Atlantic Ocean, indicating that the AMOC is transporting freshwater  
341 southward/salt northward (Fig. 7). The negative  $M_{ov}$  at 34°S is consistent with  
342 observations (Bryden et al (2011), Garzoli et al (2013), McDonagh Personal  
343 Communications based on McDonagh and King (2005)), despite being slightly  
344 weaker, and is a possible indication for a bistable AMOC (Fig. 7). After the  
345 AMOC collapses the magnitude of  $M_{ov}$ , as expected, decreases and over time  
346 adjusts to a new equilibrium (Fig. 8a). The reduction in magnitude of  $M_{ov}$  can  
347 be attributed to the reduction in AMOC transport, with changes in salinity  
348 only having a small effect (Fig. 8b). These changes lead to an anomalous  
349 northward transport of freshwater south of 45°N in the Atlantic Ocean (Fig.  
350 8b). Even more important, however, is the sign of the divergence of  $M_{ov}$  instead  
351 of the sign of  $M_{ov}$  itself, since it is the divergence that determines whether  
352 or not there will be a freshening or salinification in the region of interest.  
353 The subtropical North Atlantic has an increase of 0.132 Sv of freshwater due  
354 to the changes in the divergence of  $M_{ov}$  (Fig. 6c, Table 1). The associated  
355 increase in freshwater is twice the amount of freshwater required to balance  
356 the anomalous salinification caused by changes in PER (Table 1). Changes in

357 the divergence of  $M_{az}$  and  $M_{eddy}$  need to enhance the salinification caused by  
358 PER and thereby balance the changes in divergence of  $M_{ov}$ .

359 The salinity decrease after the AMOC collapse is largest at the eastern side  
360 of the basin, which does not only hold for the surface (Fig. 3a) but, also at  
361 depth (not shown). This decrease in salinity is strongest over the southward  
362 branch of the subtropical gyre and northward branch of the subpolar gyre,  
363 near the eastern boundary, leading to changes in  $M_{az}$ . This results in a de-  
364 crease in  $M_{az}$  in the subtropical gyre and an increase in  $M_{az}$  in the subpolar  
365 gyre, while changes at the gyre boundaries are small (e.g 10°N and 45°N)  
366 (Fig. 8c,d). Relative to the climate models in Yin and Stouffer (2007), which  
367 have a coarser resolution, HadGEM3 has larger amplitude in  $M_{az}$  divergence,  
368 also leading to larger changes in its divergence after the collapse. This is due  
369 to the increase in model resolution leading to stronger gyres (Tréguier et al,  
370 2005; Spence et al, 2013) and less east-west difference in salinity bias (Yin and  
371 Stouffer (2007) their Fig. 1), likely due to the Gulf Stream separation being  
372 too far north in lower resolution models. The change in divergence of  $M_{az}$  for  
373 both the subtropical and subpolar North Atlantic reduces the amount of fresh-  
374 water being transported into these regions (Fig. 6, Table 1). This anomaly in  
375 freshwater transport partially balances the additional freshwater being added  
376 to the subtropical North Atlantic by changes in  $M_{ov}$  (Table 1, Fig. 6c) and  
377 changes in the subpolar gyre PER and mixing (Fig. 6b).

378 The resolution of the model used in this study allows for the analysis of the  
379 effect of eddies on the freshwater budget from the equator to mid latitudes.

380 Here, freshwater transport due to eddies is defined as the difference between  
381 total freshwater transport and freshwater transport calculated by using the  
382 seasonal fields only (see appendix for more details). The main effect of the  
383 eddies is to exchange water between the subtropical and subpolar gyres, fresh-  
384 ening the former and salinifying the latter (Fig. 8e and Table 1). Immediately  
385 after the AMOC collapse the salinity gradient at the edge of the hosing re-  
386 gion becomes very large leading to a large increase in the southward freshwater  
387 transport by the eddies at 20°N. Within a few decades after the freshwater hos-  
388 ing  $M_{eddy}$  becomes relatively small again compared to  $M_{ov}$  and  $M_{az}$  with val-  
389 ues similar to the control integration (Fig. 8e). In the eddy-permitting model  
390 the freshening of the subtropical North Atlantic by  $M_{eddy}$  and the increased  
391 freshening by a larger  $M_{az}$  play a similar role to the flux adjustment in coarser  
392 resolution climate models in the control integration (e.g. GFDL R30 model in  
393 Yin and Stouffer (2007)). This helps to stabilise the freshwater budget by al-  
394 lowing for a larger negative  $M_{ov}$  in the control integration and subsequently  
395 a larger change in  $M_{ov}$  after the AMOC collapses. The change in  $M_{ov}$  is now  
396 large enough to balance all other terms in the freshwater budget without leav-  
397 ing a strong positive salinity trend in the subtropical North Atlantic. As model  
398 resolution increases further towards eddy-resolving the magnitude of  $M_{eddy}$  is  
399 expected to become even larger (Tréguier et al, 2012), further adding to the  
400 stabilising effect of the eddies.

### 401 3.3.2 Total Freshwater Budget

402 The total freshwater trend in the subtropical North Atlantic still shows a small  
403 salinification over the 311-410 year period, slightly stronger than the salini-  
404 fication in the control run (Table 1, Fig. 6). Despite the salinification of the  
405 subtropical North Atlantic, the subpolar North Atlantic shows a freshening  
406 trend, enhancing the salinity gradient between the two (Table 1, Fig. 6). The  
407 anomalous freshening trend of the subpolar North Atlantic can be attributed  
408 to the combination of decreased evaporation in this region, the anomalous  
409 northward freshwater transport at the gyre boundaries and an increased mix-  
410 ing term (i.e  $M_{mix}$ ). The gradient in salinity across the North Atlantic, despite  
411 being stronger than in the control integration, does not lead to large-scale in-  
412 stabilities that suddenly give rise to very strong salinity transports as seen  
413 in Yin and Stouffer (2007). The eddies are likely helping to keep the gradient  
414 small enough to avoid a sudden large-scale instability to develop and to restart  
415 the convection in the high latitude sinking regions.

## 416 4 Discussion

417 The AMOC response to freshwater perturbations has been previously investi-  
418 gated in a large CMIP/PMIP coordinated experiment (Stouffer et al, 2006).  
419 A freshwater hosing of 0.1 Sv and 1 Sv was applied for 100 years, versus a  
420 hosing of 10 Sv over 10 years in the present experiment. Of the nine models  
421 involved in the 100 Sv-year hosing experiment, seven models had started the

422 transition from the off state back to the on state before 100 years after the  
423 completion of the hosing. Two models remained in the off state; one model of  
424 intermediate complexity, Uvic, and one older GFDL model, GFDL R30. The  
425 different behaviour between GFDL R30 and a newer version, GFDL CM2.1,  
426 was afterwards analysed (Yin and Stouffer, 2007) and it was argued that the  
427 stable off state in GFDL R30 was maintained by flux adjustment and weak  
428 atmospheric feedbacks allowing the RTHC to develop. This result led to the  
429 paradigm that newer generation climate models that no longer use flux adjust-  
430 ment and feature more realistic atmospheric dynamics are not able to maintain  
431 a stable AMOC off state (Yin and Stouffer, 2007; Liu et al, 2014). Here we  
432 show that an eddy-permitting coupled climate model is able to maintain a  
433 stable AMOC off state for 440 years after the hosing is completed, which is  
434 more than twice as long as the runs performed in the CMIP/PMIP experiment.  
435 The increase in freshwater transport into the subtropical North Atlantic due to  
436 higher-resolution eddies and increased boundary currents allow the AMOC to  
437 transport more salt northwards across the entire Atlantic basin. This stronger  
438 advective salt feedback is key for the model to be able to counteract the strong  
439 atmospheric response over the tropical/subtropical North Atlantic basin that  
440 features in complex climate models when the AMOC collapses. In a sense, ed-  
441 dies and swifter boundary currents play a similar role in the freshwater budget  
442 to the flux adjustment used in older generation climate models.

443 Some coupled climate models of lower complexity have been integrated for  
444 even longer durations with some of them having the AMOC off state become

445 unstable after many centuries (Krebs and Timmermann, 2007). We cannot ex-  
446 clude that such a transition will eventually occur in HadGEM3, but at present  
447 there is no deep water formation site returning to the high latitude North  
448 Atlantic (Fig. 4d) and the freshwater budget shows no signs of a potential  
449 recovery. While the subtropical North Atlantic is continuing to increase its  
450 salinity, albeit with a very small trend, the subpolar North Atlantic is getting  
451 relatively fresher, hampering the restart of deep convection. Also when taking  
452 the subpolar North Atlantic and the Arctic into account there is an overall  
453 freshening trend suggesting that having a return of deep convection in the high  
454 latitude North Atlantic in the near future is very unlikely.

455 When taking the salinity of the entire Atlantic into account, as was done  
456 in Sijp (2012), we do not see a difference in salinity between the hosing and  
457 control simulations. In Sijp (2012) the two states in Atlantic mean salinity are  
458 associated with the AMOC on and off states. However in Sijp (2012) an RTHC  
459 develops, which is responsible for the low salinity state, while in HadGEM3  
460 the AMOC off state still has a shallow wind-driven cell that extends into the  
461 Northern Hemisphere, preventing a low salinity state. However if we focus  
462 on the region north of 35°N only, the hosing integration is 0.7 psu fresher  
463 in the upper 3000 m than the control integration, indicating that low and  
464 high salinity states in the subpolar gyre can be associated with the AMOC on  
465 and off state in this model. This suggests that a bifurcation in basin average  
466 salinity no longer exists in HadGEM3 but bistability in subpolar gyre salinity  
467 is still existent.

468 The increase in northward salt transport by the AMOC in HadGEM3,  
469 relative to the coarser resolution climate models (Yin and Stouffer, 2007) is  
470 associated with a reduction in vertical gradient of salinity bias in the Atlantic.  
471 The model using flux adjustment in Yin and Stouffer (2007), GFDL R30,  
472 showed little bias, but the climate model that did not use flux adjustment,  
473 GFDL CM2.1, featured larger biases. In particular, the salinity bias in the  
474 GFDL CM2.1 model contained a pronounced vertical gradient with a negative  
475 salinity bias near the surface and a positive bias at deeper levels throughout  
476 most of the Atlantic. Combined with an AMOC that transports surface water  
477 northward and deep water southward this salinity bias leads to  $M_{ov}$  being  
478 strongly biased towards positive values. With a positive  $M_{ov}$ , when the AMOC  
479 collapses, more saline water will be transported into the Atlantic, aiding the  
480 recovery of the AMOC, as is clearly the case with GFDL CM2.1 in Yin and  
481 Stouffer (2007). These results are supported by the analysis of Liu et al (2014),  
482 where they see a larger negative salinity bias in the surface for the un-flux  
483 adjusted models relative to flux adjusted models. This led to a less negative  
484  $M_{ov}$  at 34°S, reducing the likelihood of bistability. For the model used in this  
485 study, HadGEM3, the salinity bias has a weak negative vertical gradient in the  
486 Southern Atlantic in the depths corresponding with the North Atlantic Deep  
487 Water (NADW) cell of the AMOC and a mostly positive bias in the upper  
488 1000 m throughout the rest of the Atlantic (Fig. 9). This weaker salinity bias  
489 is likely due to the fact that the model is eddy permitting and has swifter  
490 and narrower boundary currents. In GFDL CM2.1 the positive salinity bias

491 peaks near  $20^\circ\text{N}$  (Yin and Stouffer, 2007), while in HadGEM3 the model bias  
492 is smaller there (Fig. 9) since  $20^\circ\text{N}$  coincides with a convergence in freshwater  
493 transport due to the eddies (Fig. 8e). The vertical structure of the salinity  
494 bias in HadGEM3 is too small to affect the sign of  $M_{ov}$ : it only has a minor  
495 effect on  $M_{ov}$  south of the equator and an even weaker effect between the  
496 equator and  $30^\circ\text{N}$  (Fig. 7). However, a further reduction of the salinity bias  
497 would move the model values of  $M_{ov}$  even closer to the estimates based on  
498 observations of  $M_{ov}$  throughout the Atlantic (Fig. 7).

499 At  $26^\circ\text{N}$   $M_{ov}$  is  $-0.601$  Sv in the control integration of HadGEM3 (about  
500  $-0.6$  Sv GFDL CM2.1 Yin and Stouffer (2007)) and  $-0.78$  Sv in observations  
501 (McDonagh et al, 2015-in press). A larger difference between HadGEM3 and  
502 the models analysed in Yin and Stouffer (2007) occurs at the southern bound-  
503 ary of the subtropical gyre ( $10^\circ\text{N}$ ). In HadGEM3  $M_{ov}$  is largely negative at  
504 those latitudes,  $-0.361$  Sv, while in GFDL CM2.1  $M_{ov}$  has about half the  
505 amplitude, approximately  $-0.2$  Sv. Both models agree on  $M_{ov}$  being slightly  
506 negative at the subtropical-subpolar boundary, around  $-0.2$  Sv. Thus the dif-  
507 ferent values at the southern boundary of the subtropical gyre in the models  
508 determines the sign of the divergence of  $M_{ov}$  over the subtropical gyre and  
509 the sign of the advective salt feedback in this area when the AMOC weak-  
510 ens or collapses. Unfortunately there are no estimates of  $M_{ov}$  near  $10^\circ\text{N}$ , but  
511 the reduced salinity bias in HadGEM3 suggests that a negative  $M_{ov}$  at those  
512 latitudes is the more likely.

513 Of some concern is the absence of an RTHC in the AMOC streamfunction  
514 after hosing is applied. Stability analysis of coarse-resolution ocean-only mod-  
515 els suggests that the collapsed AMOC is an unstable steady state, dividing  
516 the attractor space between a stable on state and a stable RTHC reaching  
517 to the bottom of the Atlantic (Dijkstra, 2007). Furthermore, the studies of  
518 Saenko et al (2003) and Sijp et al (2012) point out that it is the density dif-  
519 ference between the NADW and the Antarctic Intermediate Water (AAIW)  
520 formation regions which are important for the existence of an RTHC. In this  
521 study the density of the NADW formation region is not reduced enough after  
522 the initial hosing to become lighter than the water in the AAIW formation  
523 region as RTHC is not maintained. This study and the results of Yin and  
524 Stouffer (2007) suggest that the development of the RTHC is suppressed by  
525 atmospheric feedbacks. However, there is at present insufficient analysis to  
526 conclude whether atmospheric feedbacks really prevent a stable RTHC to de-  
527 velop, or whether there are other reasons for why it is absent in HadGEM3.  
528 For HadGEM3, we believe there are two possibilities; 1) the AMOC off state,  
529 despite the maintaining an AMOC off state for much longer than the models  
530 used in the PMIP experiment of Stouffer et al (2006), will eventually return to  
531 an AMOC on state, or 2) the AMOC off state is a stable solution of coupled  
532 climate models at eddy-permitting or higher resolution.

533 In HadGEM3 the presence of eddies and swifter boundary currents (stronger  
534 gyres) allows for stronger northward salt transport of the AMOC, stabilising  
535 the off state (Fig. 8). An even higher-resolution (1/12 degree), eddy-resolving

536 ocean model features even larger northward salt transport by  $M_{eddy}$  than the  
537 eddy-permitting version (Tréguier et al, 2012), implying an AMOC off state  
538 could potentially be favoured by even stronger advective salt feedbacks. On  
539 the other hand, the latitudinal structure of  $M_{ov}$  in HadGEM3 seems broadly  
540 consistent with the few estimates we have at different latitudes (Fig. 7) and  
541 we anticipate only a small improvement in this respect when going to higher  
542 resolution in the ocean component of climate models.

## 543 **5 Conclusions**

544 The goal of the model run analysed in this study was to rapidly collapse the  
545 AMOC and study the stability of the AMOC off state. Several other studies  
546 have been done choosing a freshwater hosing setup that more realistically rep-  
547 resents what could happen in the climate system (Weijer et al, 2012; Spence  
548 et al, 2013; Swingedouw et al, 2013). These studies have all shown that it is  
549 possible to weaken the AMOC using a more realistic hosing setup. On top of  
550 that Weijer et al (2012) and Spence et al (2013) have shown that when using  
551 higher resolution the amount by which the AMOC weakens is larger relative  
552 to their coarse resolution models used in those studies. However, these studies  
553 often only have been run for 50 years in the high resolution setting. These re-  
554 sults plus the results presented in this study support the possibility of coupled  
555 models being more likely to model abrupt climate changes as model resolutions  
556 continue to improve. At higher resolution a stronger advective salt feedback  
557 associated with the AMOC, leading to a freshening of the subtropical North

558 Atlantic, overcomes the damping feedback that salinifies this region, associated  
 559 with the atmospheric response to an AMOC collapse. This changed balance  
 560 between the different feedbacks makes the transition to a stable AMOC off  
 561 state possible, when the freshwater transports at high latitudes in the North  
 562 Atlantic increases. This is illustrated by the eddy-permitting climate model,  
 563 HadGEM3, being able to maintain an AMOC off state for 440 years.

#### 564 Appendix: Freshwater Budget Calculation

565 The freshwater budget calculation used in this study is based on the method presented  
 566 in Drijfhout et al (2011) with modifications to include the effects of a northern and southern  
 567 boundary, as well as specifics to the version of NEMO used (GO5, version 3.4 of NEMO)  
 568 (Megann et al, 2013). Mean flow transports are based on 3 month means, while total trans-  
 569 ports (i.e.  $vS$ ) are calculated online and are updated after each ocean model time step, which  
 570 are later averaged over the years of interest removing the effects of the seasonal cycle on the  
 571 budget. Following Drijfhout et al (2011), the equation for the volume budget is as follows:

$$V_t = T_S - T_N - T_{Med} + PER - Res_V, \quad (2)$$

572 where  $V_t$  is the rate of change of the volume,  $T_{(N/S)}$  are volume transports through the  
 573 northern and southern boundaries,  $T_{Med}$  is the volume transport through the Strait of  
 574 Gibraltar,  $PER$  is the precipitation minus evaporation plus runoffs and  $Res_V$  is the error  
 575 generated by the choice of differencing scheme and temporal resolution of the data. The  
 576 value of  $Res_V$  is computed as a residual to close the budget. Since the model has a free  
 577 surface  $V_t$  is equivalent to the changes in the sea surface height using backwards differencing.  
 578 The main differences between eqn. 2 and eqn. 4 in Drijfhout et al (2011) are that we have left  
 579 the choice of the northern and southern boundaries as arbitrary as opposed to choosing  $34^\circ S$   
 580 and the Bering Strait and we have included a term,  $T_{med}$  for the volume transports through  
 581 the Strait of Gibraltar. In this configuration of NEMO the transports are computed without  
 582 taking the changes in sea surface height into account. For the regions of interest used in this  
 583 study the values of  $Res_V$  are relatively small resulting in  $O(10^{-4} \text{ Sv})$  for the North Atlantic

584 subtropical gyre and  $O(10^{-5}$  Sv) for the North Atlantic subpolar gyre, which in both cases  
 585 is the smallest term in the budget with the remaining terms ranging from  $O(10^{-3}$  Sv) to  
 586  $O(1$  Sv). Choosing instantaneous values of sea surface height from the model restart files in  
 587 the computation of  $V_t$  leads to  $Res_V$  having the same order as the precision in which the  
 588 data is stored but, not all model restart files were available.

589 Similarly the salinity budget in terms of freshwater becomes the following:

$$M_{trend} - V_t = M_S - M_N - M_{Med} + M_{Mix} - Res_V + H, \quad (3)$$

590 where  $M_{trend}$  is the rate of change of freshwater in the region of interest,  $M_{(N/S)}$  are the  
 591 northward/southward freshwater transports,  $M_{med}$  is the freshwater transport through the  
 592 Strait of Gibraltar,  $H$  represents the freshwater hosing and  $M_{mix}$ , computed as a residual,  
 593 closes the budget capturing mixing and errors introduced by the temporal resolution of the  
 594 data, as well as, the choice of reference salinity,  $S_o$ . The conversion between salinity based  
 595 terms to the freshwater based terms in eqn. 3 is done through multiplying all the terms in  
 596 the equation by  $-1/S_o$ . Note that we have dropped the negative sign before  $M_{trend}$  in eqn.  
 597 3, contrary to Drijfhout et al (2011) so that positive values indicate an increase in freshwater  
 598 not salinity. In this case the hosing is included in the salinity budget and not the volume  
 599 budget since it is computed as a redistribution of salinity in this model study. Combining  
 600 eqns. 2 and 3 gives the following expression for the fresh water budget:

$$M_{trend} = (M_S + T_S) - (M_N + T_N) - (M_{Med} + T_{Med}) + M_{Mix} + PER + H. \quad (4)$$

601 The  $M_{(N/S)}$  terms can be divided into eddy and mean flow components since the ocean  
 602 model output includes  $vS$  computed at every model time step. The eddy contribution to  
 603 the freshwater transport is defined as follows:

$$M_{(eddy(N/S))} = \frac{-1}{S_o} \int_{N/S} (\overline{vS} - \overline{v}\overline{S})dA = M_{(N/S)} - M_{(mean(N/S))}, \quad (5)$$

$$\rightarrow M_{(N/S)} = M_{(mean(N/S))} + M_{(eddy(N/S))}, \quad (6)$$

604 where the integral is taken over each zonal section of the Atlantic basin,  $\overline{vS}$  is the total  
 605 seasonal mean transport,  $\bar{v}$  and  $\bar{S}$  are the seasonal mean meridional velocity and salinity and  
 606  $M_{(mean(S/N))} = -1/S_o \int_{N/S} \bar{v}\bar{S}dA$  represents the non-eddy transports, with the overbar  
 607 denoting a mean computed over 3 months. A map of the eddy kinetic energy (Fig. 10)  
 608 shows that the eddy field in HadGEM3 is very similar to other models of similar resolution  
 609 (Delworth et al, 2012), perhaps even slightly closer to what is expected from observations.  
 610 The eddy contribution is computed in a very similar way to Tréguier et al (2012), in which it  
 611 was also shown that the eddy contribution will be even stronger at higher model resolutions.  
 612 Since the current model resolution is eddy-permitting it is not possible to completely resolve  
 613 eddies at all latitudes, therefore caution must be taken in interpreting the role of the eddies  
 614 in the high latitudes. Similar to what is done in Drijfhout et al (2011),  $M_{(mean(S/N))}$  can  
 615 be divided into an overturning  $M_{(ov(S/N))}$ , azonal  $M_{(az(S/N))}$  and the volume transport  
 616  $T_{(S/N)}$  terms as follows:

$$M_{mean(N/S)} = M_{ov(N/S)} + M_{az(N/S)} - T_{(N/S)}, \quad (7)$$

$$M_{ov(N/S)} = \frac{-1}{S_o} \int_{N/S} v^* \langle S \rangle dA, \quad (8)$$

$$M_{az(N/S)} = \frac{-1}{S_o} \int_{N/S} v' S' dA, \quad (9)$$

617 where  $\langle f \rangle = \int f dx / \int x$  is the zonal mean,  $f' = f - \langle f \rangle$  is the difference from the zonal  
 618 mean,  $\hat{f} = \int f dA / \int dA$  is the zonal section mean or barotropic component and  $f^* = \langle f \rangle - \hat{f}$   
 619 is the zonal mean baroclinic component for  $f = \bar{v}$  or  $f = \bar{S}$ . Substituting eqns. 6 and 7 into  
 620 eqn. 4 gives the final form for the zonal freshwater budget equation:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + \Delta M_{Med} + M_{Mix} + PER + H, \quad (10)$$

621 where  $\Delta M_{ov} = M_{(ov(S))} - M_{(ov(N))}$ ,  $\Delta M_{az} = M_{(az(S))} - M_{(az(N))}$ ,  $\Delta M_{eddy} = M_{(eddy(S))} -$   
 622  $M_{(eddy(N))}$  and  $\Delta M_{Med} = -M_{Med} - T_{Med}$ .

623 There are several possible valid choices of the reference salinity; the mean salinity over the  
 624 entire volume of the region used in the budget calculation, the mean salinity over the section

625 used as the northern (southern) boundary or the mean salinity from the Strait of Gibraltar.  
626 For this study it was chosen to use the mean salinity at the boundary between the North  
627 Atlantic subtropical and subpolar gyres for  $S_o$ , the reference salinity. Choosing one of the  
628 other salinities as a reference salinity creates a maximum difference of  $O(10^{-4} \text{ Sv})$ , which is  
629 less than 10% of the smallest value represented in our budget analysis. To further simplify  
630 the budget analysis only times when there is no hosing being applied are considered and  
631 the freshwater transport through the Strait of Gibraltar is combined with the mixing term,  
632 resulting in the following final equation for the budget analysis:

$$M_{trend} = \Delta M_{ov} + \Delta M_{az} + \Delta M_{eddy} + M_{mix} + PER. \quad (11)$$

633 **Acknowledgements** We acknowledge use of the MONSooN system, a collaborative facility  
634 supplied under the Joint Weather and Climate Research Programme, a strategic partnership  
635 between the UK Met Office and the Natural Environment Research Council. We would also  
636 like to thank Matt Mizlienski for helping setup the model as well as, Jeff Blundell and Adam  
637 Blaker for technical support. Finally we wish to thank two anonymous reviewers for their  
638 comments which improved this manuscript.

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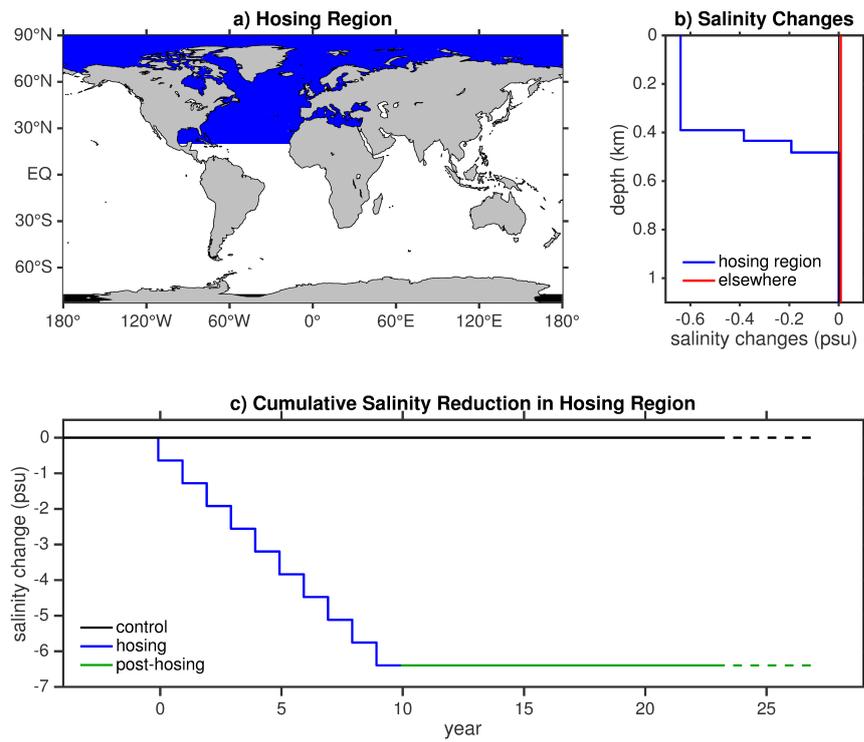
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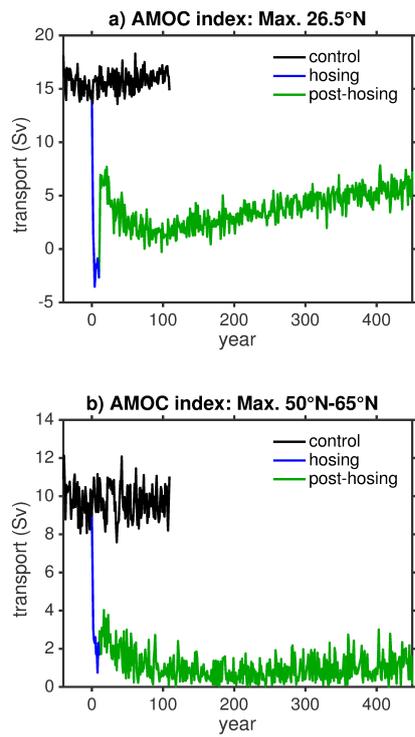
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**Table 1** Summary of freshwater budget for subtropical (10°N-45°N) and subpolar North Atlantic (45°N-70°N). All values are given in Sv with positive values indicating an addition of freshwater into the region. The bottom row of each section is the anomalous change in freshwater (i.e. Hosing (311-410) - Control).

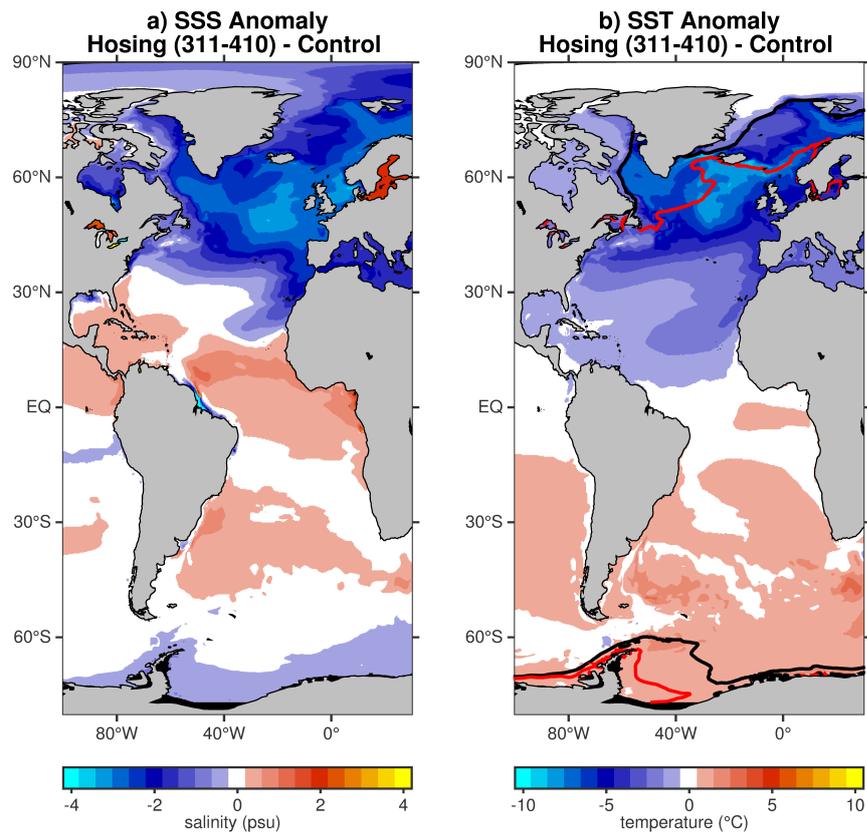
		<b>Overturning</b> $\Delta M_{ov}$	<b>Azonal</b> $\Delta M_{az}$	<b>Eddy</b> $\Delta M_{eddy}$	<b>P-E+R</b> $PER$	<b>Mixing</b> $M_{mix}$	<b>Total</b> $M_{trend}$
<b>Subpolar N. Atlantic</b>	Control	-0.170	-0.090	-0.033	0.241	0.041	-0.012
	Hosing (311-410)	-0.165	-0.187	-0.031	0.283	0.093	-0.007
	Anomaly	0.006	-0.097	0.001	0.042	0.052	0.005
<b>Subtropical N. Atlantic</b>	Control	-0.164	0.524	0.102	-0.504	-0.009	-0.051
	Hosing (311-410)	-0.032	0.490	0.061	-0.551	-0.024	-0.055
	Anomaly	0.132	-0.033	-0.041	-0.047	-0.015	-0.004



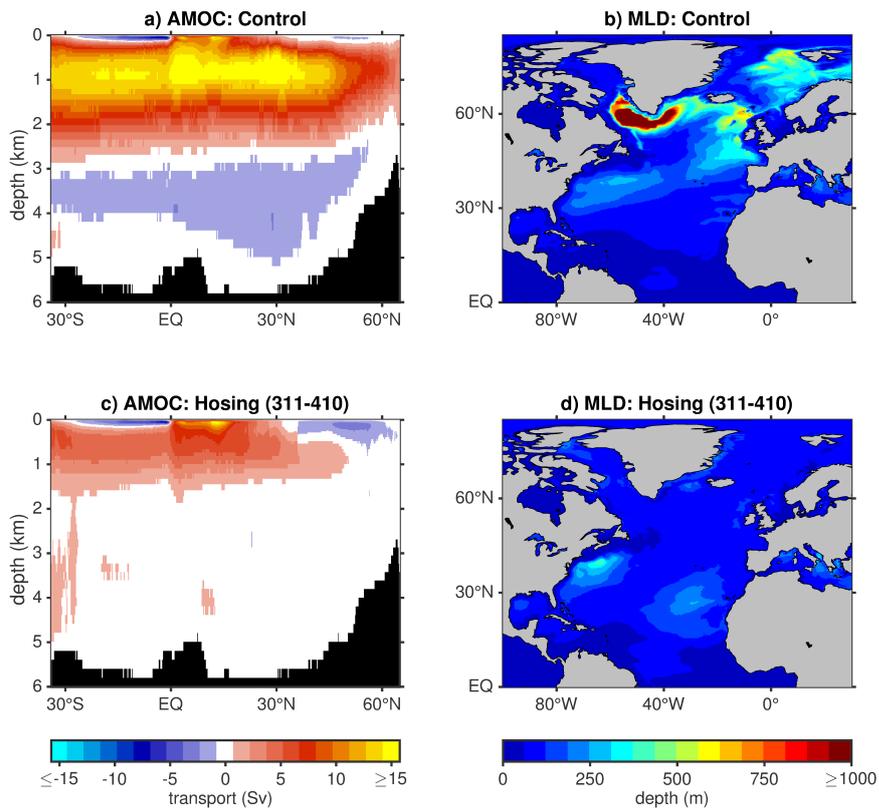
**Fig. 1** (a) The region where the freshwater hosing is applied. (b) The redistribution of salinity in the hosing region (blue) and everywhere else (red). (c) The cumulative salinity reduction in the hosing region (upper 350 m) in the model experiments for the control (black), hosing (blue) post-hosing (green).



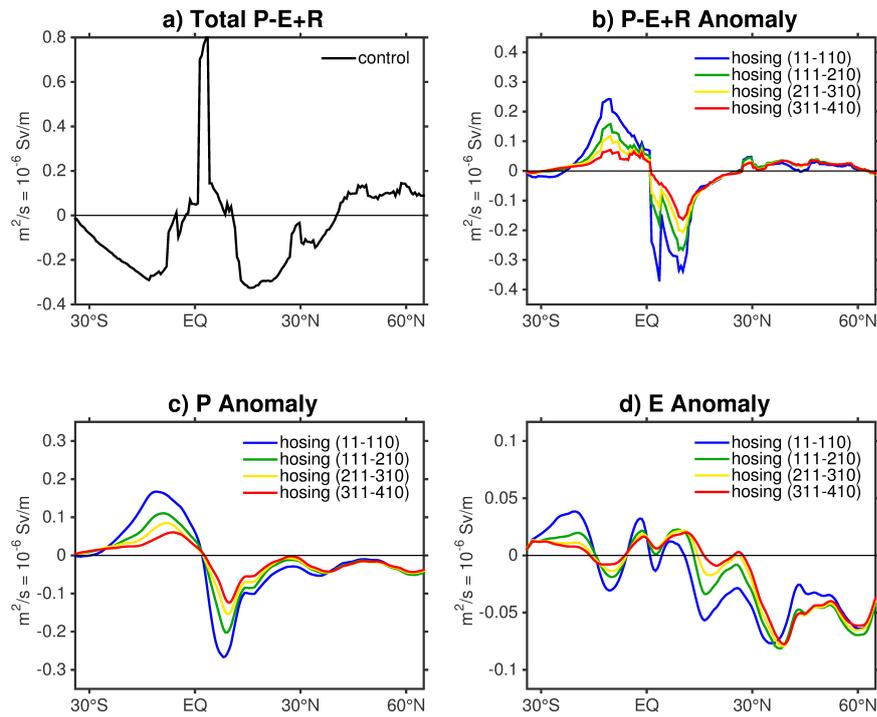
**Fig. 2** (a) The AMOC index computed as the maximum AMOC streamfunction at  $26.5^\circ\text{N}$  below a depth of 500 m and above 2000 m for the control experiment (black), hosing period (blue) and post-hosing period (green). (b) same as a expect computed between  $50^\circ\text{N}$  and  $65^\circ\text{N}$ .



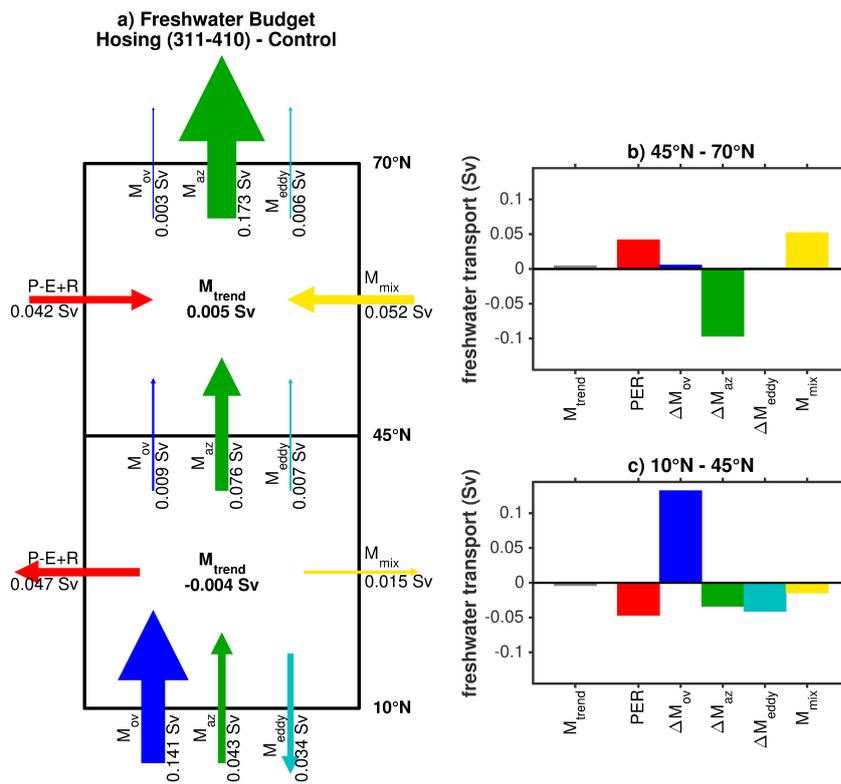
**Fig. 3** (a) Mean SSS from years 311-410 of the hosing simulation minus the mean SSS from the control simulation. (b) same as in (a) but for SST with the black contour indicating the annual maximum sea-ice extent in the control simulation and the red contour from years 311-410 of the hosing simulation.



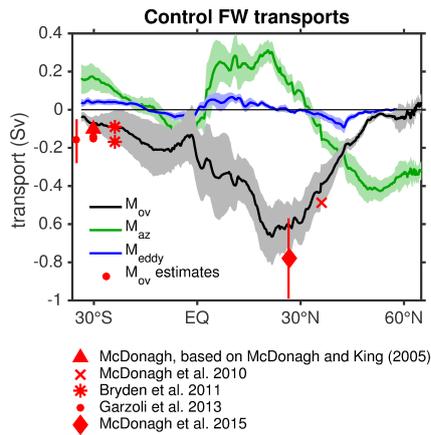
**Fig. 4** (a) The mean AMOC streamfunction and (b) the mean annual maximum mixed layer depth from the control simulation. (c) and (d) same as a and b but for years 311-410 of the hosing simulation.



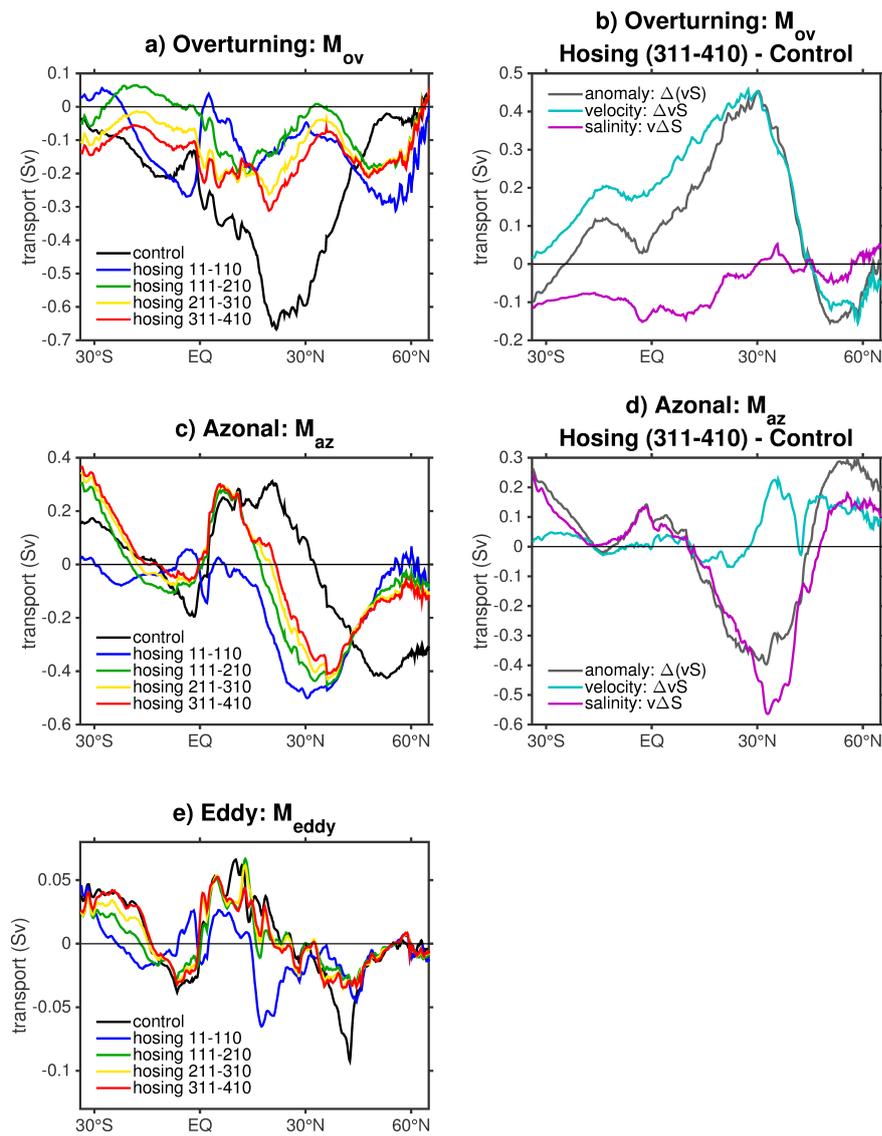
**Fig. 5** (a) The zonally integrated P-E+R from the control simulation normalized to Sv per meter in latitude. (b) the anomalous P-E+R from various 100 year means in the hosing simulation, (c) same as b but for precipitation only, (d) same as b but for evaporation only with blue years 11-110, green years 111-210 yellow years 211-310 and red years 311-410. All data is smoothed using a  $2^\circ$  latitude window to reduce the spikes from the river runoffs.



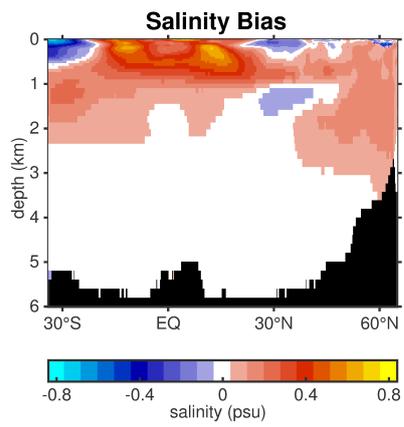
**Fig. 6** (a) Anomalous freshwater budget boxes for the subtropical (10°N-45°N) and subpolar (45°N-70°N) North Atlantic. The width of the arrows and arrow heads have been scaled according to the strength of the freshwater transport anomalies. (b) Summary of the anomalous freshwater budget for the subpolar North Atlantic. (c) Same as (b) but for the subtropical North Atlantic.



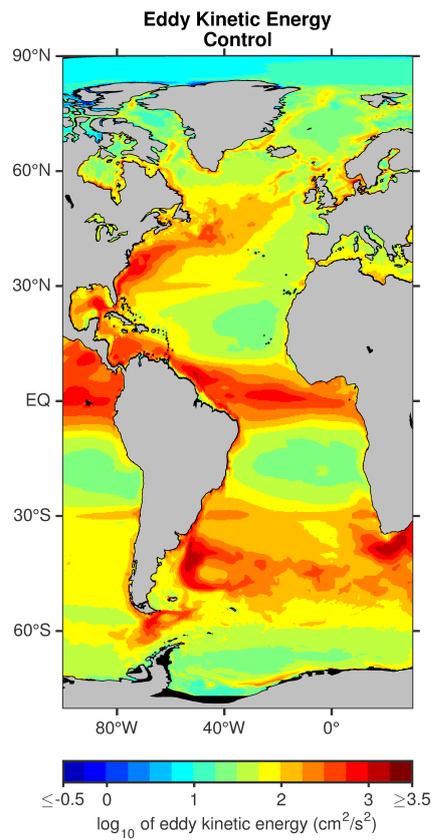
**Fig. 7** Mean  $M_{ov}$  from control simulation with  $\pm$  one standard deviation of seasonal data (black/grey shading), mean  $M_{az}$  from control simulation with  $\pm$  one standard deviation of seasonal data (green/green shading) and mean  $M_{eddy}$  from control simulation with  $\pm$  one standard deviation of seasonal data (blue/blue shading). Estimates of  $M_{ov}$  from observations (red): triangle based on McDonagh and King (2005); cross McDonagh et al (2010); stars Bryden et al (2011); circles Garzoli et al (2013) with vertical line representing the range in estimates; and diamond McDonagh et al (2015-in press) with the vertical line indicating the standard deviation of 10 day timeseries. Note that the standard deviations/range are computed using data available on different timescales.



**Fig. 8** The freshwater transports along latitude bands in the Atlantic. (a) Freshwater transport due overturning  $M_{ov}$ . The different colours represent different means over various years; control (black), hosing 11-110 (blue), hosing 111-210 (green), hosing 211-310 (yellow) and hosing 311-410 (red). (b) Decomposition of  $M_{ov}$  anomalies (hosing years 311-410 minus control) into contributions from velocity (cyan) and salinity (magenta) compared to total anomaly (dark gray). (c and d) same as (a and b) but for  $M_{az}$  and (e) same as (a) but for  $M_{eddy}$ . Note the different scales on panels a-e.



**Fig. 9** Zonal mean salinity bias of the control experiment relative to EN3 data (Ingleby and Huddleston, 2007).



**Fig. 10** Logarithm of the surface eddy kinetic energy in the control simulation. The eddy kinetic energy was computed from the model's surface velocity fields using the difference between the instantaneous velocities and seasonal mean velocities before averaging over all years of the simulation.