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- 2 Deglacial changes in flow and frontal structure through the Drake Passage
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- 14 Keywords: Cold Water Route; Sortable Silt; Deglacial; Alkenones; Sub-Antarctic Front
- 15 Abstract: The oceanic gateways of the Drake Passage and the Agulhas Current are critical locations
- 16 for the inflow of intermediate-depth water masses to the Atlantic, which contribute to the shallow
- 17 return flow that balances the export of deep water from the North Atlantic. The thermohaline
- 18 properties of northward flowing intermediate water are ultimately determined by the inflow of
- 19 water through oceanic gateways. Here, we focus on the less well-studied "Cold Water Route"
- 20 through the Drake Passage. We present millennially-resolved bottom current flow speed and sea
- 21 surface temperature records downstream of the Drake Passage spanning the last 25,000 years. We
- 22 find that prior to 15 ka, bottom current flow speeds at sites in the Drake Passage region were
- 23 dissimilar and there was a marked anti-phasing between sea surface temperatures at sites upstream
- 24 and downstream of the Drake Passage. After 14 ka, we observe a remarkable convergence of flow
- 25 speeds coupled with a sea surface temperature phase change at sites upstream and downstream of
- 26 Drake Passage. We interpret this convergence as evidence for a significant southward shift of the
- 27 sub-Antarctic Front from a position north of Drake Passage. This southward shift increased the
- 28 through-flow of water from the Pacific, likely reducing the density of Atlantic Intermediate Water.
- 29 The timing of the southward shift in the sub-Antarctic Front is synchronous with a major re-
- 30 invigoration of Atlantic Meridional Overturning Circulation, with which, we argue, it may be linked.

#### Introduction

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32 The export of North Atlantic Deep Water (NADW) to the South Atlantic requires a compensating 33 intermediate-depth northward flow. Volume transport estimates suggest that Antarctic 34 Intermediate Water (AAIW) is the main component of this upper layer return flow (Poole and 35 Tomczak, 1999). Thermocline waters contributing to Atlantic AAIW enter the Atlantic either through 36 Drake Passage (the Cold Water Route) or via the Agulhas Current (the Warm Water Route) (Poole 37 and Tomczak, 1999). The water mass properties of these AAIW contributions are markedly different, 38 with relatively cold-fresh water entering the southwest Atlantic via the Cold Water Route and warmer-saltier thermocline waters entering the southeast Atlantic via the Warm Water Route 39 40 (Gordon et al., 1992). The relative contribution and thermohaline properties of these water masses 41 determine the Atlantic's meridional heat and freshwater fluxes, and thus have implications for global 42 climate. Whilst there have been a large number of studies that have focused on the changes in the flow and 43 44 thermohaline properties of water masses entering the southeast Atlantic via the Warm Water Route, there are relatively few proxy reconstructions of past changes in Cold Water Route inflow. 45 The inflow of water via the Cold Water Route is strongly influenced by the position of the oceanic 46 47 fronts within the Drake Passage. Low-density Sub-Antarctic Mode Water (SAMW), formed in the southeast Pacific, is transported into the Atlantic basin north of the sub-Antarctic Front (SAF, the 48 49 most northerly jet of the Antarctic Circumpolar Current, ACC) (Talley, 1999). Along its flow path through the Drake Passage, SAMW mixes with Antarctic-derived waters to ultimately form the cold-50 51 fresh AAIW end-member in the Atlantic. Determining the position of the SAF is therefore critical for understanding changes in the inflow of low-density intermediate water via the Cold Water Route in 52 the past. 53 54 Sea surface temperature (SST)-based reconstructions of the SAF during the last glacial period suggest 55 that it was 5-10° northwards of its present position (e.g. Gersonde et al., 2005); however, 56 reconstructions of the position of the SAF in the Drake Passage region remain scarce. Estimates of 57 bottom current speed (McCave et al., 1995) provide a more direct means to determine relative 58 changes in Cold Water Route through-flow. To-date, the only two flow speed reconstructions from 59 within (Lamy et al., 2015) and downstream (McCave et al., 2014) of the Drake Passage are relatively low resolution across the last deglaciation and their interpretation of deglacial changes in flow via 60 61 the Cold Water Route are contradictory.

Here, we use a multi-proxy approach to determine changes in the structure and flow of water through the Cold Water Route over the last deglaciation. We present sub-millennially-resolved bottom current flow speed and alkenone-based SST reconstructions from a site immediately downstream of Drake Passage spanning the last 25 kyr. Comparison with other SST and bottom current flow speed records from the Drake Passage region reveals that the Last Glacial Maximum (LGM) and early deglacial was characterised by a very different spatial pattern of SSTs and bottom current flow speeds relative to the last 14 kyr, suggesting Cold Water Route through-flow during the glacial period was reduced.

### 2. Materials and Methods

### 2.1 Core Material

 This study is based on the sediment core GC528 (53°0.8′S, 58°2.4′W, 598 m) located on the southern flank of the Falkland Plateau collected on the cruise JR244 of the RRS *James Clark Ross*. The core was taken from a contourite drift (South Falkland Slope Drift; Koenitz et al., 2008), that is today located under the core of a strong westerly jet associated with the SAF (Figure 1) and is bathed by AAIW.

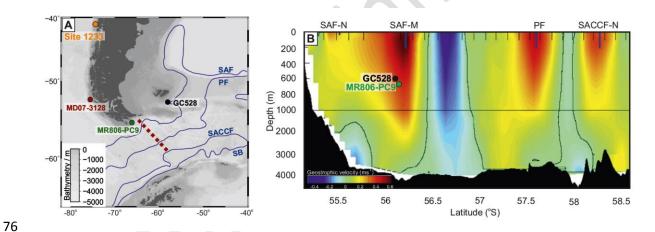


Figure 1: (A) Bathymetric map of Drake Passage showing the annual mean position of the circumpolar fronts in blue (from N to S: SAF = Sub-Antarctic Front; PF = Polar Front; SACCF = Southern Antarctic Circumpolar Current Front; SB = Southern Boundary of the ACC) (Orsi et al., 1995), core sites and location of flow speed profile shown in Fig. 1B (dotted red line); (B) Geostrophic flow velocities on transect through Drake Passage (Renault et al., 2011). Core positions have been projected onto this transect.

The top 50cm of GC528 (Holocene sediments) consist of foraminifera ooze containing >10% CaCO<sub>3</sub> and alkenone concentrations of >1  $\mu$ g/g. The rest of the sequence is mainly composed of silt-bearing clay with occasional dropstones present (particularly towards the LGM). Weight percent carbonate is ~1% and the alkenone concentration is 0.5-1  $\mu$ g/g.

87 The age model for core GC528 is based on reservoir-age corrected AMS <sup>14</sup>C dates on monospecific 88 benthic foraminifera and has been described in full in Roberts et al. (2016). The sedimentation rate 89 varies from 7 cm/kyr in the Holocene to 50 cm/kyr during the glacial. 90 2.2 Methods 91 2.2.1 Sortable silt grain size analysis 92 Sediment grain size distribution in the silt fraction (10-63 µm) is strongly controlled by the 93 geostrophic speed occurring above the bottom mixed layer (McCave et al., 1995). A detailed grain 94 size analysis of the silt fraction (with an average temporal resolution of 200 years), in which carbonate and biogenic silica had been removed (McCave et al., 1995), was performed using a 95 96 Coulter Counter (Multisizer 3) (Bianchi et al., 1999) with a 200 µm aperture. The set-up was adapted 97 following Moffa-Sanchez et al. (2015). The measurement error in the sortable silt mean grain size  $(\overline{SS})$  was 0.48 µm (1 S.D., n=8). 98 99 2.2.2 Ice-rafted Debris 100 Counts of terrigenous grains >300 µm (used as a proxy for coarse IRD) were made every 1cm in order 101 to determine potential changes in the source of sediment and the effect on  $\overline{SS}$ . Ice-rafted debris 102 comprises all grain sizes but we use counts of grains >300 µm as a proxy for IRD because grains in 103 this fraction are unlikely to be transported by other processes. 104 No obvious gravitational down-slope deposits were observed in the core, suggesting no influence 105 from downslope transport that could have otherwise emplaced sediment unmodified by current 106 transport at the site. Shards of Southern Andean volcanic ash or tephra have the potential to be 107 transported significant distances by the south-westerly winds (SWWs), therefore grains of tephra 108 were not counted. 2.2.3. Alkenone-derived Sea Surface Temperatures 109 Sea surface temperatures were calculated from the U<sup>K</sup><sub>37</sub> index (Prahl et al., 1988) from alkenone 110 111 analysis of core GC528 with an average sample resolution of 250 years. Lipids were extracted from ~3 g of homogenised, freeze-dried sediment using a CEM microwave system with 12 mL of 112 Dichloromethane (DCM):MeOH (3:1, v/v) (Kornilova and Rosell-Melé, 2003). Internal standards were 113 114 added for quantification ( $5\alpha$ -cholestane, dotriacontane and tetracontane). The relative abundances of di-, tri-, and tetra-unsaturated C<sub>37</sub> alkenones were measured with a Trace Ultra gas 115 116 chromatograph directly coupled to a Thermo DSQ single quadrupole mass spectrometer, fitted with a programmed temperature vaporising (PTV) injector. The target m/z were: 300 (nonadecanone), 117 118 544 ( $C_{37:4}$ ), 546 ( $C_{37:3}$ ), 548 ( $C_{37:2}$ ), 560 ( $C_{38:3Et}$  and  $C_{38:3Me}$ ), 562 ( $C_{38:2Et}$  and  $C_{38:2Me}$ ), 564 and 578 ( $C_{36}$ 

119 alkyl alkenoates) (Rosell-Mele et al., 1995). Several samples were split and extracted separately to determine the procedural error of  $U_{37}^{K} = \pm 0.034$  (n= 12; equating to an error in SST  $\pm 1.2$ °C). The 120 Prahl et al., (1988)  $U_{37}^{K}$ -SST calibration produced the best match between the surface sample and 121 122 modern annually-averaged sea surface temperatures at site GC528. This is because the Prahl et al., (1988) calibration is based on the  $U_{37}^{K}$  index (as opposed to the  $U_{37}^{K}$ ), which includes  $C_{37:4}$ . This 123 124 alkenone is abundant in high latitude regions and it has been shown that the U<sup>K</sup><sub>37</sub> index provides more robust SST reconstructions in high Southern latitudes than  $U^{K'}_{37}$  (Ho et al., 2012). 125 2.2.4.  $\delta^{13}C_{TOC}$  and C/N ratio 126 Total organic carbon  $\delta^{13}C_{TOC}$  and C/N ratio are used here to track changes in surface ocean 127 productivity and organic carbon source. Decarbonated samples (at 5 cm resolution) were measured 128 for organic carbon content following the method described in Könitzer et al., (2012). Percentage 129 130 carbon and nitrogen and  $\delta^{13}C_{TOC}$  analysis were performed by combustion on pre-weighed samples in an online system comprising a Costech ECS4010 elemental analyser (EA) coupled with a VG 131 132 TripleTrap and a VG Optima dual-inlet mass spectrometer at the NERC Isotope Geosciences Laboratory. Each analytical run contained 10 replicates of the internal NIGL standard BROC2 and 2 133 134 replicates of the external standard SOILB.  $\delta^{13}C_{TOC}$  is reported relative to the VPDB standard, with a 135 precision of 0.1% (1 S.D.). 136 2.3. Numerical Modelling A numerical model is used to determine the effect of a change in sea level on bottom current flow 137 138 speeds. The numerical model (Princeton Ocean Model, Blumberg and Mellor, 1987) is a regional 139 nested implementation of the Regional Ocean Modeling System. At the base there is a parent 140 model, which extends from 20°N to 70°S and from 180°E to 180°W with a horizontal resolution of ¼°. The parent model is nudged to the monthly mean climatological values of Simple Ocean Data 141 142 Assimilation (SODA) at its northern (open) boundary. Nested into this parent model there is a child 143 model, which covers all the southwestern Atlantic and southeastern Pacific sectors with a horizontal 144 resolution of 1/12°. The nested model configuration is forced with climatological mean wind stress 145 forcing derived from the ERA-Interim data set. Surface heat and freshwater fluxes are derived from 146 the COADS dataset. See Combes and Matano 2014 and Palma et al. 2008 for details. 147 An "LGM" sea level simulation was run such that the only difference from the control was a 120 m 148 reduction in sea level. Note that this "LGM" simulation does not take into account topographic 149 changes related to isostatic adjustment. The LGM and control simulations were spun-up to 150 dynamical equilibrium and run in diagnostic mode for 5 model years. Average values of the last year of the diagnostic run are used in the present analysis. 151

## 3. Results 152 153 3.1 Grain size analysis 154 The $\overline{SS}$ at site GC528 almost doubles across the last deglaciation, from a minimum of 18.2 $\mu$ m at 19-18 ka to a maximum of 34.7 $\mu$ m at 1 ka [Figure 2b]. There are two step-wise changes in $\overline{SS}$ across the 155 deglaciation; (i) between 21-19 ka, $\overline{SS}$ decreases by 5.2 µm, (ii) between 14.5-14.0 ka, $\overline{SS}$ increases 156 by 5.0 $\mu$ m. After 10 ka, there is a steady increase in $\overline{SS}$ spanning the entirety of the Holocene. 157 There is a $\overline{SS}$ -bottom current flow speed calibration for the Scotia – Weddell Sea region (SS=158 0.59U + 12.24) based on six points for which the sensitivity is 1.69 cm s<sup>-1</sup>/ $\mu$ m. However the greater 159 confidence in the universal sensitivity of sortable silt mean grain size to bottom current flow speed 160 based on 24 points of 1.47 ± 0.20 cm s<sup>-1</sup>/ µm (by Coulter Counter) (McCave et al., submitted) is 161 162 preferred. The total deglacial change in $\overline{SS}$ of 16.5 $\mu$ m (34.7-18.2 $\mu$ m) implies a change in bottom current flow speeds of 19.8-26.4 cms<sup>-1</sup>. Given that the largest grains in the silt size fraction (60 μm) 163 transition into an erosion regime in flow velocities greater than 25 cm s<sup>-1</sup>, the magnitude of this 164 deglacial change in $\overline{SS}$ must correspond to an increase in bottom current velocity from <5 cm s<sup>-1</sup> to 165 ~20 cm s<sup>-1</sup> (using the lower limit of sensitivity). 166 167 3.2 Fidelity of $\overline{SS}$ as record of bottom water flow 168 Several processes act on glacial-interglacial timescales that could affect the sedimentation at GC528, 169 170 such as changes in flow strength, sea level and ice-rafting, which we discuss here. 171 Reconstruction of the LGM Patagonian coastline (based on modern bathymetry) places it 200 km closer to GC528 than it is today (Figure S1), potentially increasing sediment supply. This could impact 172 on the reliability of $\overline{SS}$ as a bottom current flow speed proxy if the rate of direct sediment 173 174 introduction was faster than the current could rework. The sediment mass accumulation rate at 175 GC528 was higher during the LGM than the Holocene (Figure S1b), likely reflecting a closer sediment 176 source during glacial times as a result of the sea level lowstand. However, it is interesting to note 177 that the mass accumulation rate at GC528 does not decrease linearly with increasing sea level across 178 the deglaciation, suggesting that any relationship between sea level and sediment supply to GC528 is 179 not straightforward. Increased rates of sea level rise at 20-19 ka and 14.5-13 ka coincide with 180 elevated accumulation rates at GC528 (Figure S1b-c), perhaps suggesting that the initial flooding of

fluctuations on glacial-interglacial timescales. At site GC528, high concentrations of IRD are observed

At high latitudes, ice-rafted sediment provides additional material to the core site with significant

exposed shelf may have affected sediment supply to the slope.

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Figure 2: Deglacial grain size analysis, organic carbon and biomarker records from site GC528 in the southwest Atlantic. (A) EPICA Dome C (EDC)  $\delta D$  record on the AICC2012 age scale (Veres et al., 2013); (B) Mean sortable silt grain size  $(\overline{SS})$ ; (C) Counts of grains >300  $\mu$ m, used as a proxy for IRD; (D) Total organic carbon (TOC)  $\delta^{13}C$ ; (E) Organic carbon C/N ratio; (F) Alkenone-derived SST based on the  $U^k_{37}$  index (Prahl et al., 1988); (G) Total  $C_{37}$ :  $C_{38}$  alkenones used here to track changes in the alkenone-producing community. Grey and green bars show the typical range of E.huxleyi and G.oceanica  $C_{37}$ : $C_{38}$  ratio based on culture studies (Volkman et al., 1995). In each record, the  $2\sigma$  analytical precision is shown by the error bars, and a 1500 yr moving average and  $1\sigma$  moving standard deviation of each record is shown by the solid line and dashed lines respectively.

Do the variations in the sedimentation rate across the last deglaciation reduce confidence in  $\overline{SS}$  as an indicator of bottom current flow speed? The key assumption underlying the  $\overline{SS}$  proxy is that the sediment must be current sorted (McCave et al., 1995). Lamy et al., (2015) used the correlation between  $\overline{SS}$  and %SS to demonstrate current sorting at sites within the Drake Passage. In GC528, the cross-plot of  $\overline{SS}$  and %SS (measured on a subset of the samples, with a temporal resolution of 500 years) reveal that the sortable silt fraction of all sediments within GC528 has been current sorted, even within intervals that have been strongly influenced by IRD (Figure S2a), and there is no correlation between  $\overline{SS}$  and IRD concentration (Figure S2b). Therefore, we are confident that the  $\overline{SS}$  presented here can be interpreted as a reliable record of bottom current flow speeds.

## 3.3 Source of IRD in the southwest Atlantic

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Variations in IRD concentrations occur across the last deglaciation at site GC528 (Figure 2c). The two most significant potential source regions of this IRD are Patagonia and Antarctica. Acoustic imaging of the shelf sediments north of the Falkland islands have identified iceberg scours at present-day water depths of 240-480m (Brown et al., 2017). These authors argue that the depth of these scours requires large icebergs that could only have been derived from Antarctica. The anti-clockwise nature of the Antarctic coastal current means that icebergs calving off Antarctic glaciers typically accumulate within the Weddell Sea (Stuart and Long, 2011). From the Weddell Sea, most of the Antarctic icebergs escape into the ACC via "Iceberg Alley". Weber et al. (2014) identified several large iceberg discharge events from the Weddell Sea during the last deglaciation. However, consistent with icebergs drifting eastwards with the ACC from the Scotia Sea region, the timing of these iceberg discharge events (18-12 ka) do not coincide with the maximum IRD concentration at site GC528 (21-19 ka) (Figure S3), indicating that the source of IRD at GC528 was distinct from the Scotia Sea. We consider it likely that the dominant source of IRD deposited at GC528 was derived from marine-terminating glaciers on the west coast of Patagonia and transported through the Drake Passage. However, an exclusively Patagonian origin of IRD in this region is inconsistent with the inference of Brown et al. (2017) based on iceberg size, suggesting that icebergs sourced from East Antarctic may also have approached our site via the Drake Passage. A detailed IRD provenance study would be required to further test this hypothesis.

# 3.3 Productivity changes at site GC528

The  $\delta^{13}C_{TOC}$  reflects variations in both surface water productivity and source of organic matter, whereas the C/N ratio is primarily driven by the organic carbon source. In core GC528, organic carbon  $\delta^{13}C_{TOC}$  increases from values ranging between -23.0 to -22.6 % prior to 15 ka, to -21.5 to -21.0 % after 14 ka [Figure 2d]. In contrast the C/N ratio gradually decreases across the deglaciation

230 and Holocene with no marked shift in the C/N ratio at 15-14 ka [Figure 2e]. Taken together, we suggest the shift at 15-14 ka to more positive  $\delta^{13}C_{TOC}$  reflects a change in surface ocean productivity. 231 232 3.4 Alkenone-derived SST 233 Alkenone-derived SST across the last deglaciation is highly variable [Figure 2e]. The Holocene (10-0 234 ka) average SST (7.6 ± 2.0 °C) is in good agreement with modern annual average SSTs at the core 235 location. In contrast to other southern hemisphere alkenone records (e.g. Caniupán et al., 2011; Kaiser et al., 2005), the LGM (19-22 ka) average alkenone-derived SST at GC528 is relatively warm 236 237  $(7.0 \pm 3.9 \, ^{\circ}\text{C})$ . However, this average value masks the large fluctuations that are apparent in the 238 alkenone-SST record across the LGM. The LGM SSTs fluctuate between a minimum of 1.3°C at 25-23 239 ka and a maximum of 8°C at 19 ka, and there are two intervals spanning the end of the last glacial 240 period (20-18.5 ka and 23-22 ka) during which SSTs in the southwest Atlantic were elevated by 4-6°C 241 above the glacial mean. Replicate analyses on samples over this interval confirm that these warm 242 LGM temperature signals are robust. In the following, we evaluate factors that may result in the 243 elevated LGM alkenone-derived SSTs. 244 3.5 Fidelity of alkenone- U<sup>K</sup><sub>37</sub> as record of sea surface temperature Variations in the dominant alkenone-producing haptophyte could require the use of different 245 246 regressions between  $U_{37}^{K}$  and SST. The abundance of  $\Sigma C_{37}$ :  $\Sigma C_{38}$  alkenones within a sediment sample 247 has been used to identify shifts in the dominant alkenone-producing haptophyte population 248 (McClymont et al., 2005). We identify no significant shift in the  $\Sigma C_{37}$ :  $\Sigma C_{38}$  ratio over the interval studied [Figure 2f] with values averaging 1.16, in the range of Emiliania huxleyi (0.91-2.26) rather 249 250 than the other dominant alkenone-producing haptophyte Gephyrocapsa oceanica (0.59-0.81) 251 (Volkman et al., 1995). This suggests that no change in the dominant alkenone producing haptophyte population occurred, and thus the  $U_{37}^{K}$ -SST relationship should have remained constant. 252 253 Preferential degradation of components with a greater degree of unsaturation may bias the 254 alkenone-SST record towards warmer values (Flügge, 1997). We consider the percentage of the cold 255 tetra-unsaturated C<sub>37:4</sub> alkenone that would need to be removed in order to produce LGM SSTs as 256 low as during the late deglaciation period (2.2°C at 11.9 ka). We find that to account for a diagenetic 257 bias of 6°C, at least 70% of the 'original' C<sub>37:4</sub> would have to be removed from the samples. Such a significant amount of alteration has not been observed experimentally or within sediments studied 258 259 from this time period (Flügge, 1997). 260 Advection of warm water alkenones has been used to explain a warm LGM alkenone-SST signal in 261 some regions of the ocean. Modern alkenone studies from sites influenced by the Falkland Current 262 (Benthien and Müller, 2000) demonstrate that alkenone-derived SSTs from this region are biased

263 towards colder temperatures as a result of strong northward transport by the Falkland Current and 264 deep western boundary current. However core-top alkenone-SST data from site GC528 (7.7 °C) 265 shows good correlation with modern SSTs suggesting that advection does not strongly bias the 266 alkenone-SST signal at this site. 267 Reworking of alkenones from older warm intervals (such as the Eemian) could overprint the original 268 alkenone-SST and bias the alkenone signal to warmer temperatures. However, we do not believe 269 this to be a significant issue at site GC528 because (i) neither the TOC or the alkenone 270 concentrations during the two "warm" intervals during the LGM are significantly higher than the LGM average [Figure S3b-c], and (ii) other proxies such as planktonic  $\delta^{18}$ O [Figure S3d] show no 271 272 evidence of reworking. In addition, the spatial distribution of "warm" alkenone-derived LGM SST 273 records is globally distributed (Barrows et al., 2011). If all the sites were influenced by older 274 reworked alkenones, they would all have to be simultaneously affected by (specifically warm) 275 reworked alkenones during the LGM only. This seems an improbable explanation. 276 Changes in the seasonality of alkenone production can affect the temperatures recorded by the 277 alkenones. In high latitude regions, the limitation of light and the extreme surface conditions restrict 278 alkenone production to a short summer growth season (Ternois et al., 2000). Site GC528 is located in 279 the sub-Antarctic, and is sufficiently far north that it is not light limited. However, could there be 280 other processes in play during the LGM that affects the seasonality of alkenone production? Whilst it 281 is unlikely that sea ice was able to accumulate (even seasonally) across the Drake Passage, there 282 exist very indirect arguments supporting the idea that sea ice may have been present between the 283 Falkland Islands and Argentina during the LGM (Austin et al., 2013); however, we find no evidence of 284 any sea ice biomarkers (i.e. highly-branching isoprenoids (Collins et al., 2013)) in the glacial samples 285 of GC528. Alternatively, the accumulation of large icebergs in an "iceberg graveyard" around the 286 Falkland shelf (Brown et al., 2017) could have generated shallow meltwater stratification in the 287 surface ocean, potentially biasing summer SSTs (Peck et al., 2008). High IRD concentrations 288 coincident with both apparent warming events lend weight to this hypothesis [Figure 2c]. 289 Despite our concern that seasonality could have led to a warm bias in the glacial alkenone SST at site 290 GC528, a global comparison of seasonality in alkenone flux demonstrated that at a global scale, the 291 mean annual SST signal still dominates the sedimentary record (Rosell-Melé and Prahl, 2013). 292 Furthermore, where high latitude seasonality in alkenone production has been shown (Conte et al., 293 2006), the bias relative to the expected mean annual SST was only 2.5°C. Subtracting this value from 294 our glacial-stage warm intervals does not compensate the 4 °C anomalies relative to the other 295 samples; thus, the intervals of warming remain. As we show later (Section 4.1), seemingly logical 296 spatial trends in SSTs can be inferred from the comparison of this SST record with other alkenoneSST records in the Drake Passage region (see Section 4.1). We therefore argue that whilst a seasonality overprint of the glacial alkenone-SST record could have contributed to the magnitude of the warmings, the overall pattern in SSTs remains robust.

### 4. Discussion

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4.1 The 15-14 ka event

The  $\overline{SS}$  record from site GC528 shows a step-wise increase in bottom current flow speeds ( $\Delta U = 6.4$ cms<sup>-1</sup>) at 14.5-14 ka [Figure 2b]. Coupled with the change in bottom current velocity, there is evidence of a significant increase in  $\delta^{13}C_{TOC}$  [Figure 2d], potentially reflecting increased surface ocean productivity. Today, the Falkland Current (the most northerly ACC jet associated with the SAF) acts as the primary source of nutrients to the surface ocean above site GC528, supporting high concentrations of chlorophyll along the shelf break. While the Falkland Current remains in-situ, surface productivity is relatively unresponsive to changes in dust flux (unlike other sub-Antarctic sites). This idea is supported by our record of  $\delta^{13}C_{TOC}$ , which shows little response to the well-known decrease in dust between 18-17 ka (Lambert et al., 2008). This indicates that the prominent increase in  $\delta^{13}C_{TOC}$  at 14.5-14 ka was likely driven by reorganisation of the oceanic currents in the region. Oceanic current reorganisation would impact the supply of nutrients and/or the thermohaline properties of the surface ocean, affecting surface productivity. Furthermore, reorganisation would impact the bottom current speeds, so would be expressed in the  $\overline{SS}$  record. Understanding the nature of the inferred change in ocean circulation at 14.5-14 ka requires comparison with other sites in the region. Site MR806-PC9 is located at the northern margin of the Drake Passage [Figure 1a], and is today immediately upstream of site GC528 within the jet associated with the SAF [Figure 1b]. Comparison of the  $\overline{SS}$  records between site GC528 and MR806-PC9 (Lamy et al., 2015) reveals a remarkable similarity in bottom current speed at the two sites between 14-0 ka [Figure 3b]. In contrast, prior to 14 ka, there was divergence between the two  $\overline{SS}$ records [Figure 3b], with faster bottom current flow speeds in Drake Passage than on the South Falkland slope. The fact that sites GC528 and MR806-PC9 presently lie within the core of the main SAF jet [Figure 1b] and have similar bottom current flow speeds over 14-0 ka [Figure 3b] suggests a common response to changes in the intensity of the SAF jet. Both records suggest strengthening of the SAF jet through the Drake Passage over the late deglaciation and Holocene (which we discuss in Section 4.2). In contrast, the disparity in bottom current flow speeds prior to 14 ka - with faster flow speeds in Drake Passage relative to the downstream site - could be the product of one of two different

scenarios; (i) a lower glacial sea level meant that the SAF did not cross the topographic high of the

North Scotia Ridge and thus site GC528 was 'shielded' from the influence of the SAF, or (ii) the Southern Ocean frontal system shifted northwards relative to its current position such that the SAF did not extend through Drake Passage but instead was truncated by South America, analogous to the present Sub-Tropical Front. The faster flow speed at MR806-PC9 may then be the result of proximity to a more northerly position of the Polar Front.

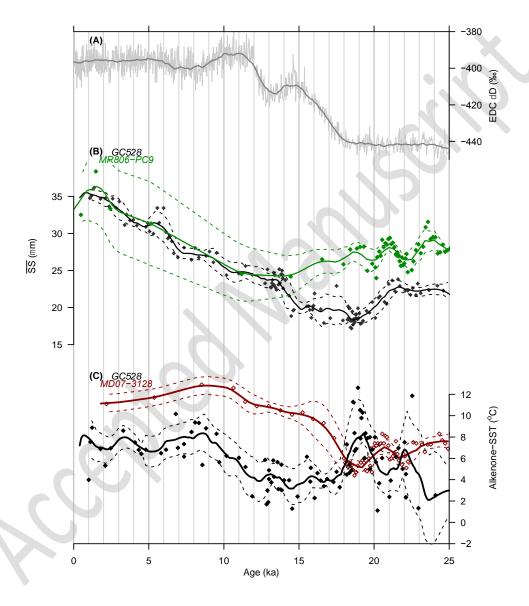


Figure 3: Comparison of proxy records from GC528 with upstream sites. (A) EPICA Dome C (EDC)  $\delta$ D record on the AICC2012 age scale (Veres et al., 2013); (B) Mean sortable silt grain size  $\overline{(SS)}$ , GC528 (black) and MR806-PC9 (green; Lamy et al., 2015); (C) Alkenone-derived SST records from upstream (MD07-3128 - red, Caniupán et al. 2011) and downstream (GC528 - black) of the Drake Passage. A 1500 yr moving average and  $1\sigma$  moving standard deviation of each record is shown by the solid line and dashed lines respectively.

In order to test Scenario 1 (a lower sea level shielded GC528 from the SAF), we modelled the effect of a lower sea level on bottom current flow speeds. The model predicts an overall increase in bottom current flow within the Drake Passage [Figure 4]. This increase is because the ACC is an equivalent barotropic jet and therefore, bottom current velocities are inversely proportional to the depth of the fluid. In contrast to the general strengthening within the Drake Passage, the model predicts a decrease in bottom current flow speeds of 10-15 cms<sup>-1</sup> at site MR806-PC9 [Figure 4].

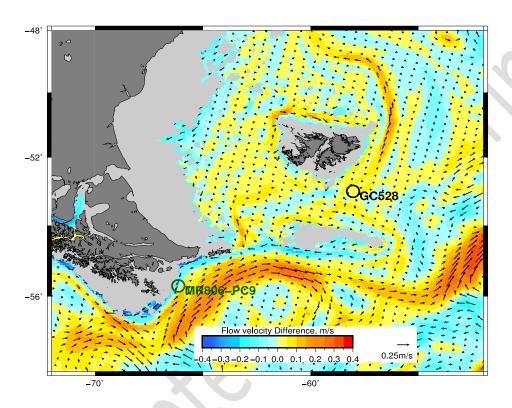


Figure 4: Model simulations of changes in bottom water flow speeds as a result of a lowering of sea level. Plot shows the relative change in bottom current flow velocity as a result of a -120m reduction in sea level. A reduction in bottom current velocity is highlighted by the blue colours. Green and black open circles show the position of MR806-PC9 (in Drake Passage) and GC528 (south of the Falkland Islands) respectively. Pale grey shows the model coastline in the low sea level scenario.

This is driven by the development of a frictional recirculation cell that occupies the space between the continent and the intensified offshore flow, and is highly dependent on the model configuration, and is therefore not a robust prediction. At site GC528, the model predicts an increase of 6 cm s<sup>-1</sup> in northward advection of bottom water during sea level lowstands [Figure 4]. This predicted increase in transport at GC528 is driven by a reduction in on-shelf transport, resulting in an increase in transport along the shelf-break. Because the Patagonian shelf is so expansive, a reduction in sea level results in a large reduction in on-shelf transport (and subsequent increase in shelf-break

transport) relative to the minor reduction in the transport across the North Scotia Ridge. Therefore a relative strengthening of bottom currents is predicted at GC528. The model predictions are inconsistent with our  $\overline{SS}$  reconstructions across the 15-14 ka transition, which suggest a significantly weaker bottom current before 15 ka at site GC528 and a stronger bottom water flow at MR806-PC9. Whilst we emphasise the uncertainty in the model prediction at MR806-PC9, we are confident in the model result at GC528. The discrepancy between the model prediction and the  $\overline{SS}$  reconstructions implies that sea level change is not the driver of the change in bottom current flow speeds observed in the records at GC528 (and possibly also MR806-PC9). Instead, oceanographic changes other than eustatic sea level change must have controlled the observed bottom current speed changes.

Front.

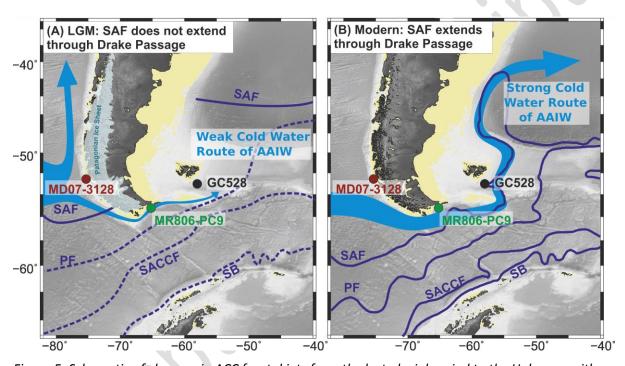


Figure 5: Schematic of changes in ACC frontal jets from the last glacial period to the Holocene with inferred inter-basinal exchange between Pacific and Atlantic. (A) Glacial ACC flow showing SAF located northwards of its present position and little inter-basinal exchange; (B) Holocene ACC flow showing the SAF extending through the Drake Passage and enhanced inter-basinal exchange. The positions of the ACC fronts are shown by the blue lines (more speculative locations denoted by dashed lines). The "Cold Water Route" water shown by light blue arrow. SAF = Sub-Antarctic Front; PF = Polar Front; SACCF = Southern Antarctic Circumpolar Current Front; SB = Southern Boundary

An alternative hypothesis invokes a change in the position of the oceanic fronts. LGM front reconstructions (e.g. Gersonde et al., 2005) suggest that the Southern Ocean fronts were located 5-10° northward of their present position. We propose that during the last glacial and early deglacial period (until 15 ka), the SAF did not extend through the Drake Passage but was instead truncated by

the South American continent (similar to the modern Sub-Tropical Front; Figure 5a). If the SAF did not extend through the Drake Passage, slower bottom current flow velocities at both GC528 and MR806-PC9 would be expected. The fact that bottom current flow speeds at MR806-PC9 are significantly faster than GC528 prior to 15 ka [Figure 4b] suggests that this is not the entire story, and may be influenced by a more northerly located Polar Front [Figure 5a] or a decrease in sea level (given the model uncertainty at this location).

Support for the idea of a northward shifted SAF can be found in the Pacific-Atlantic SST phase relationships over the last deglaciation. The comparison of alkenone-SST records from GC528 and two sites upstream of the Drake Passage [Figure 1a], site MD07-3128 (Caniupán et al., 2011) and ODP 1233 (Kaiser et al., 2005) reveals a striking SST anti-correlation prior to 15 ka [Figures 3c, 6c]. Intervals of warming downstream of the Drake Passage (e.g. 23-22 ka and 20-18 ka) coincide with significant cooling upstream. From 14 ka onwards, this anti-phase relationship is absent [Figures 3c, 6b]. In contrast, planktonic Mg/Ca-derived SSTs downstream of GC528, at the Brazil Margin site GL-1090 (24°S, 42°W; Santos et al., 2017) are in-phase with site GC528 throughout the last deglaciation [Figure 6]. The idea of a "thermal see-saw" between the southeast Pacific and the southwest Atlantic

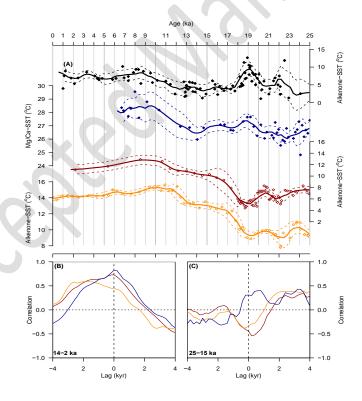


Figure 6. (A) Sea surface temperature records from GC528 (black) and GL-1090 (blue, Santos et al., 2017) from the southwest Atlantic, and sites MD07-3128 (red, Caniupán et al., 2011) and 1233 (orange, Kaiser et al., 2005) from the Chile Margin. (B) Cross-correlation between GC528 and MD07-3128 (red), GC528 and 1233 (orange), and GC528 and GL-1090 (blue) for the interval 14-2 ka; (C) Cross correlation for 25-15 ka, colours as in (B).

during the last glacial and early deglacial period (25-15 ka) supports the inference of northward shifted SAF during the LGM and early deglaciation. In the southeast Pacific, SSTs are controlled by the balance between cold sub-Antarctic Surface Water of the ACC advected northeastwards and warm surface waters transported southwards by the Chilean Coastal Current (Strub et al., 1998). Similarly, the SSTs in the southwest Atlantic represent a balance between cold southern-sourced water transported through the Drake Passage (Talley, 1999) and warm northern-sourced water carried southwards via the Brazil Current (Peterson and Stramma, 1991). Hypothesising that the SAF was absent from the Drake Passage during the LGM, a significant proportion of the cold water transported via the ACC to the Chilean margin would be redirected northwards into the Pacific (Lamy et al., 2015). This results in a relative cooling in the southeast Pacific and warming in the southwest Atlantic [Figure 5a]. Note that, based on our hypothesis, we argue that the warm LGM SST records from the Brazil Margin (Santos et al., 2017) are the result of a reduction in cold water advected through the Drake Passage, rather than an increase in warm water advected south from the Equator. The transition to in-phase SSTs upstream and downstream of Drake Passage at ~14 ka is consistent with the southward migration of the SAF to pass through the Drake Passage [Figure 5b]. A southward shift of the SAF would increase the inter-basin exchange of water between the Pacific and Atlantic, enabling temperature perturbations upstream of Drake Passage to be propagated rapidly to regions downstream of this gateway. In summary, sortable silt, SST and  $\delta^{13}C_{TOC}$  records from GC528 suggest a significant change in ocean circulation in the southwest Atlantic at 15-14 ka. Comparison of bottom current flow speeds downstream and within the Drake Passage provides evidence for a significant reorganisation of the frontal structure in the Drake Passage region across this interval. In particular, we suggest that the SAF did not extend through the Drake Passage until 15-14 ka. Anti-phased alkenone-SST records upstream and downstream of the Drake Passage support the idea of a reduced inter-basinal connection prior to 15 ka.

## 4.2 Strengthening of flow in the SAF after 14 ka

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Bottom current flow speeds downstream (GC528) and within the Drake Passage (MR806-PC9) evidence a significant increase ( $\overline{\Delta SS}$  of 11.5 µm suggests an increase of ~16 cm s<sup>-1</sup>) since 14 ka. Given that both sites record a very similar  $\overline{SS}$  record throughout this period, we infer that both sites are responding to an in-situ increase in the strength of the jet associated with the SAF. A strengthening of the SAF is supported by  $\epsilon$ Nd records from the Brazil Margin (Howe et al., 2016), which show increasingly radiogenic Nd isotopic signatures across the Holocene, indicating a greater through-flow of Pacific-derived AAIW.

The cause of such a large increase in the strength of the SAF jet is difficult to determine. The density difference either side of a front results in a strong pressure gradient, which is balanced by the Coriolis force producing a strong eastward jet. At a basic level, a stronger density gradient across a front will strengthen the associated jet (Thompson, 2008). The SAF is the location where cold-fresh dense Antarctic-derived water subducts beneath less dense water to the north to form Antarctic Intermediate Water (Hartin et al., 2011). The increase in the strength of the jet associated with the SAF may thus be the result of an increase in the density contrast between sub-Antarctic Surface Water north of the SAF and Antarctic surface waters south of the SAF. Proxy data of SSTs from north and south of the SAF do not provide strong support for the idea of an increased north-south temperature-driven surface density gradient during the Holocene. Sea surface temperature records of Antarctic surface water across the Holocene in the South Atlantic (Figure 7D; Nielsen et al., 2004) suggest an early Holocene warming followed by a cooler interval between 7-4 ka, and a relative warming from 4 ka to the present. In contrast, SST records north of the SAF from the Chilean margin suggest an initial warming in the Early Holocene followed by a general cooling trend (Figure 7C; Kaiser et al., 2005). Over the course of the Holocene, there is significant fluctuation in the north-south SST gradient; however, the long-term Holocene trend suggests a general decrease in the SST gradient [Figure 7E]. Based on this evidence, we cannot attribute the increase in the intensity of the SAF jet to an increase in a density gradient driven by temperature across the front. Note that the sites that we have used to reconstruct the SST gradient in Figure 7 are not proximal to the SAF and so may not accurately reflect SST gradients in the vicinity of the SAF. Alternatively, there remains the possibility of salinity-driven changes related to ice-melt, but we have no data to examine this aspect, nor (to our knowledge) are there any proxy data reconstructions of salinity across the SAF. However, it should be noted that the melt-related salinity gradients in the Southern Ocean were probably at a maximum during the deglaciation (18-12 ka) when the majority of ice was lost from Antarctica and Patagonia, and not during the Holocene. Finally, based on current understanding of the ACC, we observe that it is difficult to determine how changes in the intensity or position of the south westerly winds (SWW) might drive changes in the strength of the SAF jet. The SWWs transfer momentum into the surface ocean of the ACC, which sets up the barotropic component of flow (e.g. Allison et al., 2010). However, the exact relationship between the overlying wind forcing and the response of the ACC remains a matter of debate and is model dependent (Hogg et al., 2008 and references therein). Eddy compensation and eddy saturation result in non-linear responses of the ACC to changes in wind forcing, and jets are a finer scale that is not often well resolved in models. Furthermore, even if the response of the ACC could be predicted, proxy reconstructions of SWW strength show that the SWWs did not increase

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monotonically across the Holocene (Lamy et al., 2010). On balance, wind stress cannot adequately explain the seemingly linear increase in bottom current flow speeds suggested by data presented here

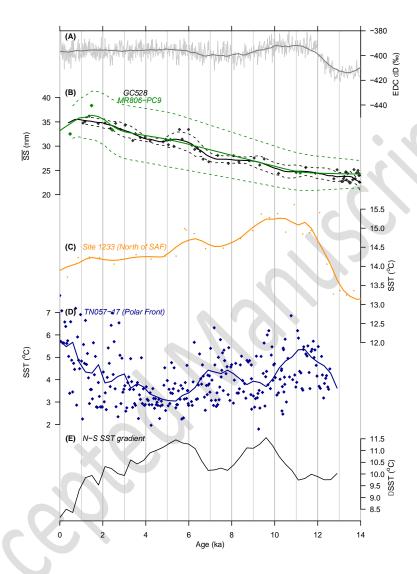


Figure 7: Reconciling the increase in strength of the SAF with changes in SST gradients over the Holocene. (A) EPICA Dome C (EDC)  $\delta D$  record on the AICC2012 age scale (Veres et al., 2013); (B) Mean sortable silt grain size, GC528 (black) and MR806-PC9 (green; Lamy et al., 2015); (C) Alkenone-SST record from the Chilean Margin, north of the SAF (orange; ODP site 1233, Kaiser et al., 2005); (D) Diatom transfer function based SST reconstruction from the Atlantic Sector of the Southern Ocean, at the Polar Front (blue; Core TN057-17, Nielsen et al., 2004); (E) Difference between the two SST records (C and D), used here as an indication of the density gradient across the SAF. A 1500 yr moving average and 1 $\sigma$  moving standard deviation of each record is shown by the solid line and dashed lines respectively.

In summary, the increase in bottom current velocities at sites in the Drake Passage and on the South Falkland slope since 14 ka suggests an increase in the strength of the SAF jet. However, the underlying cause of the increase in SAF jet strength remains enigmatic. Improved knowledge of sea surface salinity and temperature in the Drake Passage region could provide a key to understanding the increase in SAF jet flow speed across the Holocene. 4.3 Implications for the Atlantic Meridional Overturning Circulation (AMOC) A key idea presented here is that of a thermal seesaw between the southeast Pacific and southwest Atlantic during the last glacial period and early deglacial (25-15 ka). This, we argue, was related to a reduction in the interbasinal connection via the Drake Passage as a result of a northward shifted SAF. The following discussion concerns the wider effect of these changes on ocean circulation. The relative proportion of water flowing into the Atlantic basin via the Cold and Warm Water Routes affects the thermohaline properties of the shallow northward return flow. Extension of the SAF through the Drake Passage would have transmitted a greater volume of cold SAMW into the southwest Atlantic via Drake Passage, with potential consequences for global ocean circulation. The onset of the southward shift of the SAF at 15-14 ka is synchronous with a 'spin-up' of the AMOC during the Bølling-Allerød interstadial (McManus et al., 2004; Skinner et al., 2013). This AMOC spinup has commonly been interpreted to reflect changes in North Atlantic freshwater forcing, but the flow of water through the Drake Passage might also have played a role. Reduced throughput of cold low-salinity SAMW during the LGM would result in denser Atlantic AAIW (relative to today). If the density of Atlantic AAIW advected into the North Atlantic exceeds the density of NADW, then AMOC is potentially unstable (Keeling and Stephens, 2001). In contrast, a southward shift of the SAF at 15-14 ka would increase the transport of low-salinity SAMW into the Atlantic and decrease the density of AAIW relative to NADW, producing a more stable AMOC (Keeling and Stephens, 2001), and could explain the re-invigoration in AMOC observed during the Bølling-Allerød interstadial (McManus et al., 2004; Skinner et al., 2013). Support for this hypothesis comes from modelling studies (Weaver et al., 2003) which show that, given an initial circulation state in which the density of AAIW is greater than that of NADW, it is possible to spin up AMOC circulation by freshening AAIW in the vicinity of the Drake Passage. In this scenario, Drake Passage through-flow would act as a driver of AMOC circulation. 5. Conclusion

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In this study, detailed grain size data and SST records from sites along the northern margin of Drake

Passage were used to determine changes in the flow and frontal structures within the Drake Passage

since the LGM. In particular, we focus on the interval between 15-14 ka when an increase in bottom water current speeds is accompanied by increased in  $\delta^{13}C_{TOC}$  in the southwest Atlantic, indicating significant reorganisation of the currents. Comparison of bottom current flow speeds in the southwest Atlantic under the influence of the SAF jet to an upstream site in the Drake Passage reveals very similar trends after 14 ka, indicating that both sites are responding to changes in the strength of the SAF jet. In contrast, prior to 15 ka, bottom current speeds at the two sites are dissimilar. We propose that during the interval prior to 15 ka, the SAF did not extend through Drake Passage but instead lay further north and was truncated by South America, similar to the modern STF. This hypothesis is supported by alkenone-based SSTs upstream and downstream of the Drake Passage that suggest through-flow via the Cold Water Route during the last glacial was reduced relative to today. We suggest that this reduction in the inflow of low-density SAMW during the LGM potentially had wider reaching implication for AMOC; the higher density of Atlantic AAIW relative to NADW may have contributed to a more sluggish overturning circulation in the Atlantic. In contrast, the subsequent southward migration of the SAF at 15-14 ka enabled a greater through-flow of lowdensity SAMW into the Atlantic, reducing the density of Atlantic AAIW and potentially contributing to the spin-up of AMOC during the Bølling-Allerød interstadial. If correct, our findings have significant implications for the importance of the Drake Passage in controlling AMOC stability.

# Acknowledgments

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# **Supporting Material**

<u>Understanding the sortable silt flow speed proxy in marine systems with varying supply</u>

The sortable silt proxy is predicated upon the premise that the grain-size of deposited sediment at a point is primarily controlled by the short term (a few tens of years) time history of flow speeds at that point acting on a broad spectrum of fine sediment sizes delivered along the transport path. Control of deposition is expressed in the classical Krone equation (Einstein and Krone, 1962; Krone, 1962) for the selective deposition of fine suspended material. The amount deposited is given by:

 $\Sigma R_i t = \Sigma C_i w_{si} (1 - \tau_o / \tau_{di}) t$ 

where  $R_i$  is the rate of deposition (dimensions mass/area.time;  $ML^{-2}T^{-1}$ ), thus  $\Sigma R_i t$  is ( $ML^{-2}$ ), to be summed over i size fractions. The controls are; settling velocity  $w_{si}$ , critical depositional stress  $\tau_{di}$ , boundary shear stress  $\tau_o$  and the concentration  $C_i$  for each fraction according to the size distribution supplied in suspension from upstream.

This has been set out extensively in the publications of McCave, 2008, 2007; McCave et al., 1995; McCave and Hall, 2006. All the material that is considered in the analysis of the proxy is derived from terrestrial sources because care is taken to remove biogenic components (carbonate and silica) of marine origin. These sediments are delivered to the ocean beyond the shelf via a variety of routes, primarily by gravity flows (turbidity currents and debris flows) and shelf-edge resuspension/ spill-over, but also in a number of cases by aeolian fallout and via ice rafting in polar areas. Direct fluvial supply is rare, particularly at present under high sea-level stand. Even under glacially lowered sea-levels many deltas did not reach the shelf edge and many others debouched directly into the heads of submarine canyons, feeding turbidity currents which led into the deep sea (see reviews by e.g. McCave, 2002; Thomsen et al., 2002).

A test of whether sediments are current sorted or not is derived from cross-plots of  $\overline{SS}$  against SS%. Under a current sorted regime,  $\overline{SS}$  correlates positively with SS%, whereas unsorted sediments show no correlation. The latter case is evident in a fluvially-supplied lake (Gammon et al., 2017) or in the sediments from the Mississippi delta top (Xu et al., 2016) (Figure S5). In contrast, the marine sediment cores in the Drake Passage in Lamy et al., (2015) show a high degree of sediment sorting (Lamy et al., (2015), their Figure S3). Similarly, GC528 (presented in this study) shows a strong correlation between  $\overline{SS}$  and SS% (Figure S2a) suggesting that the sediment has been current sorted. Furthermore, even though there are high IRD concentrations in some intervals, there is no

correlation between IRD concentration and  $\overline{SS}$  (Figure S2b) and the samples containing IRD do not deviate from the  $\overline{SS}$ -SS% trend (Figure S2a).

## **Supplementary Figures**

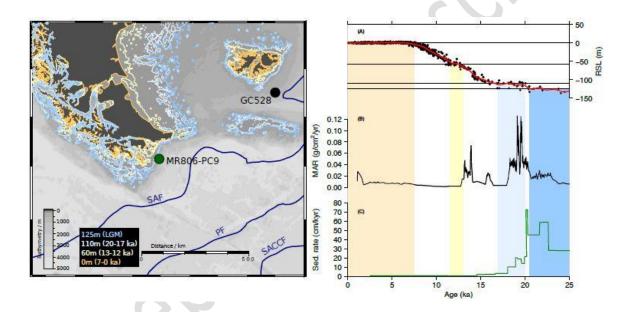
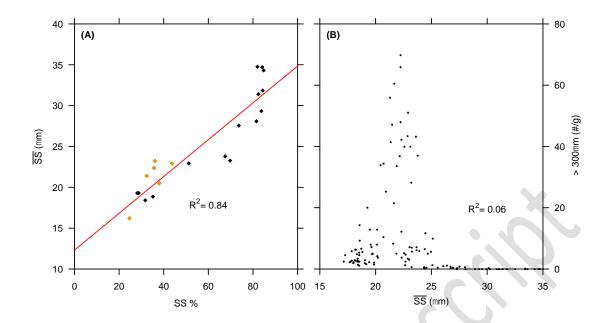


Figure S1: Assessing the impact of changes in sea level on the patterns of sedimentation. (Left) Map showing the position of the coastline at various intervals during the last deglaciation (see inset box) based on bathymetry. (Right) (A) Global relative sea level change over the last deglaciation (Lambeck et al., 2014); (B) Mass accumulation rate at site GC528; (C) Sedimentation rate at site MR806-PC9 (Lamy et al., 2015).



**Figure S2:** Assessing the current sorting in core GC528. (A) Cross-plot of weight percentage sortable silt fraction (10-63 $\mu$ m) versus the mean sortable silt grain size. Samples that contain significant quantities of IRD (grains > 300 $\mu$ m) are shown in orange. (B) Cross-plot of mean sortable silt grain size versus IRD (grains > 300 $\mu$ m) concentration.

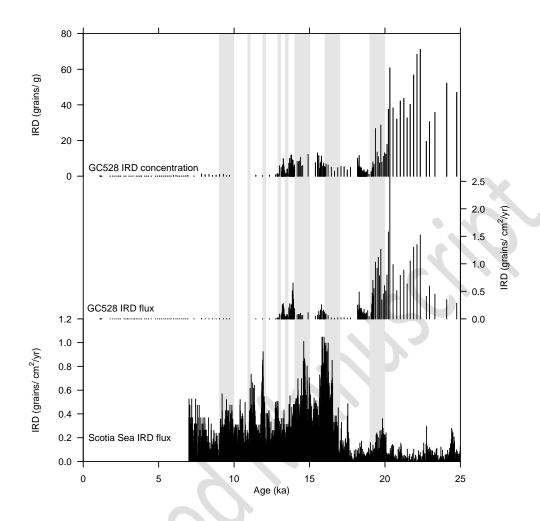
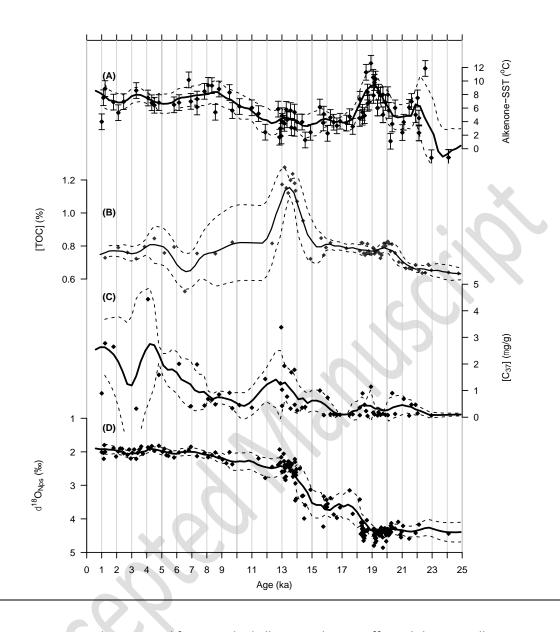
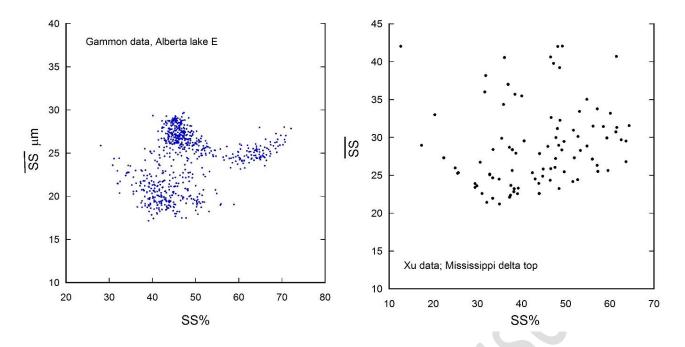


Figure S3: Comparison of the IRD records from the Southwest Atlantic (GC528) and the Scotia Sea

(Weber et al., 2014). (A) GC528 IRD concentration; (B) GC528 IRD flux (based on mass accumulation rates); (C) Scotia Sea IRD flux. Grey bars show the Antarctic Icesheet Discharge events (AIDS; Weber et al., 2014)



**Figure S4:** Assessing the potential for reworked alkenones having affected the LGM alkenone-SST record. (A) Alkenone-derived SST records from downstream (GC528 - black) of the Drake Passage; (B) Total organic carbon concentration in GC528; (C) Total alkenone concentration in GC528; (D) Planktonic foraminifera *Neogloboquadrina pacherma* (sinistral)  $\delta^{18}$ O from GC528.



**Figure S5:** Unsorted sediment cross-plots from the fluvially dominated Alberta Lake (left; Gammon et al., 2017) and the Mississippi delta top (right; Xu et al., 2016)

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