

1                    **Contribution of fine ash to the atmosphere from plumes associated with**  
2    **pyroclastic density currents**

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5    Samantha Engwell<sup>1</sup> & Julia Eychenne<sup>2</sup>

6    <sup>1</sup> Istituto Nazionale di Geofisica e Vulcanologia, Sezione di Pisa, Italy

7    <sup>2</sup> University of Bristol, School of Earth Sciences, Bristol, United Kingdom

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## 9 **1. Abstract**

10

11 Co-pyroclastic density current (co-PDC) plumes form as a mixture of fine-grained (<  
12 90  $\mu\text{m}$ ) particles and hot gas loft from the top of pyroclastic density currents. Such plumes  
13 can rise tens of kilometres high and inject substantial volumes of fine ash into the  
14 atmosphere, with significant implications for airspace disruption and populations, livestock  
15 and agriculture in downwind areas. Co-PDC deposits have a remarkably consistent grain size  
16 that remains constant with distance from source, regardless of eruption style, highlighting the  
17 complex sedimentation mechanisms that control deposition of co-PDC ash due to its fine  
18 grain size. Observations and numerical simulations of co-PDC onset emphasize the role  
19 played by the dynamics of PDCs in the development of co-PDC columns and plumes. The  
20 key differences between co-PDC and vent-derived plume source conditions and dispersion  
21 dynamics have important implications for application of remote sensing and numerical  
22 modelling methods.

23

## 24 **2. Introduction**

25

26 A critical requirement in assessing volcanic ash hazard and quantifying related risk is  
27 accurate identification and characterization of sources of atmospheric ash (Rose and Durant,  
28 2009; 2011). Ash injection and dispersion in the atmosphere during explosive eruptions is  
29 primarily associated with the formation of buoyant convective columns above volcanic vents  
30 (e.g. Bonadonna and Phillips 2003; Bursik 2001; Carey and Sparks 1986; Sparks 1986). In  
31 this context, a mixture of particles and gas is ejected from the conduit in a high velocity jet  
32 following fragmentation of magma ascending from depth in the conduit (Figure 1A).  
33 Convective plumes can also form from propagating pyroclastic density currents (PDCs).

34 PDCs are dry mixtures of gas, particles and entrained air that travel as fast moving density  
35 currents away from the eruptive vent (Figure 1B). While spreading, PDCs can generate  
36 secondary plumes by buoyant rise of hot gas and fine-grained material. Such plumes, here  
37 entitled ‘co-pyroclastic density current plumes’ (co-PDC plumes, also described as co-  
38 ignimbrite plumes and Phoenix clouds in the literature, e.g. Rosi et al. 2006; Sigurdsson and  
39 Carey 1989; Sparks and Walker 1977; Walker 1972) differ from vent-derived plumes in three  
40 principal ways (Figure 1B): 1) they are not characterised by a jet region at their base, 2) they  
41 initiate by buoyancy reversal of the material transported in the PDCs, and 3) they have a large  
42 source area.

43 PDCs occur frequently during explosive eruptions, and can be associated with years to  
44 decades-long low-level eruptive activities, including dome and lava flow-forming eruptions,  
45 as demonstrated by the recent activity at Soufriere Hills (Montserrat), Sinabung (Indonesia),  
46 Santiaguito (Guatemala), and Colima (Mexico). Observations of PDC emplacement, as well  
47 as analysis of pyroclastic deposits from recent and past eruptions, have shown that co-PDC  
48 plumes are associated with the entire size spectrum of PDCs, from hot rock falls and small  
49 volume PDCs (e.g. Bonadonna et al. 2002a; Calder et al. 2002; Eychenne et al. 2012), to  
50 massive expanded pumice and ash flows (e.g. Sparks and Huang 1980; Sparks and Walker  
51 1977). Co-PDC plumes can rise through the atmosphere to altitudes of tens of kilometres, as  
52 illustrated by those formed during the May 18, 1980 Mount St. Helens and June 15, 1991  
53 Pinatubo eruptions, which attained maximum heights of 30 km a.s.l (Holasek et al. 1996;  
54 Sparks et al. 1986). In fact, the Mount St. Helens co-PDC plume reached a greater height than  
55 the vent-derived plumes formed during the subsequent Plinian phase (Criswell 1987; Moore  
56 and Rice 1984). When co-PDC contributions have been distinguished from vent-derived  
57 tephra in fallout deposits, based on detailed studies of stratigraphy, grain size and  
58 componentry, co-PDC plumes have been shown to be capable of dispersing enormous

59 volumes of fine-grained ash over entire continents (Pyle et al. 2006; Rose and Chesner 1987).  
60 For example, the large magnitude ignimbrite-forming Youngest Toba Tuff (Indonesia) and  
61 Minoan (Santorini) eruptions deposited more than 700 km<sup>3</sup> and 38 km<sup>3</sup> of co-PDC ash,  
62 respectively (Rose and Chesner 1987; Sigurdsson et al. 1990). Co-PDC ash represents ~ 60  
63 vol% of the fallout deposit from the ignimbrite-forming Campanian eruption (Engwell et al.  
64 2014), and up to 58 wt% of the fallout deposit from the 2006 Subplinian eruption of  
65 Tungurahua (Eycheenne et al. 2012). Co-PDC plumes can thus contribute a significant  
66 proportion of the distal ash associated with very large, as well as intermediate sized,  
67 explosive eruptions.

68         These examples show that co-PDC plumes represent a substantial and recurring  
69 source of atmospheric ash during explosive eruptions, and constitute a significant and  
70 widespread hazard. Volcanic ash impacts communities on the ground by destroying crops and  
71 livestock, affecting infrastructure and causing health problems in humans. At Montserrat, the  
72 respiratory health hazards of co-PDC ash has been particularly important, due to the fine  
73 grain size and abundance of crystalline silicate phases (Baxter et al. 1999; Horwell and  
74 Baxter 2006; Horwell et al. 2001). In addition, the fine-grained nature of the ash allows it to  
75 remain suspended in the atmosphere for many days (Bonadonna et al. 2002a, 2002b; Rose  
76 and Chesner 1987), posing substantial risk for aviation (Gislason et al. 2011; Kueppers et al.  
77 2014). Co-PDC plumes can also impact the global environment, by contributing substantially  
78 to climate forcing by injection of aerosols, and specifically sulphurous phases, into the  
79 atmosphere, with implications for climate many years after eruption (Guevara-Murua et al.  
80 2014; Oppenheimer 2002).

81         There are a number of fundamental differences between co-PDC and vent-derived  
82 plumes in terms of, for example, grain size of the particulate load, particle concentration and  
83 ascent velocities (e.g. Bonadonna et al. 2002a; 2002b; Sparks 1976; Sparks et al. 1997;

84 Walker 1972). Distinguishing between the two plume types is therefore essential for hazard  
85 assessment, and proved crucial for investigating eruption dynamics during the 1991 Pinatubo  
86 eruption (Koyaguchi and Tokuno 1993) and for plume dispersion modelling at Montserrat  
87 (Bonadonna et al. 2002b). Accounting for co-PDC plumes when assessing hazard and risk  
88 related to volcanic ash dispersion requires comprehension of the specific processes  
89 controlling co-PDC plume formation, their injection into the atmosphere, and transport and  
90 sedimentation of co-PDC ash.

91         Accurate ash hazard assessment requires the integration of information from PDC, co-  
92 PDC and mixed fallout deposits (i.e. those that include both a co-PDC and a vent-derived  
93 contribution), with findings from analogue experiments in the laboratory and numerical  
94 simulations of PDC liftoff and formation of co-PDC columns. In this chapter, we review  
95 published data related to co-PDC plume formation, dispersion and sedimentation. In the  
96 following sections, we present an overview of dynamics and occurrences of co-PDC plumes  
97 and summarise observations from co-PDC plume-forming eruptions that encompass a range  
98 of eruption styles. Next, we summarize co-PDC deposit characteristics - specifically  
99 geometry, grain size and componentry - to highlight key spatial trends and provide insights  
100 for co-PDC plume formation, dispersion and sedimentation. We also summarise proposed  
101 theoretical models for co-PDC plume formation and insights gained from numerical  
102 modelling of co-PDC column development from propagating PDCs. Finally, we discuss  
103 appropriate considerations required for monitoring dispersion of co-PDC plumes using  
104 satellite infrared remote sensing methods, and for modelling co-PDC plume injection,  
105 dispersion and sedimentation processes.

106         In the following, the term “fallout deposits” refers to mixed deposits including tephra  
107 from both vent-derived and co-PDC plumes. We use the term “vent” to refer to the eruptive  
108 crater of a volcano, while “source” refers to the location of co-PDC plume liftoff. The term

109 “proximal” describes the areas affected by column dynamics and PDC emplacement, “distal”  
110 corresponds to all distances at which deposit thickness is less than a few centimetres to a few  
111 millimetres depending on the eruption magnitude, and “medial” refers to all intermediate  
112 distances.

113

### 114 **3. Overview**

115

116 Co-PDC plumes develop from propagating PDCs, and thus do not have their source  
117 centred on volcanic vents (Figure 1B). These plumes form as particles sediment from the  
118 base of propagating currents, and air is entrained at the top (e.g. Andrews and Manga 2011;  
119 2012; Bursik and Woods 1996; Sparks et al. 1997). The entrained air is heated by the  
120 particles, and expands. The combination of sedimentation and entrainment causes the density  
121 of the upper portions of PDCs to decrease below that of the surrounding atmosphere. This  
122 density decrease leads to a buoyancy reversal (Figure 1B; Bursik and Woods 1996; Woods  
123 and Wohletz 1991) and the formation of buoyant plumes. Buoyancy reversal can be enhanced  
124 by interaction of PDCs with topography, which favours sedimentation and air entrainment  
125 within the current (Andrews and Manga 2011; 2012). Co-PDC plumes ascend in the same  
126 manner as vent-derived plumes (Figure 1A), reach the level of neutral buoyancy and spread  
127 laterally by a combination of buoyant forces and atmospheric processes (Bonadonna and  
128 Phillips, 2003; Bursik et al. 1992a; 1992b; Sparks et al. 1986; Sparks et al. 1992). Formation  
129 of PDCs, vent-derived and co-PDC plumes can occur simultaneously during explosive  
130 eruptions, and even during a single eruptive event. Therefore, several sources may contribute  
131 to atmospheric ash injection during an explosive eruption, leading to complex mixed plumes  
132 and fallout deposits.

133

134            *3.1. Ignimbrite-forming eruptions*

135            The term ignimbrite refers to a voluminous PDC deposit formed by a very large  
136 eruption. The phenomena of co-PDC plumes was first highlighted in the context of  
137 ignimbrite-forming eruptions, where a distinct lack of fine-grained ash, and an enrichment in  
138 crystals in the matrix of the ignimbrite deposits, was used to suggest that gas and ash rich  
139 plumes lofted from the top of currents during their emplacement (Hay 1959; Lipman 1967;  
140 Sparks and Walker 1977; Walker 1972). Ignimbrite-forming eruptions typically have a  
141 magnitude  $M > 6.5$  (on the scale from Pyle (1995) where  $M = \log_{10}(\text{erupted mass, kg}) - 7$ ),  
142 and produce mixed fallout deposits which range from a few to thousands of  $\text{km}^3$  in DRE  
143 (Dense Rock Equivalent) volume (Table 1) and disperse several thousands of kilometres  
144 away from source (Pyle et al. 2006; Rose and Chesner 1987). Such large eruptions have not  
145 occurred during historical times, and their eruption processes are therefore inferred by  
146 analysis of resultant deposits. This reveals that ignimbrite-forming eruptions typically initiate  
147 with development of vent-derived plumes, followed by or synchronous with extensive PDC  
148 propagation blanketing the topography (Branney and Kokelaar 1997; Sparks et al. 1973;  
149 Sparks and Wilson 1976). During emplacement of the PDCs, co-PDC plumes loft from the  
150 entire extent of the current, have source radii on the order of tens to hundreds of kilometres,  
151 and reach altitudes of several tens of kilometres (Costa et al. 2012; Woods and Wohletz  
152 1991).

153

154            *3.2. Plinian eruptions*

155            PDCs are common during Plinian eruptions, often forming by collapse of the vent-  
156 derived column. Several historic eruptions have provided detailed observations of co-PDC  
157 plume formation and dispersion. The Plinian eruption of Mount Pinatubo in the Philippines  
158 on June 15, 1991, was the second largest eruption of the 20<sup>th</sup> Century, and produced both

159 Plinian and co-PDC plumes (Holasek et al. 1996; Koyaguchi and Tokuno 1993), leading to a  
160 mixed fallout deposit covering an area of  $\sim 4 \times 10^5 \text{ km}^2$  with a DRE volume of  $2.5 \text{ km}^3$   
161 (Wiesner et al. 2004; Table 1). The voluminous pumice flows generated ( $2.1 - 3.3 \text{ km}^3$  in  
162 DRE volume) overflowed topography to about 10 km away from the vent and were channel-  
163 confined for the rest of their runout distances (Rosi et al. 2001; Scott et al. 1996). A PDC  
164 travelling in a valley to 16 km from the vent produced a co-PDC plume clearly identifiable  
165 from the ground and on satellite images (Holasek et al. 1996). The plume lofted from the  
166 entire length of the PDC as an elongate sheet-like column, and showed notable variation in  
167 diameter with height, with a narrowing directly above the PDC and a sudden increase of  
168 diameter between 500 m and 1 km from the ground (Figure 2A). This co-PDC plume rose to  
169 a height of 20.5 km and dispersed at altitudes 1.5 - 2 km lower (Holasek et al. 1996).

170 Lateral blasts, which are impulsive directed explosions of a part of a volcano due to  
171 sudden decompression of cryptodomes, generate extensive PDCs associated with efficient co-  
172 PDC plume formation (Druitt 1992). Perhaps the most famous example of co-PDC plume  
173 formation is that associated with the May 18<sup>th</sup> 1980 eruption of Mount St. Helens, U.S.  
174 (Table 1, Figure 2B). A devastating PDC generated by a blast travelled to 15 km from source  
175 and blanketed an area of  $\sim 600 \text{ km}^2$  to the north of the volcano in less than 6 minutes (Druitt  
176 1992; Hoblitt et al. 1981; Moore and Sisson 1981). As the current interacted with topography,  
177 multiple plumes ascended from the PDC at velocities between 25 and 70 m/s, and rose 6 to 8  
178 km-high (Hoblitt 2000; Moore and Rice 1984; Sparks et al. 1986). The final plume (Figure  
179 2B) lofted from the entire area covered by the PDC and had a strong horizontal velocity  
180 component, presumably inherited from the northward directional energy of the blast  
181 (Eychenne et al. 2015). The plume was characterized by ascent velocities of  $\sim 110 \text{ m/s}$  and a  
182 maximum height of 35 km (Holasek and Self 1995; Sparks et al. 1986). It spread at an

183 altitude of 10 to 12 km (Eycheenne et al. 2015), controlled by gravitational processes for the  
184 first 100 km from source (Bursik et al. 1992a).

185

### 186 *3.3. Intermediate size eruptions*

187 Vulcanian and Subplinian eruptions, with magnitudes typically smaller than 5 (Table  
188 1), are characterized by unsteady vent-derived columns leading to formation of small volume  
189 (a few tens of km<sup>3</sup> in DRE volume), channelized, pumice and scoria flows (e.g. Bonadonna et  
190 al. 2002a; Calder et al. 1997; Cole et al. 2014; Hoblitt 1986; Stinton et al. 2014). Coeval  
191 formation of both vent-derived and co-PDC plume is common during such eruptions. Pumice  
192 flows generated by the short-lived, pulsating eruption of Mount St. Helens on August 7, 1980  
193 (Hoblitt 1986) produced co-PDC plumes with maximum vertical velocities of 5 to 13 m/s,  
194 despite initial velocities of only 1 to 2 m/s (Calder et al. 1997). Interaction of the pumice  
195 flows with topographic changes, either breaks in slope or bends in the channel, triggered  
196 formation of multiple individual plumes along the length of a single flow (Hoblitt 1986).  
197 With initial radii of ~ 50 to 200 m, the plumes merged during ascent and attained heights of  
198 400 to 800 m (Calder et al. 1997).

199

### 200 *3.4. Dome Collapse events*

201 Co-PDC plumes frequently rise from small volume, channel-constrained, PDCs  
202 generated by dome (e.g. Bonadonna et al. 2002a; Calder et al. 2002; Cole et al. 1998;  
203 Miyabuchi 1999; Stinton et al. 2014; Watanabe et al. 1999; Woods and Kienle 1994) and lava  
204 flow collapses (e.g. Evans et al. 2009). These collapses occur as viscous magma that has been  
205 extruded at the vent becomes gravitationally unstable, producing PDCs without development  
206 of an eruptive column at the vent. These PDCs often comprise a dilute upper part forming a  
207 cloud of hot gas and ash, which can detach from the denser part of the current on interaction

208 with topography (Watanabe et al. 1999). Discrete plumes develop sequentially from this  
209 dilute upper cloud and rise as a line of plumes that later coalesce during ascent, as observed  
210 during dome collapse activity at Soufriere Hills, Montserrat, in 1997 (Figure 2C; Bonadonna  
211 et al. 2002a; Calder et al. 2002; Cole et al. 2002). At Montserrat, the resultant co-PDC  
212 plumes occasionally detached from the source PDCs, and rose to heights of a few hundred  
213 metres to 15 km before dispersing many kilometres downwind. Co-PDC plumes associated  
214 with dome collapse events at Mount Redoubt, Alaska, in 1990 (Figure 2D) reached heights of  
215 12 km, had ascent velocities on the order of 10 m/s (Woods and Kienle 1996), and dispersed  
216 ash 100 to 200 km from vent (Scott and McGimsey 1994).

217

### 218 *3.5. Summary*

219 Observations made during historical eruptions of various styles and scales provide  
220 critical information about the dynamics of co-PDC plumes. Coeval formation of both vent-  
221 derived and co-PDC plumes during a single event is common. Co-PDC plumes can either  
222 form as multiple discrete plumes lofting from the top of channel-constrained PDCs before  
223 coalescing at height to form a single plume, or as liftoff occurs from the entire extent of  
224 PDCs. Interaction of currents with topography changes enhances the development of co-PDC  
225 plumes.

226 The effective source area of co-PDC plumes is much larger than vent-derived  
227 columns (Figure 1), and is generally controlled by the width of the parent PDCs. Reported  
228 ascent velocities are variable, on the order of many tens to hundreds of m/s for large  
229 magnitude eruptions, and of  $\sim 10$  m/s for smaller events. Maximum co-PDC plume heights  
230 can reach more than 30 km during large eruptions, while they are commonly on the order of  
231 hundreds of meters to a few kilometres during intermediate to small events. The examples of  
232 co-PDC plumes formed during Plinian events were not greatly affected by wind, and show

233 significant differences between maximum and neutral buoyancy heights (Figure 1), similar to  
234 those observed in strong vent-derived plumes (Bonadonna and Phillips 2003). Initial intrusion  
235 of the plumes into the atmosphere occurred radially, regardless of wind direction (Bursik et  
236 al. 1992a; Holasek et al. 1996; Sparks et al. 1986; Woods and Kienle 1994), which implies  
237 that initial spreading velocities were greater than typical stratospheric wind speeds. Co-PDC  
238 plumes from smaller dome collapse, Vulcanian and Subplinian events are strongly affected  
239 by wind, showing a classic bent over shape (e.g. Calder et al. 1997) typically seen in wind-  
240 affected vent-derived plumes.

241

#### 242 **4. Characteristics of co-PDC deposits**

243 In this section, the characteristics of co-PDC and mixed fallout deposits are described  
244 for eruptions where different co-PDC contributions have been inferred (Table 1).

245

##### 246 *4.1. Stratigraphy*

247 Most pyroclastic deposits associated with large magnitude ignimbrite-forming  
248 eruptions can be described by a simplified stratigraphy, that includes a coarse-grained (10's  
249 cm to a few mm) Plinian vent-derived deposit at the base (e.g. Rosi et al. 1999; Paladio-  
250 Melosantos et al. 1996; Scott et al. 1996), that is overlain by PDC deposits of up to hundreds  
251 of metres thick in proximal locations, with runout distances on the order of many tens of  
252 kilometres, covering areas of many thousands of square kilometres and representing volumes  
253 up to hundreds of cubic kilometres (see Table 1 for examples and references). Co-PDC  
254 deposits were identified both overlying and inter-bedded within individual PDC units, and  
255 were described as fine grained (<125  $\mu\text{m}$ ) ash layers rich in glassy particles (Hay 1959;  
256 Sparks et al. 1973; Sparks and Walker 1977). Ubiquitous accretionary pellets (following the  
257 definition of Brown et al. 2012) have also been found in co-PDC layers (e.g. Brown and

258 Branney 2004). During ignimbrite-forming eruptions, proximal Plinian and co-PDC deposits  
259 may be eroded or buried during emplacement of PDCs and later eruptive episodes, such that  
260 an entire eruptive sequence is often not observable at a single outcrop. At distances beyond  
261 the reach of the PDCs, it is possible to identify fine-grained ash layers corresponding to co-  
262 PDC deposits directly overlying Plinian deposits (Figure 3; Engwell et al. 2014; Ninkovich et  
263 al. 1978; Sparks and Huang 1980; Wiesner et al. 1995). With increasing distance from  
264 source, as the grain size of the Plinian deposit fines, it becomes increasingly difficult to  
265 differentiate the two deposits, both in terms of their stratigraphy and grain size (Figure 3),  
266 and a single layer of ash is generally observed.

267 Pyroclastic deposits from PDC-forming, Plinian to Vulcanian eruptions are  
268 considerably smaller in extent and volume (see Table 1 for examples and references). In such  
269 examples, PDC deposits occur as multiple, often channelized, lobes and fans, producing  
270 stratigraphic sequences that vary significantly over spatial scales of a few tens of kilometres  
271 (Bernard et al. 2014; Brand et al. 2014; Charbonnier and Gertisser 2011; Druitt 1992; Hoblitt  
272 1986). Distinct co-PDC ash layers with similar characteristics to those associated with  
273 ignimbrite deposits have been documented close to source on top of sequences of smaller  
274 PDC units (Cole et al. 2002; Criswell 1987; Eychenne et al. 2012; Kuntz et al. 1981;  
275 Miyabuchi 1999; Ritchie et al. 2002). These co-PDC deposits appear as structureless layers  
276 of fine ash ( $< 90 \mu\text{m}$ ), millimetres to a few centimetres-thick, often bearing accretionary  
277 pellets (Eychenne et al. 2012; Ritchie et al. 2002; Watanabe et al. 1999). Beyond PDC runout  
278 distances, distinct co-PDC ash layers are not easily identifiable based on stratigraphy alone,  
279 due to coeval formation and interaction of vent-derived and co-PDC plumes, common during  
280 Vulcanian, Subplinian and Plinian eruptions (Bonadonna et al. 2002a; Di Muro et al. 2008;  
281 Eychenne et al. 2012). When vent-derived and co-PDC plumes form during the same eruptive  
282 episode, medial to distal fallout deposits occur as massive tephra beds, comprising a mixture

283 of co-PDC and vent-derived material (Bonadonna et al. 2002a; Eychenne et al. 2012; 2015;  
284 Sarna-Wojcicki et al. 1981). An important signature of these unstratified mixed fallout  
285 deposits is their bimodal grain size distributions, which are further described in section 4.4.  
286 Dome collapse events, which lack vent-derived plumes, offer the best opportunities to  
287 observe and document co-PDC ash sheets to distal locations (Bonadonna et al. 2002a; 2002b;  
288 Watanabe et al. 1999).

289

#### 290 *4.2. Geometry and volume of co-PDC and mixed fallout deposits*

291 Co-PDC deposits formed during dome collapse events (Bonadonna et al. 2002a;  
292 Watanabe et al. 1999) have elongate and rather broad proximal regions located away from the  
293 eruptive vent (Figure 4). This reflects the geometry of the co-PDC columns as they rise from  
294 the length of the channelized PDC (Bonadonna et al. 2002a; Calder et al. 1997; Cole et al.  
295 2002), controlled by the orientation of the constraining drainage channel. Despite these  
296 controls however, the proximal footprint of co-PDC deposits appears to depend primarily on  
297 wind, with isopachs displaced, and sometimes even detached, from the area of PDC  
298 emplacement in the downwind direction (Figure 4). In medial and distal areas, these proximal  
299 effects become insignificant, and deposit geometries and axes are ultimately controlled by the  
300 atmospheric winds (Figure 4A, B and C). A number of co-PDC deposits show secondary  
301 maxima in thickness and mass at distance downwind from the vent (Figure 4D; Bonadonna et  
302 al. 2002a; Watanabe et al. 1999), a feature also observed in some mixed fallout deposits  
303 (Figure 5A; Eychenne et al. 2015).

304 In examples of mixed fallout deposits where the co-PDC and vent-derived  
305 contributions were distinguished (based on componentry and grain size), enrichment in co-  
306 PDC ash has been shown to strongly impact the footprint of the mixed fallout deposit.  
307 Contribution of co-PDC plumes rising from large source areas, such as those that form in

308 association with emplacement of PDCs that are not controlled by topography (Table 1), can  
309 lead to mixed fallout deposits with broad proximal geometries (Figure 5A; Eychenne et al.  
310 2015). In addition, PDCs derived from lateral blast events can produce co-PDC plumes with a  
311 strong horizontal component, leading to the up- and/or cross-wind displacement of the mixed  
312 fallout deposit (Figure 5A; Eychenne et al. 2015). Localized enrichment in co-PDC ash,  
313 which can occur downwind of a topographic barrier perpendicular to the PDCs propagation  
314 direction, can cause an over-thickening of the mixed fallout deposit leading to complex  
315 deposit shapes (Figure 5B; Eychenne et al. 2012).

316 Volumes of co-PDC deposits formed during dome collapse events are typically small,  
317 ranging from  $2 \times 10^{-5}$  to  $9 \times 10^{-4}$  km<sup>3</sup> (Table 1), and accounting for 4 to 30% of the volume of  
318 associated PDCs (Bonadonna et al. 2002a; Watanabe et al. 1999). Despite such small  
319 volumes, the fine-grained nature of co-PDC ash allows its dispersion and deposition over  
320 much larger areas than would be expected for a vent-derived eruption of the same magnitude.  
321 Co-PDC deposits associated with large magnitude ignimbrite-forming eruptions have been  
322 shown to have similar volumes to deposits from vent-derived plumes formed during the same  
323 eruption (Table 1 and references therein). Their volumes range from 6 km<sup>3</sup> to 2000 km<sup>3</sup>. In  
324 some examples, inferred co-PDC volume is considerably larger than the volume of material  
325 dispersed by vent-derived plumes (Table 1).

326

#### 327 *4.3. Componentry of co-PDC deposits*

328 The matrix of ignimbrite deposits is typically enriched in phenocrysts compared to  
329 scoriae or pumice from the same eruption (Hay 1959; Lipman 1967; Sparks 1976; Walker  
330 1972). This observation suggests that much of the vitric (glassy) material is not deposited by  
331 the PDCs and must instead settle as co-PDC ash (Hay 1959; Sparks et al. 1973; Sparks and  
332 Walker 1977; Walker 1972). Some crystals in the PDC matrix, however, originate from

333 fragmentation or comminution of accidental and accessory material (Lipman 1967).  
334 Therefore, the extent to which the degree of crystal enrichment in PDCs can be used to infer  
335 the volume of co-PDC deposits is limited.

336 In fact, co-PDC ash layers are not solely composed of vitric glass; they may also  
337 contain a non-negligible proportion of crystals (3 to 14 wt% in grain size fractions < 2 mm)  
338 and non-juvenile material (1 to 12 wt% in grain size fractions < 2 mm, Sparks 1975; 1976;  
339 Sparks and Walker 1977). In co-PDC deposits formed during dome collapse events, where  
340 the juvenile material is microlite-rich, crystals can constitute up to 50 wt% of material coarser  
341 than 63  $\mu\text{m}$ , and up to 70 wt% of the finer material (Bonadonna et al. 2002a; Watanabe et al.  
342 1999). In such samples, micro-crystalline silicate phases appear concentrated in co-PDC  
343 deposits with respect to the matrix of the PDCs, especially in the sub-10  $\mu\text{m}$  fractions (Baxter  
344 et al. 1999; Horwell et al. 2001). Finally, co-PDC plumes associated with PDCs with a high  
345 content of non-juvenile material, such as blasts (which are accompanied by the collapse of  
346 volcanoes' flanks), can settle ash dominated by lithic material (Eycheenne et al. 2015).

347 These data show that crystals and various types of non-juvenile particles can  
348 constitute a considerable proportion of co-PDC deposits, suggesting that co-PDC plumes do  
349 not entrain solely low density material (such as glass shards). The componentry of co-PDC  
350 deposits is highly variable from one eruption to another, and appears to be controlled by the  
351 nature of the source material and fractionation processes occurring during PDC propagation,  
352 and not only by the processes controlling the formation of the co-PDC plumes, which will be  
353 discussed further later.

354

355                    *4.4. Grain size of co-PDC and mixed fallout deposits*

356                    The grain size of co-PDC deposits is characterized by a unique signature, which has  
357 proved a critical tool for distinguishing co-PDC from vent-derived contributions in mixed  
358 fallout deposits.

359                    *4.4.1. Methods for analysing particle size distributions of fallout deposits*

360                    Because of the fine-grained nature of co-PDC deposits (see section 4.4.2), their  
361 particle size distributions (PSD) cannot be resolved by traditional sieving methods, and a  
362 number of more complex techniques have been employed for collection of the grain size data  
363 compiled here (Table 1). The most common method in modern studies is laser diffraction  
364 (e.g. Di Muro et al. 2008; Eychenne et al. 2012; Miyabuchi 1999), which determines the size  
365 distribution of a suspension of particles by measuring the light intensity of a laser beam at  
366 low scattering angles passing through the suspension. This method applies different light  
367 scattering theories (e.g. Mie, Fraunhofer approximation), which require an assumption  
368 regarding the optical parameters (refractive index, light absorption) and shape of the particles  
369 (often approximated to spheres). Another widely used technique is electrozone sensing (e.g.  
370 Bonadonna et al. 2002a; Brazier et al. 1982; Sparks and Huang 1980), whereby particles  
371 suspended in an electrically conductive fluid are counted, and their size is inferred by  
372 measuring the electrical resistivity of the fluid while individual particles are forced through a  
373 small opening. After calibration with known size standards, this technique gives repeatable  
374 results for non-spherical particles. Finally, a number of studies have measured the  
375 sedimentation rate of individual particles falling through a liquid of known density, with  
376 particle size estimated assuming a settling velocity law (e.g. Watanabe et al. 1999; Wiesner et  
377 al. 2004). One limitation of this method arises from difficulties in obtaining adequate  
378 experimental conditions to ensure individual settling of particles finer than 63  $\mu\text{m}$ , and the  
379 lack of constraints on the flow regime controlling their settling. In addition, assumptions

380 concerning particle shape (generally spherical) and density (variability in material is often not  
381 considered) are required to compute the settling velocity law. All of these methods require  
382 some form of approximation, which do not encompass the complexity of volcanic ash (see  
383 Cashman and Rust's chapter), and therefore caution is advised when comparing PSDs  
384 determined by different techniques.

385 Commonly grain size analyses of mixed fallout deposits (including both vent-derived  
386 tephra and co-PDC ash) is conducted by combining sieve data (for material coarser than 90  
387  $\mu\text{m}$ ) with one of the methods described above (for measuring material finer than 90  $\mu\text{m}$ )  
388 (Table 1). To obtain a complete distribution, the two resultant datasets are merged by: 1)  
389 converting the PSDs obtained by laser and electrozone techniques from volume to mass,  
390 assuming a value of particle density, and 2) calibrating sieve data with the method used for  
391 measuring the fine grains, by overlapping the measurements in a given size range (Eycheenne  
392 et al. 2012; Wiesner et al. 2004).

393 The PSDs obtained by these methods are described by a probability density function  
394 using the Krumbian phi scale ( $\Phi = -\log_2(d/d_0)$ , where  $d$  is the diameter of the particle and  $d_0$   
395 is a reference diameter equal to 1 mm), widely adopted in geology (Folk and Ward 1957).  
396 PSDs from fallout deposits can be polymodal and rarely follow a mathematical function. In  
397 early studies, distinct peaks in polymodal PSDs were separated into individual  
398 subpopulations by hand (Brazier et al. 1983). More sophisticated methods to deconvolve  
399 polymodal PSDs assume that they consist of a combination of log-normal subpopulations  
400 (Wohletz et al. 1989). Here, we have deconvolved polymodal distributions (where not  
401 processed in the original studies, see Table 1) by fitting identifiable modes of the raw  
402 distributions without assuming that they follow a specific mathematical function. We use  
403 algorithms based on the versatile Weibull function, which can reproduce the shape of many  
404 probability density functions, and thus allows extraction of non-lognormal subpopulations

405 (Caballero et al. 2014; Eychenne et al. 2012; 2015). Unimodal distributions and deconvolved  
406 subpopulations are described using median diameter (Md) and sorting ( $\sigma$ ) parameters,  
407 determined following the graphical method of Inman (1952).

408

#### 409 *4.4.2. Grain size characteristics of co-PDC and mixed fallout deposits*

410 Despite differences in geological context (chemistry of the magma, crystal content,  
411 etc.) and eruptive dynamics (eruption style, magnitude, PDC type and trigger), the co-PDC  
412 deposits have very similar grain size distributions (Figure 6, Table 1). Generally unimodal,  
413 they have a median diameter (Md) between 3.5 and 6  $\Phi$  (90 and 15.6 microns) which does  
414 not vary with distance from source (Figure 6). These consistent co-PDC grain size  
415 characteristics suggest that processes controlling the entrainment of particles in co-PDC  
416 columns are fundamentally size-selective, with particles coarser than about 90  $\mu\text{m}$  not  
417 preferentially lofted.

418 Both large and intermediate magnitude PDC-forming eruptions for which grain size  
419 analyses of the fallout deposit have been published are listed in Table 1. Most of these fallout  
420 deposits are unstratified over the area sampled, and are assumed to contain both co-PDC and  
421 vent-derived ash. The other deposits comprise a coarse lower (10's cm - mm) and fine upper  
422 (< 125 microns) unit in proximal locations, and a massive ash layer at greater distances  
423 (Figure 3). The grain size data presented here correspond to the vertically averaged deposit  
424 (where grain size of individual units are weighted according to the units thickness). All of the  
425 fallout deposits have bimodal particle size distributions at individual locations in the deposit,  
426 regardless of eruption magnitude or style. The fallout deposits documented across great  
427 distances change from bimodal to unimodal distributions in distal and off-axis regions  
428 (Figure 6).

429           The coarse and fine subpopulations retrieved by deconvolution of the bimodal PSDs  
430 display a consistent spatial pattern. The coarse subpopulation Md decreases with distance  
431 from vent, while the Md of the fine subpopulations is quasi-constant over the whole extent of  
432 the deposit (Figure 6A). Distributions are unimodal at distances where the coarse  
433 subpopulation Md converges towards the value of the fine subpopulation Md. The decay rate  
434 of the coarse subpopulation Md with distance from vent varies from one eruption to another,  
435 with the high magnitude eruptions (Plinian and ignimbrite-forming) showing a slower decay  
436 rate than low magnitude events (Vulcanian and Subplinian) (Figure 6). Only the fine  
437 subpopulation Md from the Vulcanian and Subplinian eruptions decrease slightly in size with  
438 distance from vent (Figure 6). The grain size of the fine subpopulation Md in bimodal  
439 deposits is similar to those of the co-PDC grain size distributions, ranging from 63 to 30  $\mu\text{m}$   
440 (Figure 6).

441

#### 442 *4.4.3. Origin of the bimodality in mixed fallout deposits*

443           The spatial variation of the grain size subpopulations in the bimodal distributions of  
444 mixed fallout deposits have been linked to co-PDC contribution in several studies using  
445 different lines of evidence (e.g. Di Muro et al. 2008; Eychenne et al. 2012; Sparks and Huang  
446 1980). The spatial variation of stratification observed in large magnitude eruptions (Engwell  
447 et al. 2014; Sparks et al. 1984; Wiesner et al. 2004) provides evidence for contemporary  
448 deposition from both co-PDC and vent-derived plumes. The coarse and fine proximal units  
449 are inferred to derive from a vent-derived Plinian and co-PDC plumes, respectively (Engwell  
450 et al. 2014; Sparks et al. 1984; Sparks and Huang 1980). At intermediate distances,  
451 contribution from co-PDC and vent-derived plumes cannot be distinguished stratigraphically,  
452 and a massive bimodal ash layer resulting from sedimentation of both plume types is evident  
453 (Figure 3; Engwell et al. 2014; Sparks et al. 1984; Sparks and Huang 1980). At greater

454 distances, the bimodality vanishes, and the ash deposit has a unimodal PSD, often skewed  
455 towards the fine grain sizes (Figure 3; Engwell et al. 2014). This depositional pattern suggests  
456 that mixing of the two source contributions produces bimodality in these fallout deposits to  
457 distances where the vent-derived plume is depleted in large particles, at which point the  
458 remaining vent-derived ash is of a similar size to the co-PDC ash (Engwell et al. 2014). This  
459 mechanism can also be invoked to explain the convergence of the coarse and fine  
460 subpopulation Md trends towards the unimodal distributions (Figure 6).

461 Componentry changes were compared to the spatial variations of the proportion of  
462 fine and coarse subpopulations in the mixed fallout deposit from the 1980 Mount St. Helens  
463 Plinian eruption. Accounting for the complex syn-eruptive wind field, this comparison has  
464 shown that the consistent fine subpopulation was deposited by a co-PDC plume resulting  
465 from a blast-PDC, while the coarse subpopulation was deposited by a vent-derived plume  
466 formed later during the eruptive sequence (Eychenne et al. 2015).

467

## 468 **5. Controls on co-PDC plume formation**

469

470 The formation of co-PDC plumes is intimately related to processes occurring within  
471 the parent PDC. The consistent co-PDC grain size characteristics described above (Figure 6)  
472 suggests that processes controlling the entrainment of particles in co-PDC columns allow the  
473 preferential lofting of particles finer than about 90  $\mu\text{m}$ . A number of mechanisms have been  
474 identified as leading to the formation of co-PDC plumes. Fine particles are mobilised from  
475 the dense basal portion of the flows towards the turbulent upper portions, by escaping gas, in  
476 a process called elutriation. The source of this gas may be internal - continuing release of  
477 magmatic gas from hot juvenile material (Sparks 1978) - or external to the current - from  
478 incineration of vegetation, or formation of steam on interaction with ice (Woods and Kienle

479 1994) or water bodies (Moore and Rice 1984; Sigurdsson and Carey 1989). When gas  
480 velocities are greater than particle terminal velocity, these particles elutriate from the dense  
481 basal flow into the overriding turbulent suspension. As the hot current propagates beneath the  
482 effectively stagnant air, a turbulent shear boundary forms, aiding entrainment of hot particles  
483 from the main body of the flow as ambient air is forced over the top (Denlinger 1987).  
484 Finally, buoyant lift off of the whole current is proposed to occur during emplacement of  
485 large PDCs as sedimentation and ambient air entrainment lowers the density of the flow to  
486 less than that of the ambient (Sparks et al. 1986; Sparks et al. 1997).

487

#### 488 *5.1. Source conditions*

489 Detachment and rise of the upper turbulent portion of a propagating current occurs  
490 when its density becomes equal to that of the surrounding atmosphere, a process called  
491 buoyancy reversal. Gas mass fractions of 30 - 60 % are required for the mixture to reach an  
492 equivalent density to the atmosphere (Woods and Wohletz 1991). Such fractions are  
493 considerably greater than the gas mass fraction of the mixture emitted at vent (1 - 5 %;  
494 Sparks et al. 1986; Wilson et al. 1980). The density of the flow decreases by a combination of  
495 particle sedimentation, and the entrainment, heating and thermal expansion of ambient air.  
496 These processes have been thoroughly investigated using analogue experiments, whereby a  
497 dense sediment rich fluid is injected into a sloped or flat tank of less dense saline fluid (e.g.  
498 Carey et al. 1988; Huppert et al. 1986; Sparks et al. 1993; Woods and Bursik 1994; Woods  
499 and Caulfield 1992). These experiments highlight the role of entrainment in lowering current  
500 density. They also show that increasing slope enables greater entrainment, related to  
501 increased shear mixing and consequently buoyancy reversal sooner than on shallow slopes  
502 (Woods and Bursik 1994). However, thermal expansion of entrained fluid, a key process for  
503 the formation of co-PDC plumes, is not accounted for in these experiments. Experiments

504 where heated mixtures of talc (~ 22 microns in diameter) were fed into a tank to generate  
505 turbulent dilute mixtures highlight the importance of this process, and show that mass  
506 fractionation into co-PDC plumes scales with the temperature of the current (Andrews and  
507 Manga 2011; 2012).

508 Buoyancy reversal can also occur as PDCs interact with topography, resulting in the  
509 formation of discrete co-PDC plumes along the PDC runout as described in section 3, due to  
510 enhanced mixing locally and subsequent thermal expansion of air, causing strong fines  
511 depletion in PDC deposits (Branney and Kokelaar 2002; Druitt 1998; Sparks et al. 1997).  
512 Detailed field studies of blast-derived PDC deposits show that the current becomes  
513 progressively more dilute as it propagates, with increased sedimentation ahead of topographic  
514 highs (Druitt 1992; Hoblitt et al. 1981; Moore and Sisson 1981). Enhanced mixing and  
515 entrainment of ambient air also accompany hydraulic jumps, which occur as PDCs encounter  
516 a sudden change in slope causing an increase in flow thickness and a decrease in flow  
517 velocity (Branney and Kokelaar 2002; Freundt and Schmincke 1985; Levine and Keiffer  
518 1991).

519 In large eruptions, buoyancy reversal can affect the entire current resulting in buoyant  
520 lofting from the entire PDC extent to form giant plumes (Sparks et al. 1986; 1997). The May  
521 18, 1980 Mount St Helens blast-related co-PDC plume (section 3.2) is inferred to have  
522 formed by such buoyant lofting (Figure 2B; Moore and Rice 1984; Sparks et al. 1986).  
523 Observations of horizontal velocities of 100 m/s close to the final extent of the PDC (Moore  
524 and Rice 1984) imply high mass flux rates within the flow at the time of lofting and indicate  
525 rapid transition from horizontal to vertical momentum as a result of buoyancy reversal.

526

527            *5.2. Insights from numerical models for column dynamics*

528            Numerical modelling of column dynamics is essential to identify inputs required for  
529 atmospheric dispersal models. Numerical modelling of co-PDC plume formation and rise has  
530 typically involved the modification of models normally applied to vent-derived plumes (e.g.  
531 Woods and Wohletz 1991). Co-PDC plumes have been simulated using steady-state co-PDC  
532 models (Calder et al. 1997; Woods and Wohletz 1991), by buoyant thermal theory (Calder et  
533 al. 1997; Woods and Kienle 1994), and using multiple phase (e.g. gas and particulate)  
534 complex numerical models (Dobran and Neri 1993; Herzog and Graf 2010; Neri and Dobran  
535 1994) .

536            In one-dimensional steady state and buoyant models, the rising mixture is assumed to  
537 be just buoyant relative to the atmosphere, well mixed and in thermal equilibrium (Woods  
538 and Kienle 1994; Woods and Wohletz 1991). It is also generally hypothesized that the  
539 mixture has little to no initial momentum (vertical velocity in the range 0.1 – 10 m/s), based  
540 on observations of small volume discrete co-PDC plumes associated with intermediate size  
541 eruptions (Calder et al. 1997). Yet, as stated above, much greater velocities have been  
542 observed during large eruptions where buoyant lofting occurs from the whole PDC extent  
543 (Moore and Rice 1984).

544            The application of steady state models (Calder et al. 1997; Woods and Wohletz 1991)  
545 requires the assumption that the current provides a constant source of material over time  
546 scales longer than the ascent time of the co-PDC plume. Equations for mass, momentum and  
547 thermal conservation are solved with height taking both entrainment of ambient air and  
548 particle sedimentation into account. Characterisation of entrainment is key for accurate  
549 modelling of volcanic plumes, and is specified using an entrainment coefficient in one-  
550 dimensional models. Entrainment rates are particularly poorly defined for co-PDC plumes,  
551 where the large initial radius means that entrained air may not mix completely within the

552 column (Calder et al. 1997). Two entrainment assumptions have been applied when  
553 modelling co-PDC plumes: 1) constant entrainment with height (Woods and Wohletz 1991),  
554 and 2) variable entrainment with height (Calder et al. 1997) with reduced entrainment near  
555 the ground, allowing the small scale of turbulent eddies relative to plume radius to be taken  
556 into account (Calder et al. 1997; Sparks et al 1997). Comparison of results obtained using  
557 these different assumptions shows differences in velocity at the base of the plumes, but no  
558 significant difference in the final plume height (Sparks et al. 1997). With the steady state  
559 model (Woods and Wohletz 1991), strong contraction in plume radius is observed  
560 immediately after lofting (Figure 7). Such a phenomena has been observed in the formation  
561 of large co-PDC plumes (section 3.2 and Figure 2D). The contraction is related to the  
562 decrease in plume density relative to the ambient (Figure 7). As ambient air is entrained and  
563 expands, the plume becomes buoyant and accelerates upwards (Figure 7), resulting in an  
564 increase in mass flux per unit area with height which is faster than the rate of addition of  
565 mass by entrainment. As a consequence, the plume radius decreases (Figure 7) to conserve  
566 mass (Woods and Wohletz 1991).

567 A discrete thermal model (Woods and Kienle 1994) was proposed following  
568 observations of co-PDC plumes formed during dome collapse events. The model assumes hot  
569 material is elutriated on short timescales compared to the ascent time in the atmosphere, and  
570 rises as a discrete spherical hot cloud. The model predicts that the height of the co-PDC  
571 plume is dependent on the initial mass, temperature and velocity of the cloud. While trends in  
572 velocity and density with height are similar to the steady state example (Figure 7), results do  
573 not show the same radial contraction, and instead radius increases linearly with height (Figure  
574 7). In both models, the maximum plume height is calculated as the height at which the  
575 vertical velocity is zero. Application of the steady state and discrete thermal models in 1D to  
576 small volume co-PDC plumes from intermediate size eruptions show that in reality neither

577 model completely describes plume ascent dynamics, but that the solution lies somewhere  
578 between these end members (Calder et al. 1997).

579 A two-dimensional axisymmetric two-phase numerical model has also been used to  
580 investigate the processes associated with the formation and rise of co-PDC plumes (Dobran  
581 and Neri 1993; Neri and Dobran 1994). This model accounts for thermal disequilibrium  
582 between the particles and gas, and for particle-particle interaction. The production of co-PDC  
583 plumes appears to be dependent on the fine-grained particles in the mixture, which follow the  
584 gas phase as it separates from the current. Model results reproduce the formation of multiple  
585 buoyant clouds above propagating currents, resulting in a substantial decrease in the  
586 spreading rate of the parent PDCs. This was also noted in hot granular experiments (Andrews  
587 and Manga 2012), where plume liftoff is seen to inhibit further spreading of the parent  
588 current. The modelled co-PDC plumes ascend rapidly, entraining large volumes of air. Strong  
589 contraction in plume radius immediately after lofting is also observed (Dobran and Neri  
590 1993), as seen in results from the steady state model (Figure 7).

591 Finally, application of three-dimensional multiphase flow models (Herzog and Graf  
592 2010) to simulate co-PDC plume rise highlights the importance of multiple discrete source  
593 locations on controlling plume dynamics, resulting in reduced neutral buoyancy heights  
594 compared to vent-derived examples. A key result is the difference in predicted neutral and  
595 maximum heights of co-PDC plumes, which cannot be easily differentiated using one-  
596 dimensional models, and can be on the order of tens of kilometres (see section 3), an  
597 important consideration for applying dispersal models to co-PDC plumes.

598

## 599 **6. Dispersion and sedimentation of co-PDC ash**

600

601                    *6.1. Importance of enhanced sedimentation*

602                    The grain size compilation (Figure 6) reveals that the coarse subpopulation of mixed  
603 fallout deposits decreases in size with distance from vent. In comparison, co-PDC ash and the  
604 fine grain subpopulations from mixed fallout deposits (assumed mainly co-PDC in origin): 1)  
605 have a much finer median size (about 63  $\mu\text{m}$ ), 2) show no variation of grain size with  
606 distance, and more importantly 3) settle continuously from the source to distal locations. The  
607 coarse trend converges towards the fine trend with distance. The coarse trend is consistent  
608 with particle size fractionation in a laterally spreading volcanic plume due to sedimentation  
609 governed by individual particle settling in a turbulent to intermediate regime (Bonadonna et  
610 al. 1998; Bursik et al. 1992b; Durant et al. 2009; Fisher 1964; Sparks et al. 1992). The fine  
611 trend suggests that particles  $< 63 \mu\text{m}$  in size continuously exit the plume during dispersion  
612 and that only the absolute amount of grains released, and thus the sedimentation rate,  
613 decreases with distance from source (as evidenced by the overall thinning of the deposits;  
614 Figures 4 and 5). This indicates that the fine particles do not settle following Stokes law as  
615 expected for grains of low Reynolds number whereby frictional forces control particle  
616 terminal velocity (i.e. settling in a laminar regime; Bonadonna et al. 1998). Stokes law  
617 predicts that particles smaller than 100  $\mu\text{m}$  have a relatively long residence time in the  
618 atmosphere and only start settling several hours or days after eruption (Rose and Durant  
619 2009), meaning that deposition would not begin until hundreds of kilometres from source  
620 (due to the advection of the plume by winds). Stokes law also predicts the deposits should  
621 display a decreasing grain size with distance (Bonadonna et al. 1998). The grain size  
622 characteristics in Figure 6 suggest that coarse and fine particles, dominantly from vent-  
623 derived and co-PDC plumes respectively, settle at similar distances from source and possibly  
624 at the same time (but not necessarily from the same height, e.g. Di Muro et al. 2008) despite

625 their contrasting sizes; such a phenomena is not captured by settling theory for individual  
626 grains.

627 Different mechanisms of enhanced sedimentation can be invoked to explain  
628 sedimentation of fine particles closer to the vent than expected for individual particle settling,  
629 and local increase of sedimentation rates. Aggregates, formed by particle clustering or  
630 adhering to the surface of larger clasts due to wet or electrostatic forces (Brown et al. 2012;  
631 Rose and Durant 2011; Taddeucci et al. 2011; Van Eaton et al. 2012), are often observed in  
632 proximal fallout deposits. They are assumed to form in vent-derived columns and plume  
633 corners where particle interaction is more intense due to high particle concentrations (Brown  
634 et al. 2012; Van Eaton and Wilson 2013). Observations of accretionary pellets in proximal  
635 co-PDC deposits indicate that aggregation is relevant for co-PDC columns too (Bonadonna et  
636 al. 2002a; Eychenne et al. 2012; Watanabe et al. 1999). Water is known to play a crucial role  
637 in aggregate formation (Brown et al. 2012), and co-PDC plumes are likely to entrain more  
638 water than vent-derived plumes due to the interaction of the PDCs with water bodies (lake,  
639 rivers, glaciers, ocean, etc., e.g. Hoblitt 2000) and vegetation. A number of other processes  
640 could potentially contribute to continuous sedimentation of fine particles over great distances.  
641 One such process relates to entrainment of fine particles in the wake of settling coarser grains  
642 (Eychenne et al. 2015; Rose et al. 2008; Di Muro et al. 2008). This is particularly relevant in  
643 the context of synchronous co-PDC and vent-derived plume dispersion at different altitudes.  
644 The passage of large particles through a suspension of fine particles is likely and could  
645 produce the consistent trend in fines with distance seen in Figure 6. Another model is linked  
646 to the development of layers of fine-grained material at the base of spreading plumes  
647 (Carazzo and Jellinek 2012; 2013; Manzella et al. 2015), which periodically collapse due to  
648 gravitational instabilities. The development of these layers is favoured by interaction of  
649 particles of different sizes and a high content of fine grains (Carazzo and Jellinek 2013), and

650 by growth of ash-bearing ice hydrometeors (Durant et al. 2009). Release of material driven  
651 by gravitational instabilities could explain the secondary maximum in thickness and mass  
652 observed in some co-PDC and mixed fallout deposits (Figures 4 and 5; Bonadonna et al.  
653 2002a; Durant et al. 2009; Eychenne et al. 2015; Watanabe et al. 1999). Given the high  
654 concentration of fine ash in co-PDC plumes, sedimentation by gravitational instabilities may  
655 also explain the settling of these grains close to source.

656

## 657 *6.2. Source parameters for numerical modelling of dispersion and sedimentation*

658 Observations and simulations of co-PDC plumes (sections 3 and 5) consistently  
659 highlight the complexities of co-PDC source conditions relative to vent-derived plumes, with  
660 implications for numerical modelling of co-PDC ash dispersal. The initial plume radius is  
661 much larger or commonly described by multiple source areas making estimation of mass flux  
662 into the spreading plume problematic. During eruptions with synchronous vent and co-PDC  
663 activity, multiple plumes may be present in the atmosphere at different heights at the same  
664 time (e.g. Di Muro et al. 2008; Engwell et al. 2014; Eychenne et al. 2012; Rose et al. 2008).  
665 Specific attention must be paid to the co-PDC particle release height, as wind speeds and  
666 directions can vary drastically at different atmospheric levels, and therefore have a significant  
667 impact on model results (e.g. Bonadonna et al. 2002a; 2002b). Although co-PDC deposits  
668 formed from different types of eruptions share similar grain size trends (Figure 6), source  
669 radius and number of contributing discrete clouds varies with eruption. Critically, the total  
670 grain size distribution of the particulate material in co-PDC plumes is much finer than in  
671 vent-derived plumes (Bonadonna and Houghton 2005).

672 Intrusion of large co-PDC plumes into the atmosphere at neutral buoyancy occurs as a  
673 radially spreading gravity current (Holasek et al. 1996; Sparks et al. 1986; Woods and Kienle  
674 1994). Buoyancy related processes can dominate plume dispersion for many hundreds of

675 kilometres from source and have a great effect on downwind dispersion and sedimentation  
676 patterns (Baines and Sparks 2005). Advection-diffusion processes only become dominant at  
677 distance from source, dependent on the mass flux of the intruding plume (Pouget et al. 2013),  
678 and the strength of the wind. Numerical models commonly utilised in simulating ash  
679 dispersion assume that a mixture of material with defined grain size characteristics is released  
680 into the atmosphere from a point or line source, and dispersed by advection-diffusion  
681 following the dominant wind direction (e.g. Beckett et al. 2014; Bonadonna 2006; Folch et al.  
682 2010). Such numerical applications have been applied to simulate dispersion from several co-  
683 PDC plumes (Bonadonna et al. 2002b; Costa et al. 2012; Folch et al. 2010; Johnston et al.  
684 2012; Matthews et al. 2012). In a small number of examples, gravity current sedimentation  
685 models have been applied to co-PDC plumes, and have proved successful in inferring  
686 eruptive timescales and reproduce grain size trends (Kandlbauer et al. 2013). While the  
687 advection-diffusion models do not accurately describe near source plume intrusion dynamics,  
688 the gravity current models do not describe distal dispersion well and a combination of both  
689 processes is required to accurately reproduce deposit trends (Pouget et al. 2013). Neither  
690 model type is capable of accounting for the complex sedimentation processes key for  
691 description of co-PDC ash deposition (see section 6.1), which has a strong impact on  
692 predicted dispersion (paths and concentrations).

693         Results from numerical models are typically compared with actual deposit trends  
694 (specifically the mass per unit area at individual locations) in order to infer the source  
695 parameters of past eruptions, for example height of dispersion and eruption duration. By  
696 grouping vent-derived with co-PDC eruptive phases, large uncertainties may be incorporated  
697 when estimating eruptive source parameters. Volumes of co-PDC deposits can be  
698 significantly greater than those from vent-derived plumes in large magnitude eruptions (Pyle  
699 et al. 2006), with volumes the same order of magnitude of the parent PDCs (e.g. Scarpati et

700 al. 2014; Sparks and Walker 1977). Such assumptions and simplifications in modelling co-  
701 PDC events may thus limit our understanding of dispersal processes significantly, not only  
702 producing inconsistencies in operational simulation of atmospheric ash dispersion, but also  
703 making interpretation of deposits in the volcanic record problematic.

704

### 705 *6.3. Co-PDC plume retrieval by satellite infra-red methods*

706 Despite the fine-grained nature of co-PDC ash, detecting dispersing co-PDC plumes  
707 using satellite ash retrieval methods remains challenging. Grain size data (Figure 6) shows  
708 that co-PDC plumes transport ash particles over great distances which are considerably  
709 coarser (median 63 – 15.6 microns) than the grain sizes accurately described by satellite  
710 retrieval methods (particles finer than 15  $\mu\text{m}$ , Guéhenneux et al. 2015; Prata 1989; Stevenson  
711 et al. 2015). Furthermore, co-PDC plumes may be richer in water compared to some vent-  
712 derived ash clouds (Darteville et al. 2002), especially if interaction with ice, or water bodies  
713 has occurred (Scott and McGimsey 1994; Woods and Kienle 1994). As the plume rises, such  
714 water condenses and generates ice hydrometeors (Durant and Rose 2009; Durant et al. 2009;  
715 Van Eaton et al. 2012), complicating retrieval signatures, specifically in the use of 2-Band  
716 methods (e.g. Rose 2000). This results in difficulties distinguishing volcanic and  
717 meteorological clouds (Guéhenneux et al. 2015; Simpson 2000), and can lead to  
718 underestimates in plume size, as observed for the co-PDC plume associated with the  
719 November 6, 1997 dome collapse event on Montserrat (Bonadonna et al. 2002b).

720

## 721 **7. Conclusions**

722 Observations of co-PDC plume formation, along with the characterisation of their  
723 deposits, and insights from analogue and numerical modelling, highlight co-PDC plumes as  
724 an exceptionally efficient mechanism for transporting fine ash into the atmosphere. In

725 intermediate size eruptions, co-PDC and vent-derived plumes are often coeval (with the  
726 exception of dome collapse events), and co-PDC deposits generally occur mixed with tephra  
727 from vent-derived plumes, producing fallout deposits with bimodal grain sizes. Comparison  
728 of grain size trends from a number of co-PDC events of varying scale show remarkable  
729 similar characteristics, despite different eruption styles and magnitudes. Specifically, ash  
730 transported in co-PDC plumes is very fine (< 90 microns) and dispersed over widespread  
731 areas without noticeable change in grain size. The small variation in grain size with distance,  
732 suggest that enhanced sedimentation associated with the formation of aggregates or  
733 gravitational instabilities dominates in co-PDC plumes, rather than individual, size-  
734 dependent, particle settling. When both form during explosive eruptions, vent-derived and co-  
735 PDC plumes can disperse synchronously within the atmosphere, complicating sedimentation  
736 behaviour.

737 Observations of co-PDC plume formation and results from analogue models highlight the  
738 importance of topographic control (e.g. slope changes, barriers, bends and constrictions in  
739 channels) on the formation of multiple discrete co-PDC plume pulses, and the complexities of  
740 source conditions. Multiple source regions are common, and the initial radius may be many  
741 tens of kilometres, many orders of magnitudes greater than for vent-derived plumes.  
742 Differences between vent-derived and co-PDC plume phenomena provide a cautionary note  
743 for the research community modelling and monitoring plume dispersion, and more broadly  
744 for assessing volcanic ash hazard. Co-PDC plumes need to be accounted for when assessing  
745 hazard during explosive eruptions producing PDCs, and tailored source conditions are  
746 required when numerically modelling co-PDC plume dispersion and sedimentation.

747

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758

759

760 **Captions**

761 **Table 1.** Examples and details of PDC-forming eruptions with associated co-PDC plume  
762 deposits for which grain size data exists in the literature. PDC Type: C – channelized deposit,  
763 B – deposit blanketing topography. Methods of grain size analysis used in each example  
764 (described in section 4.4.1): S – sieving, LD – laser diffraction, EC – electrozone counter, SR  
765 – settling rate. \*Data from LaMeve database (Crosweller et al. 2012) and calculated  
766 according to Pyle (1995). \*\* Polymodal grain size distributions deconvolved in this study.  
767 <sup>+</sup>Converted volume bulk to Dense Rock Equivalent (DRE, i.e. no porosity) assuming a  
768 deposit density of 1000 kg/m<sup>3</sup>, <sup>++</sup>Converted from mass to volume DRE assuming a solid  
769 density of 2500 (Mount St. Helens), 2600 (Tungurahua; Eychenne et al. 2012, and  
770 Montserrat; Bonadonna et al. 2002a) and 2200 kg/m<sup>3</sup> (Unzen; Watanabe et al. 1999).  
771 <sup>(1)</sup>Watanabe et al. (1999), <sup>(2)</sup>Miyabuchi (1999), <sup>(3)</sup>Bonadonna et al. (2002a), <sup>(4)</sup>Druitt et al.  
772 (2002), <sup>(5)</sup>Eychenne et al. (2013), <sup>(6)</sup>Bernard et al. (2014), <sup>(7)</sup>Eychenne et al. (2012), <sup>(8)</sup>Rose et  
773 al. (2008), <sup>(9)</sup>Davies et al. (1978), <sup>(10)</sup>Eychenne et al. 2015, <sup>(11)</sup>Sarna-Wojcicki et al. (1981),  
774 <sup>(12)</sup>Hoblitt et al. (1981), <sup>(13)</sup>Durant et al. (2009), <sup>(14)</sup>Gardner et al. (1998), <sup>(15)</sup>Beget (1983),  
775 <sup>(16)</sup>Brazier et al. (1983), <sup>(17)</sup>Carey and Sigurdsson (1986), <sup>(18)</sup>Macias et al. (1998), <sup>(19)</sup>Rose and  
776 Durant (2009), <sup>(20)</sup>Carey et al. (1995), <sup>(21)</sup>Wiesner et al. (2004), <sup>(22)</sup>Scott et al. (1996), <sup>(23)</sup>  
777 Sigurdsson et al. (1990), <sup>(24)</sup>Sparks and Huang (1980), <sup>(25)</sup>Williams (1942), <sup>(26)</sup>Williams and  
778 Goles (1968), <sup>(27)</sup>Bacon (1983), <sup>(28)</sup>Kandlbauer and Sparks (2014), <sup>(29)</sup>Sigurdsson and Carey  
779 (1989), <sup>(30)</sup>Pyle et al. (2006), <sup>(31)</sup>Fisher et al. (1993), <sup>(32)</sup>Engwell et al. (2014), <sup>(33)</sup>Rose and  
780 Chesner (1987), <sup>(34)</sup>Gatti and Oppenheimer (2012), <sup>(35)</sup>Knight and Walker (1986).

781

782 **Figure 1.** Comparison of the main characteristics of (A) vent-derived versus (B) co-PDC  
783 plumes. A. Vent-derived plumes are emitted from a point source corresponding to the top of  
784 the conduit and can be separated into three regions. The gas thrust region corresponds to the

785 first few hundred metres above the vent where the plume can be described as a jet due to its  
786 very high velocities, and its high density in comparison to the ambient. As the jet rises,  
787 ambient air is entrained, is heated and expands resulting in a density reduction relative to the  
788 atmosphere and the mixture rises convectively in the convective region. The final region, the  
789 umbrella, describes the height at which the plume spreads laterally into the atmosphere, and  
790 is the height at which dispersion occurs. **B.** Co-PDC plumes form as a propagating PDC  
791 sediments particles and entrains ambient air resulting in buoyancy reversal, whereby the PDC  
792 density become less than that of the ambient. The buoyant mixture detaches from the dense  
793 portion of the current, and entrains more air as it rises. Because initial vertical velocities are  
794 small and the mixture density is equivalent to the ambient, co-PDC plumes do not have a jet  
795 component, but are solely convective. In the same manner as vent-derived plumes, once the  
796 plume reaches neutral buoyancy, it spreads laterally.

797

798 **Figure 2.** Examples of co-PDC plumes. **A.** Co-PDC plume rising from a channel-confined  
799 PDC during the Plinian eruption of Pinatubo, Indonesia, on June 15, 1991 (photo by R.  
800 Hoblitt; PDC in the Sacobia river valley 20 km north of Pinatubo at 6:01 local time). **B.**  
801 Blast-related co-PDC plume during the Plinian eruption of Mount St. Helens, U.S., on May  
802 18, 1980 (photo by R. Kolberg taken at ~ 15:54 GMT looking from the north-west; in Sparks  
803 et al. 1986). **C.** Co-PDC plume generated by a dome collapse on November 4, 1997 at  
804 Soufriere Hills, Montserrat (photo by E. Calder in Bonadonna et al. 2002a). **D.** Co-PDC  
805 plume generated by a dome collapse on April 21, 1990 at Redoubt volcano, Alaska (photo by  
806 R.J. Clucas). The red arrows indicate the main PDC propagation direction. The thick black  
807 arrows represent the direction and level of plume spreading. The thin black arrows indicate  
808 discrete uplift.

809

810 **Figure 3.** Grain size characteristics of the fallout deposit from the 39 ka Campanian  
811 ignimbrite-forming eruption from across the Mediterranean, showing the source location of  
812 the eruption (star), dispersal extent (dashed line, from Pyle et al. 2006), and locations of  
813 displayed grain size information. **A.** Grain size distributions of the basal vent-derived tephra  
814 layer and the upper co-PDC ash layer in a lake sediment core (photograph courtesy of Jens  
815 Mingram) 120 km from source. **B, C, and D.** Grain size changes with increasing distance  
816 from source within the unstratified Campanian ash layer observed in deep sea cores beyond  
817 400 km from source. The red dashed line represents the size threshold beyond which particles  
818 are considered as fine ash (see Cashman and Rust chapter).

819

820 **Figure 4.** Dispersal of co-PDC deposits from dome collapse events at Soufriere Hills,  
821 Montserrat (**A** and **B**) and Mount Unzen, Japan (**C** and **D**), showing the controls of the  
822 geometry of the PDC deposits and the wind direction on the geometry and extent of the co-  
823 PDC deposits. **A** and **B.** Isopach maps (in mm) reproduced from Bonadonna et al. (2002a). **C**  
824 and **D.** Isomass maps (in  $\text{g}/\text{m}^2$ ) reproduced from Watanabe et al. (1999). tr: trace of ash-fall.

825

826 **Figure 5.** Dispersal of mixed fallout deposits from the May 18, 1980 Mount St. Helens (**A**)  
827 and August 16, 2006 Tungurahua (**B**) eruptions showing the effect of the co-PDC  
828 contributions on the geometry of the mixed deposits. Isomass lines from (**A**) Sarna-Wojcicki  
829 et al. (1981) and (**B**) Eychenne et al. (2013). **A.** Note the displacement of the overall deposit  
830 to the north of the vent (blue arrow) due to the northern momentum of the co-PDC plume  
831 generated by the blast. **B.** Note the steep high in topography (about 500 m to 1 km high)  
832 located perpendicular to the PDC propagation direction which enhanced the formation of co-  
833 PDC plumes.

834

835 **Figure 6.** Grain size variations of co-PDC and mixed fallout deposits with distance from  
836 vent. The Median Diameter of the unimodal particle size distributions and coarse and fine  
837 grain size subpopulations of mixed fallout deposits are distinguished. See Table 1 for the list  
838 of eruptions and references.  
839

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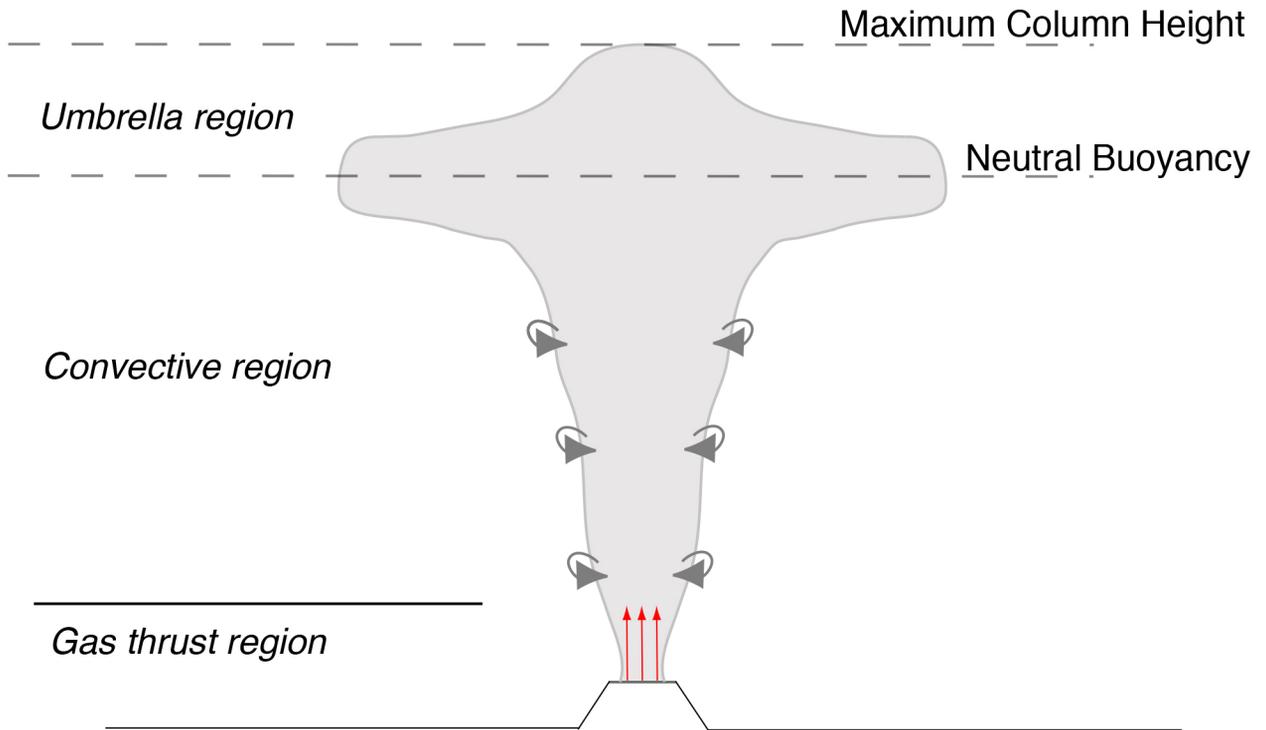
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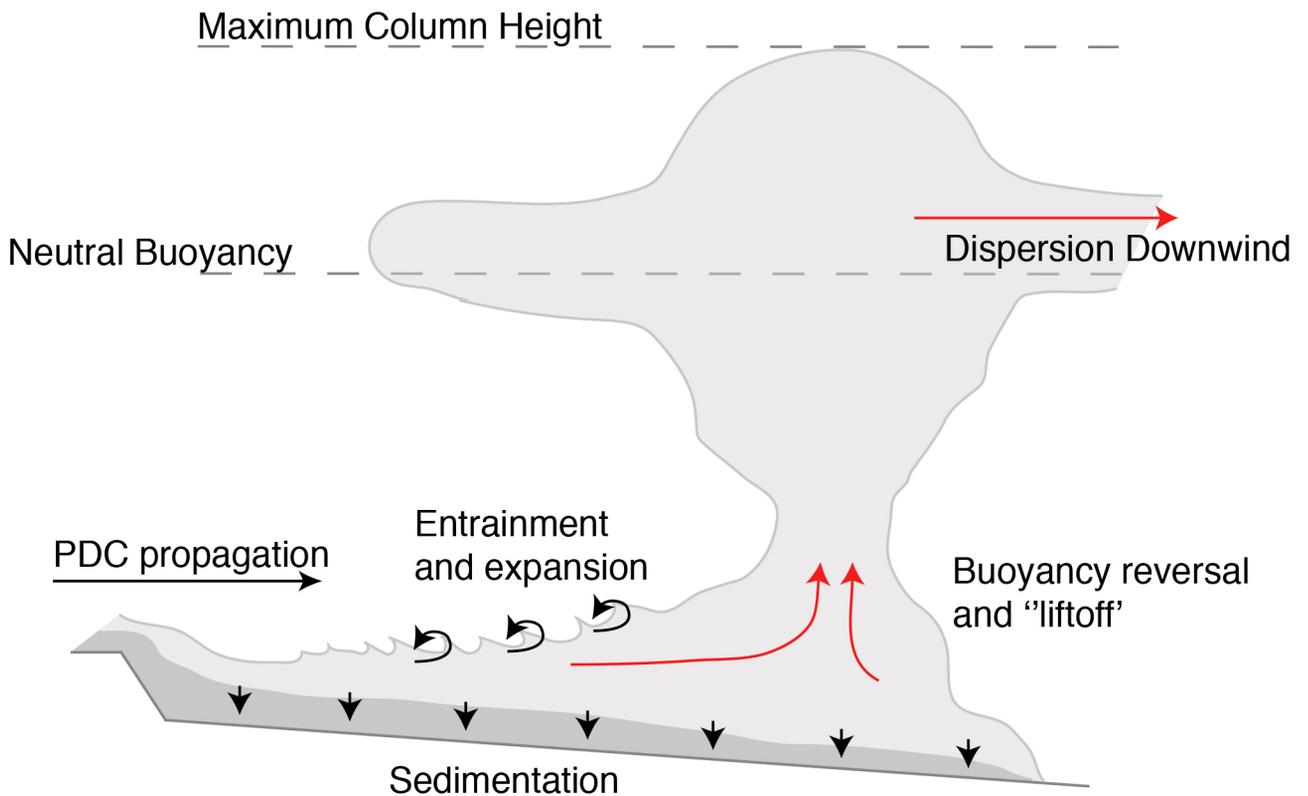
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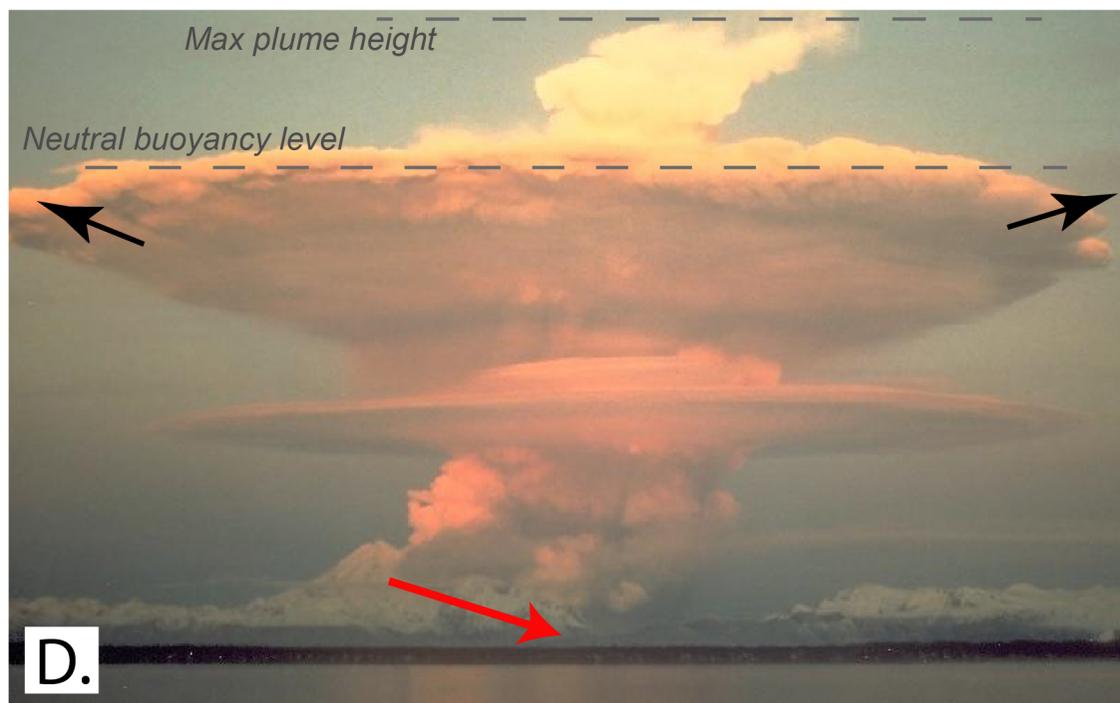
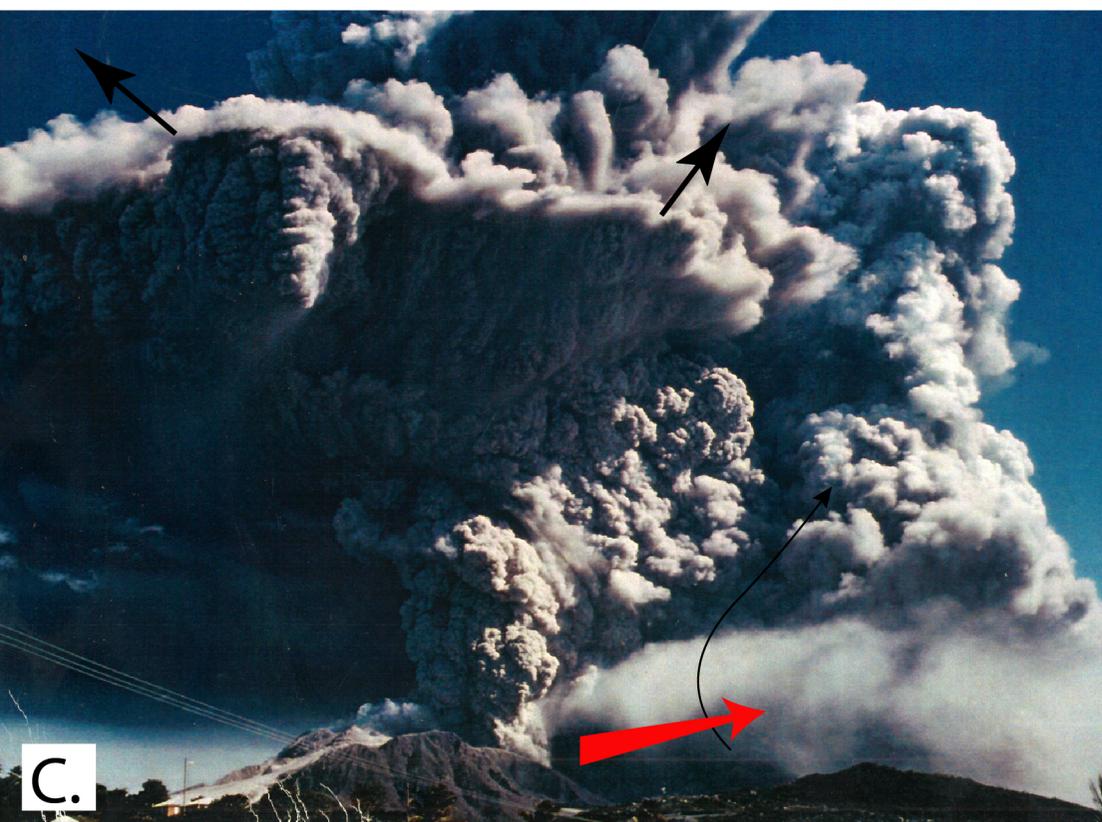
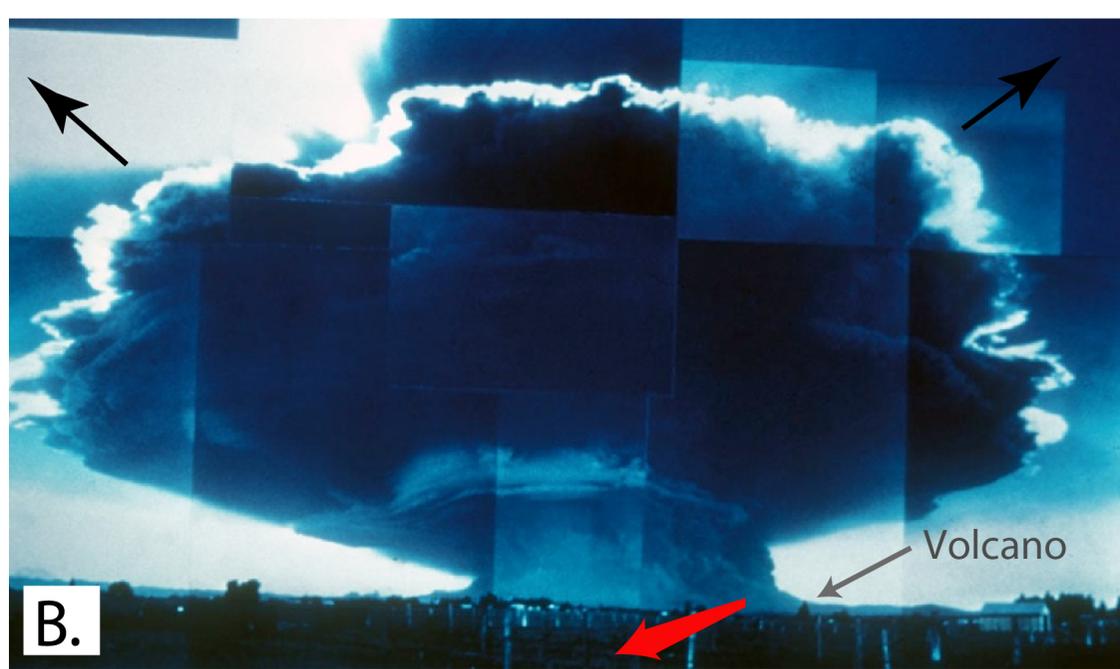
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## A: Vent-derived plume

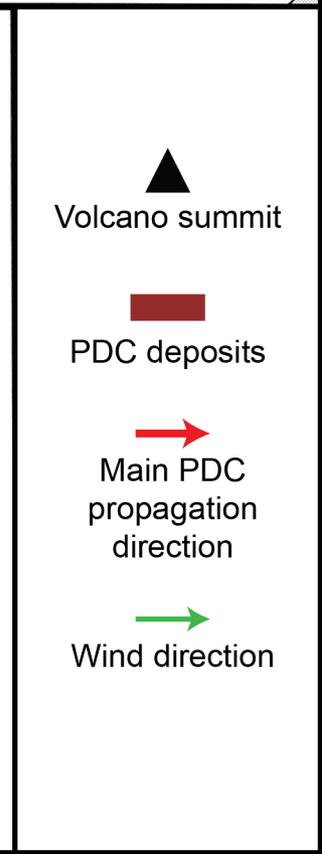
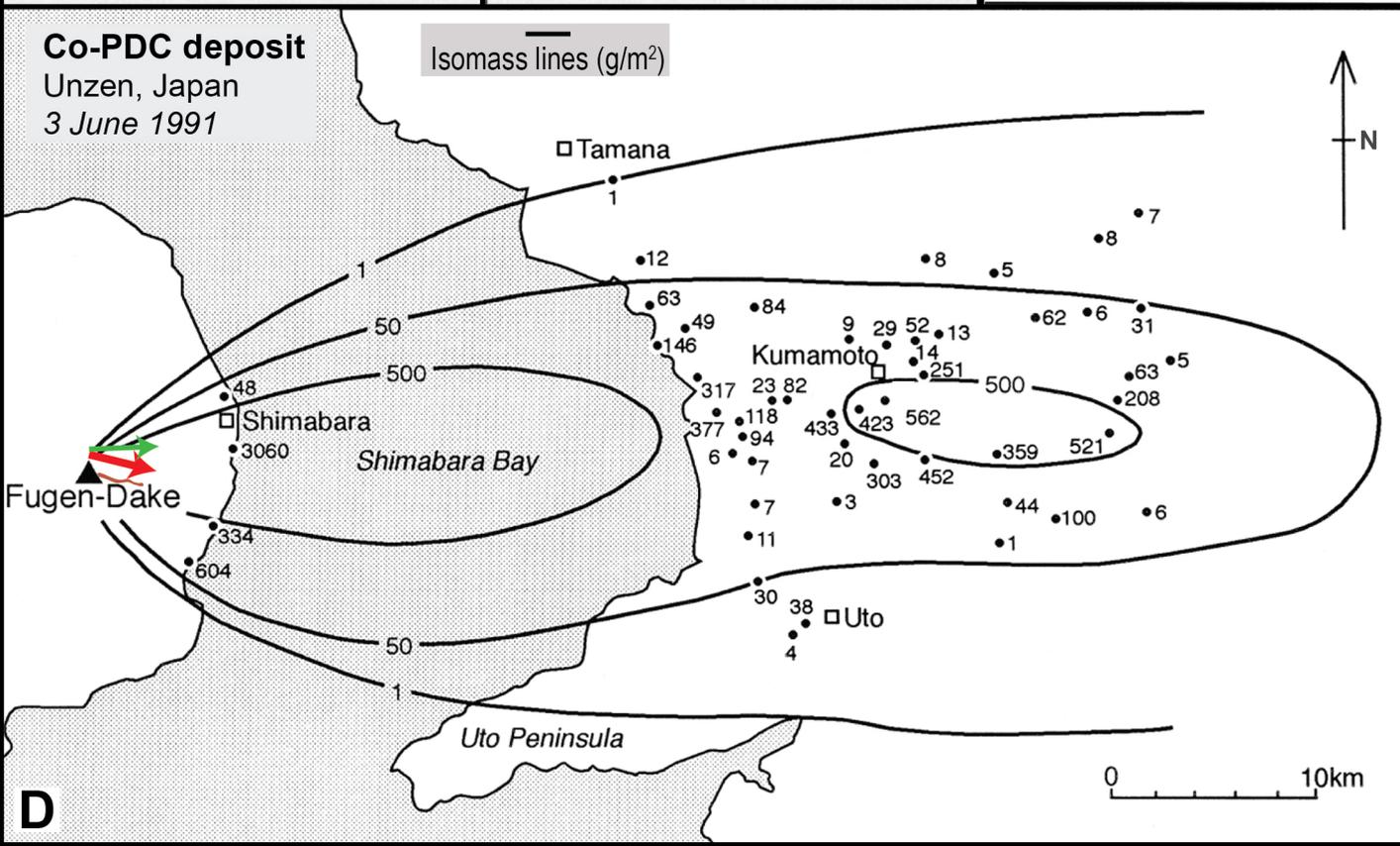
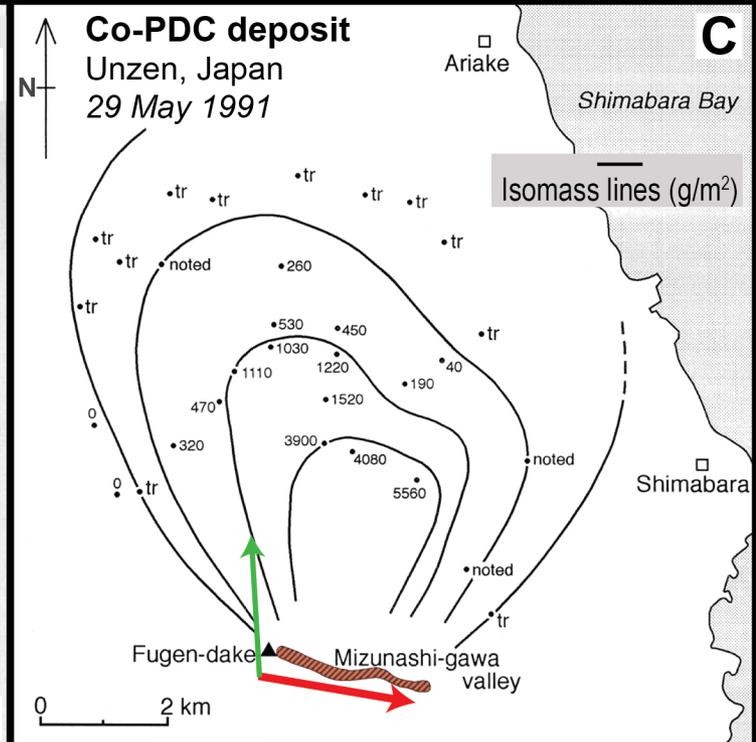
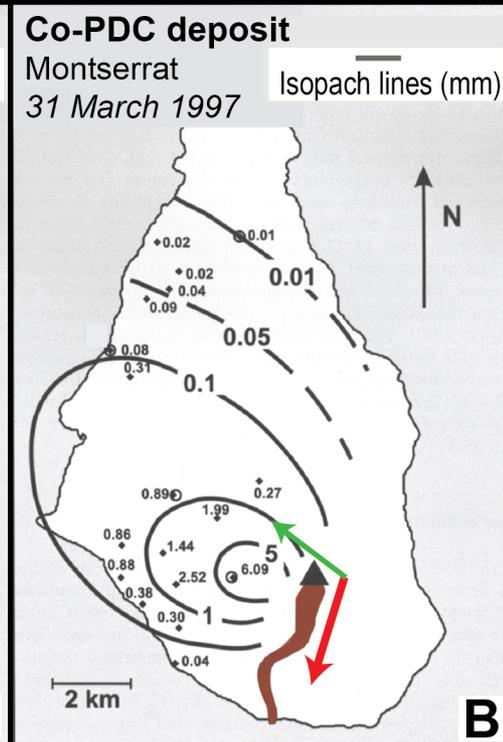
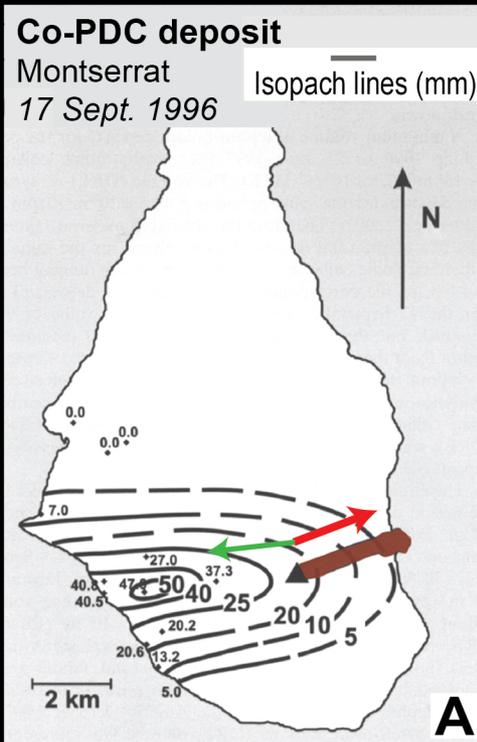


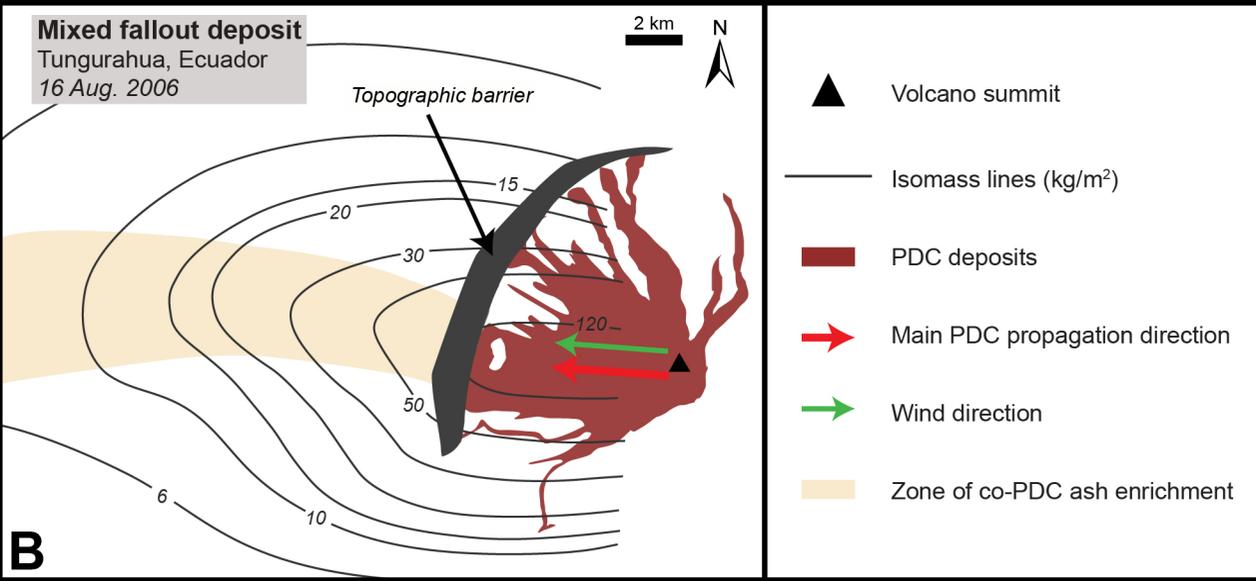
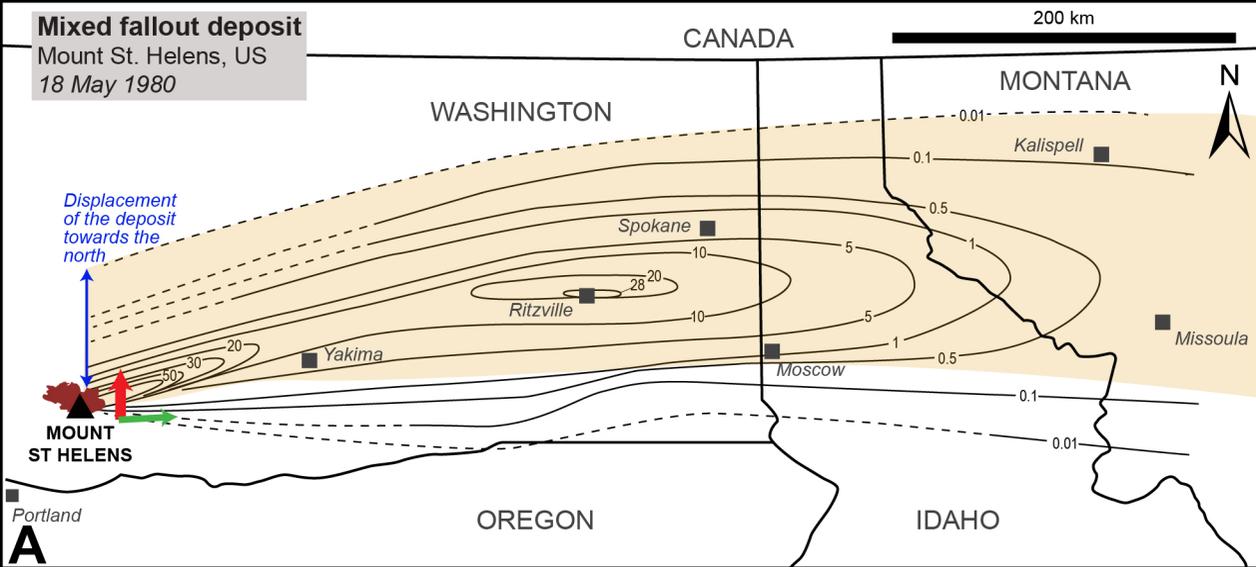
## B: Co-PDC plume











× Co-PDC deposits

Mixed fallout deposits:

○ Coarse peak from bimodal distr.

● Fine peak from bimodal distr.

▲ Unimodal distr.

Unzen

Montserrat

Dome collapse

Montserrat

Tungurahua

Fuego

Vulcanian -  
Subplinian

MSH 1980

Glacier Peak

MSH Yn

Pinatubo

Plinian

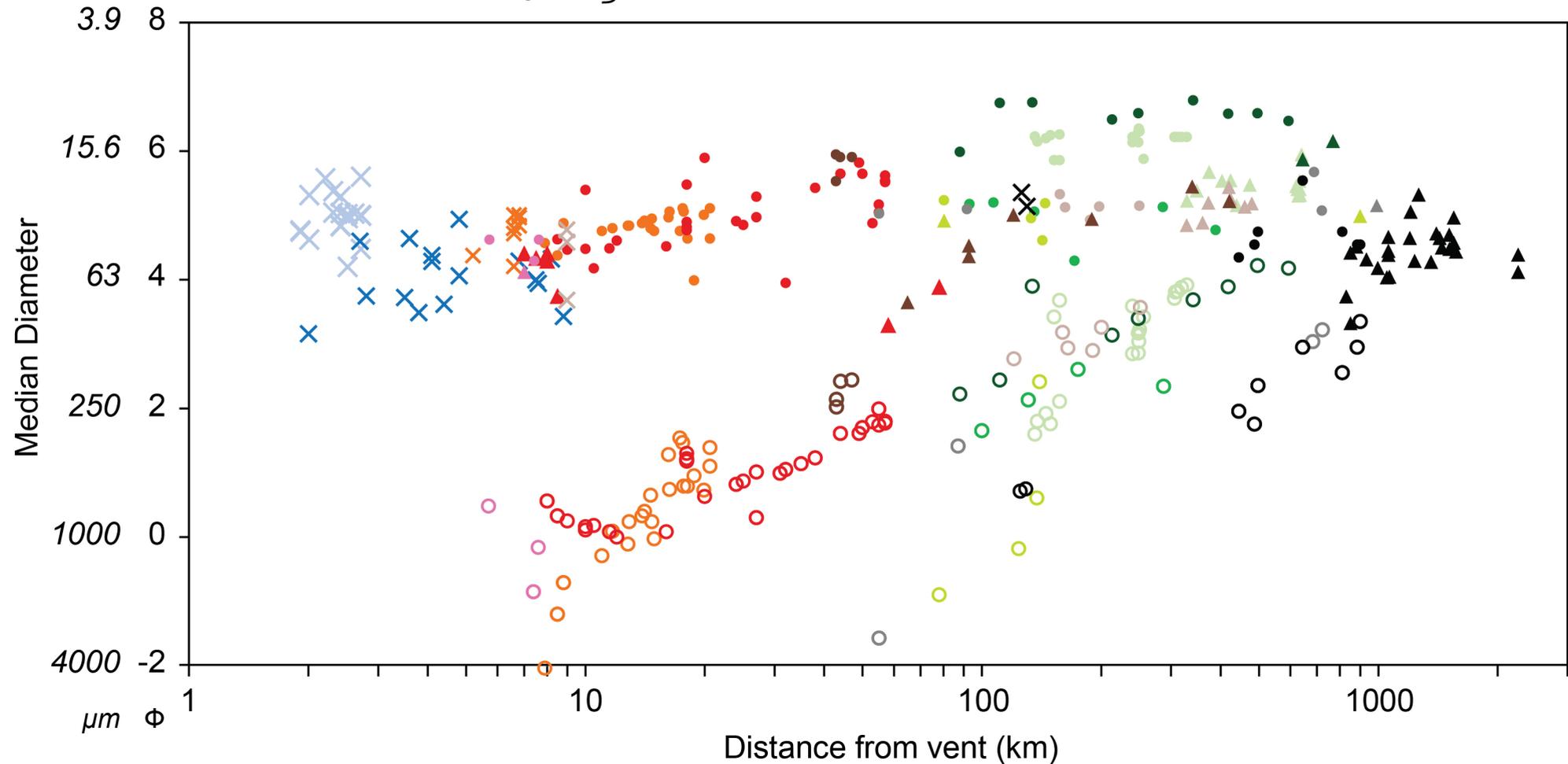
Minoan

Mazama

Tambora

Campanian

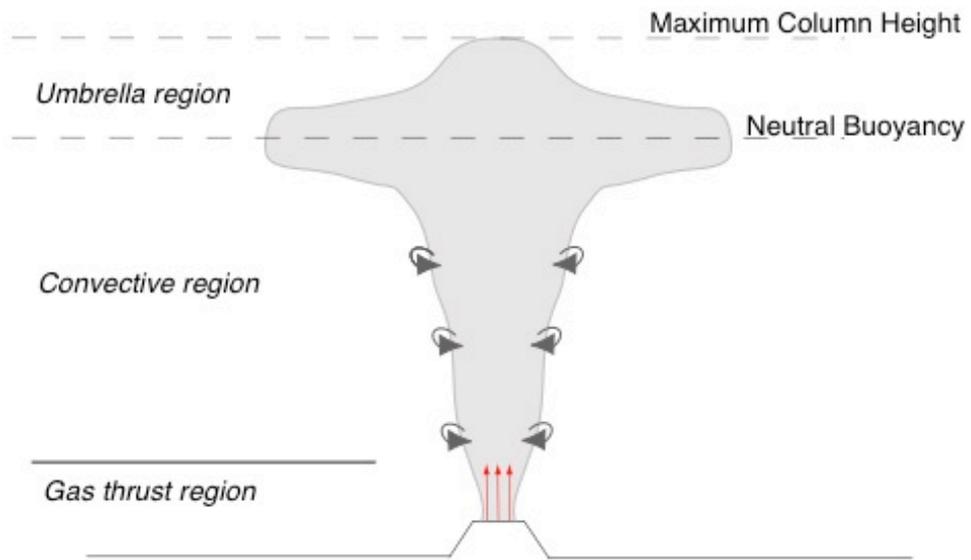
Ignimbrite-forming



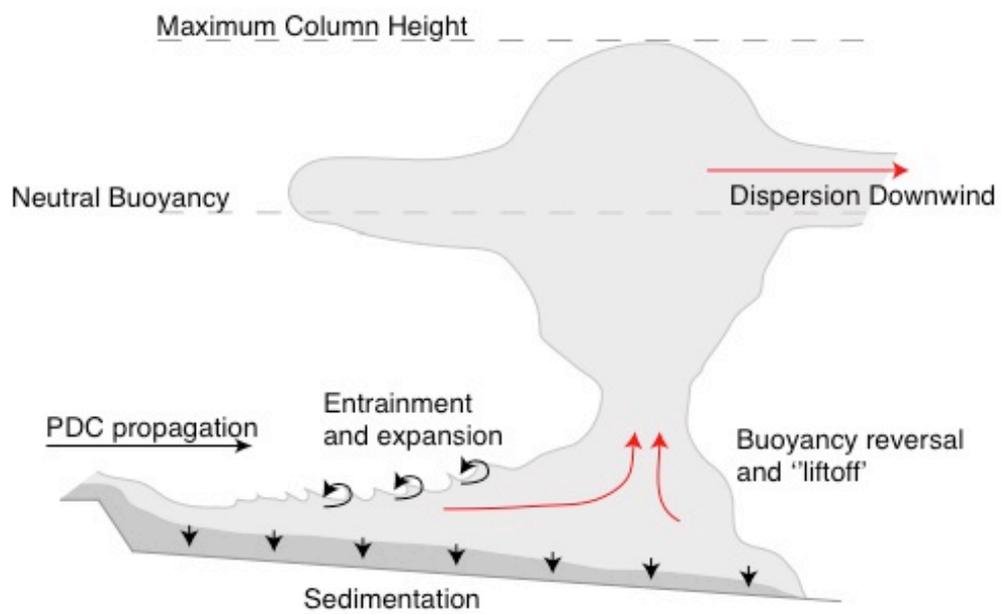


**Table 1.** Examples and details of PDC-forming eruptions with associated co-PDC plume deposits for which grain size data exists in the literature. PDC Type: C – channelized deposit, B – deposit blanketing topography. Methods of grain size analysis used in each example (described in section 4.4.1): S – sieving, LD – laser diffraction, EC – electrozone counter, SR – settling rate. \*Data from LaMeve database (Croweller et al. 2012) and calculated according to Pyle (1995). \*\* Polymodal grain size distributions deconvolved in this study. <sup>+</sup>Converted volume bulk to Dense Rock Equivalent (DRE, i.e. no porosity) assuming a deposit density of 1000 kg/m<sup>3</sup>, <sup>++</sup>Converted from mass to volume DRE assuming a solid density of 2500 (Mount St. Helens), 2600 (Tungurahua; Eychenne et al. 2012, and Montserrat; Bonadonna et al. 2002a) and 2200 kg/m<sup>3</sup> (Unzen; Watanabe et al. 1999). <sup>(1)</sup>Watanabe et al. (1999), <sup>(2)</sup>Miyabuchi (1999), <sup>(3)</sup>Bonadonna et al. (2002a), <sup>(4)</sup>Druitt et al. (2002), <sup>(5)</sup>Eychenne et al. (2013), <sup>(6)</sup>Bernard et al. (2014), <sup>(7)</sup>Eychenne et al. (2012), <sup>(8)</sup>Rose et al. (2008), <sup>(9)</sup>Davies et al. (1978), <sup>(10)</sup>Eychenne et al. 2015, <sup>(11)</sup>Sarna-Wojcicki et al. (1981), <sup>(12)</sup>Hoblitt et al. (1981), <sup>(13)</sup>Durant et al. (2009), <sup>(14)</sup>Gardner et al. (1998), <sup>(15)</sup>Beget (1983), <sup>(16)</sup>Brazier et al. (1983), <sup>(17)</sup>Carey and Sigurdsson (1986), <sup>(18)</sup>Macias et al. (1998), <sup>(19)</sup>Rose and Durant (2009), <sup>(20)</sup>Carey et al. (1995), <sup>(21)</sup>Wiesner et al. (2004), <sup>(22)</sup>Scott et al. (1996), <sup>(23)</sup>Sigurdsson et al. (1990), <sup>(24)</sup>Sparks and Huang (1980), <sup>(25)</sup>Williams (1942), <sup>(26)</sup>Williams and Goles (1968), <sup>(27)</sup>Bacon (1983), <sup>(28)</sup>Kandlbauer and Sparks (2014), <sup>(29)</sup>Sigurdsson and Carey (1989), <sup>(30)</sup>Pyle et al. (2006), <sup>(31)</sup>Fisher et al. (1993), <sup>(32)</sup>Engwell et al. (2014), <sup>(33)</sup>Rose and Chesner (1987), <sup>(34)</sup>Gatti and Oppenheimer (2012), <sup>(35)</sup>Knight and Walker (1986).

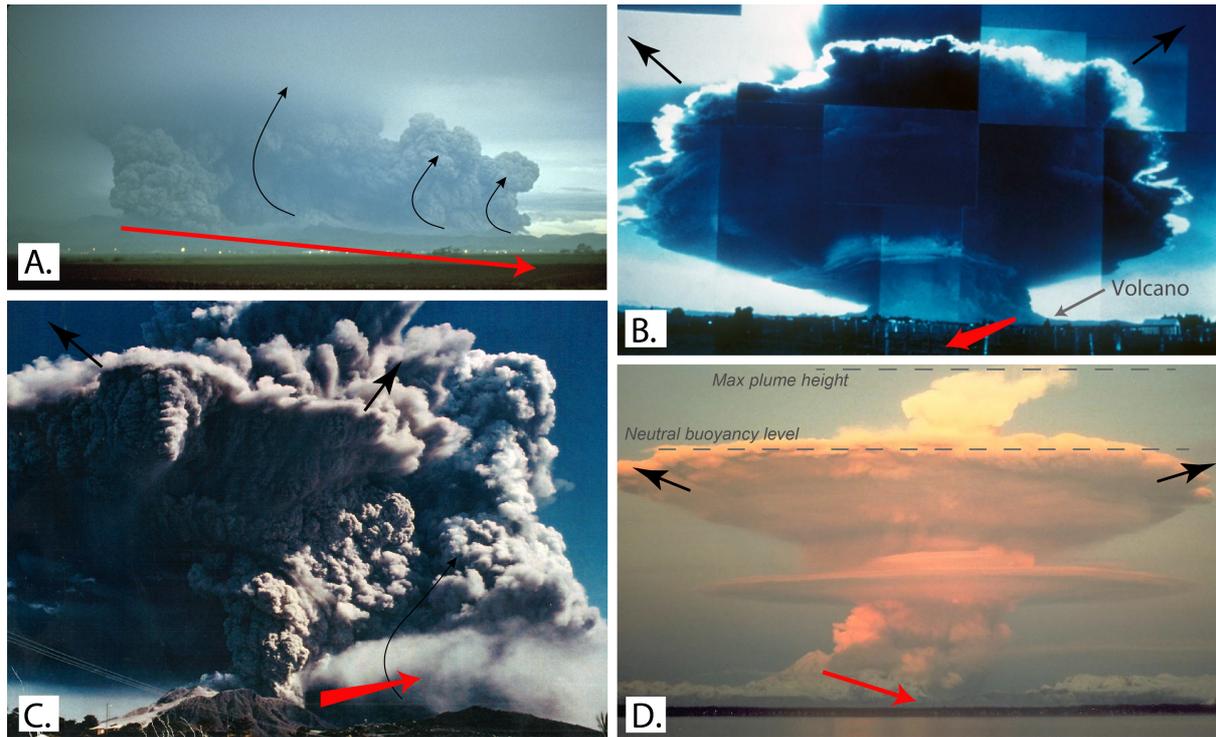
## A: Vent-derived plume



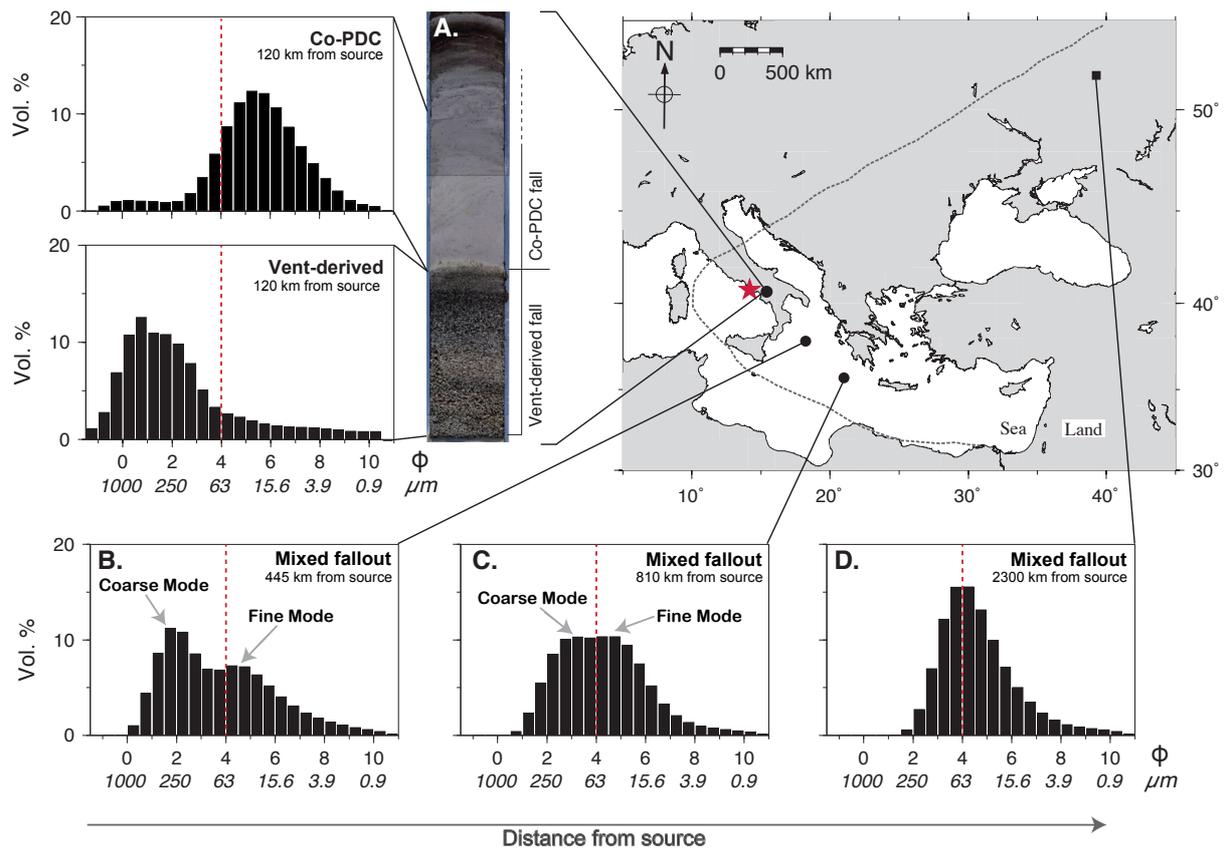
## B: Co-PDC plume



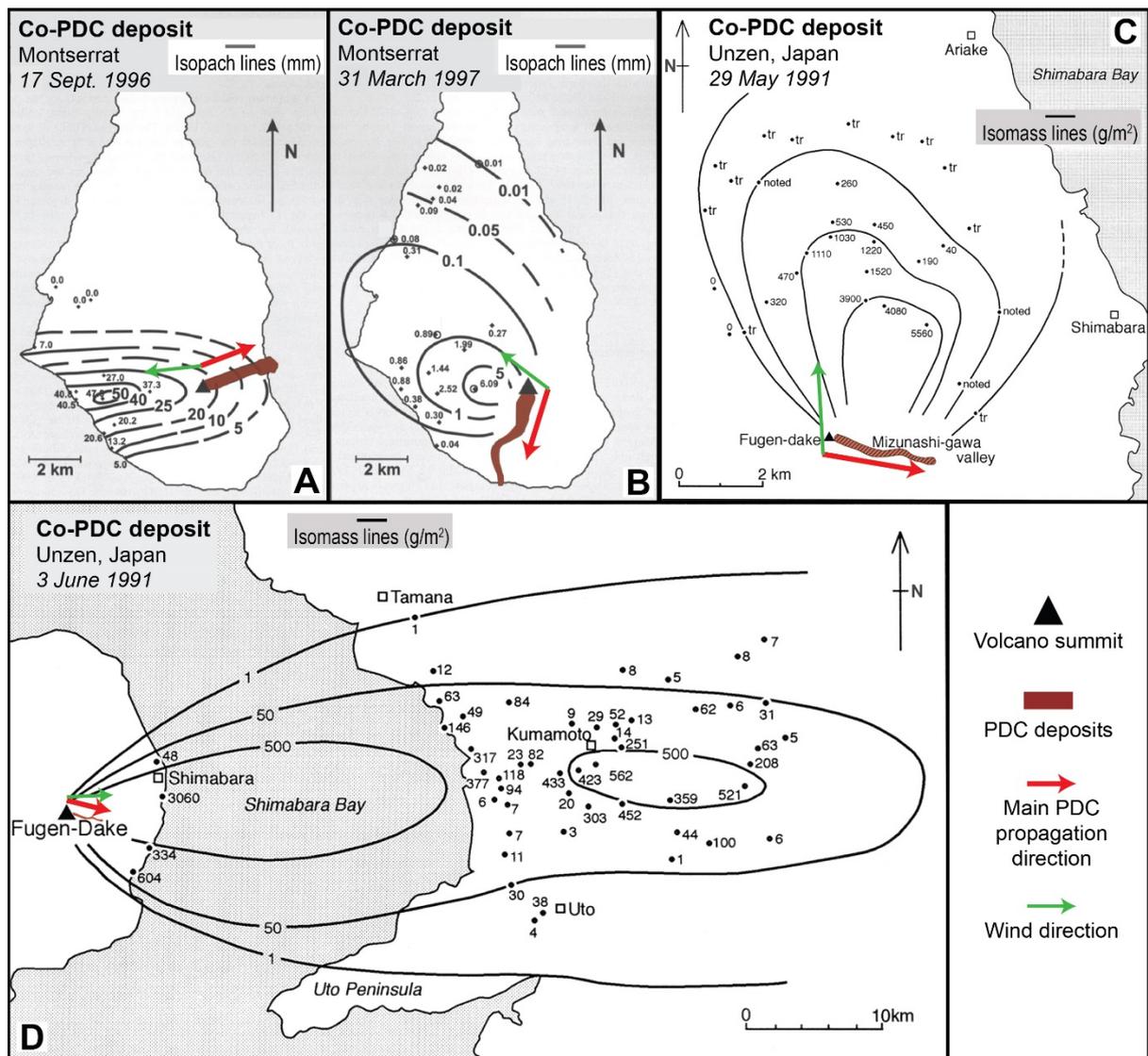
**Figure 1.** Comparison of the main characteristics of vent-derived **(A)** versus co-PDC plumes **(B)**. **A.** Vent-derived plumes are emitted from a point source corresponding to the top of the conduit and can be separated into three regions. The gas thrust region corresponds to the first few hundred metres above the vent where the plume can be described as a jet due to its very high velocities, and its high density in comparison to the ambient. As the jet rises, ambient air is entrained, is heated and expands resulting in a density reduction relative to the atmosphere and the mixture rises convectively in the convective region. The final region, the umbrella, describes the height at which the plume spreads laterally into the atmosphere, and is the height at which dispersion occurs. **B.** Co-PDC plumes form as a propagating PDC sediments particles and entrains ambient air resulting in buoyancy reversal, whereby the PDC density become less than that of the ambient. The buoyant mixture detaches from the dense portion of the current, and entrains more air as it rises. Because initial vertical velocities are small and the mixture density is equivalent to the ambient, co-PDC plumes do not have a jet component, but are solely convective. In the same manner as vent-derived plumes, once the plume reaches neutral buoyancy, it spreads laterally.



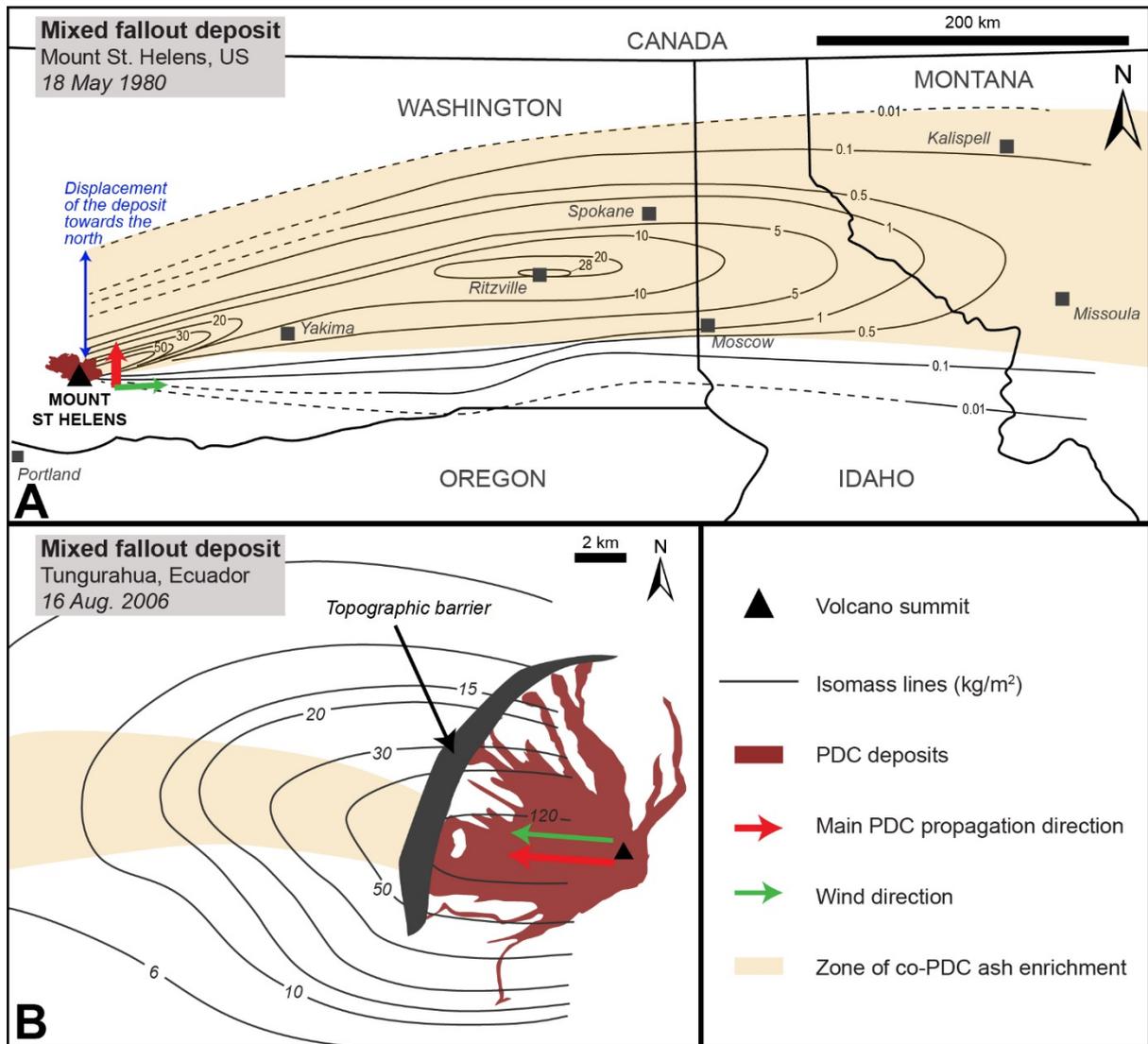
**Figure 2.** Examples of co-PDC plumes. **A.** Co-PDC plume rising from a channel-confined PDC during the Plinian eruption of Pinatubo, Indonesia, on June 15, 1991 (photo by R. Hoblitt; PDC in the Sacobia river valley 20 km north of Pinatubo at 6:01 local time). **B.** Blast-related co-PDC plume during the Plinian eruption of Mount St. Helens, U.S., on May 18, 1980 (photo by R. Kolberg taken at ~ 15:54 GMT looking from the north-west; in Sparks et al. 1986). **C.** Co-PDC plume generated by a dome collapse on November 4, 1997 at Soufriere Hills, Montserrat (photo by E. Calder in Bonadonna et al. 2002a). **D.** Co-PDC plume generated by a dome collapse on April 21, 1990 at Redoubt volcano, Alaska (photo by R.J. Clucas). The red arrows indicate the main PDC propagation direction. The thick black arrows represent the direction and level of plume spreading. The thin black arrows indicate discrete uplift.



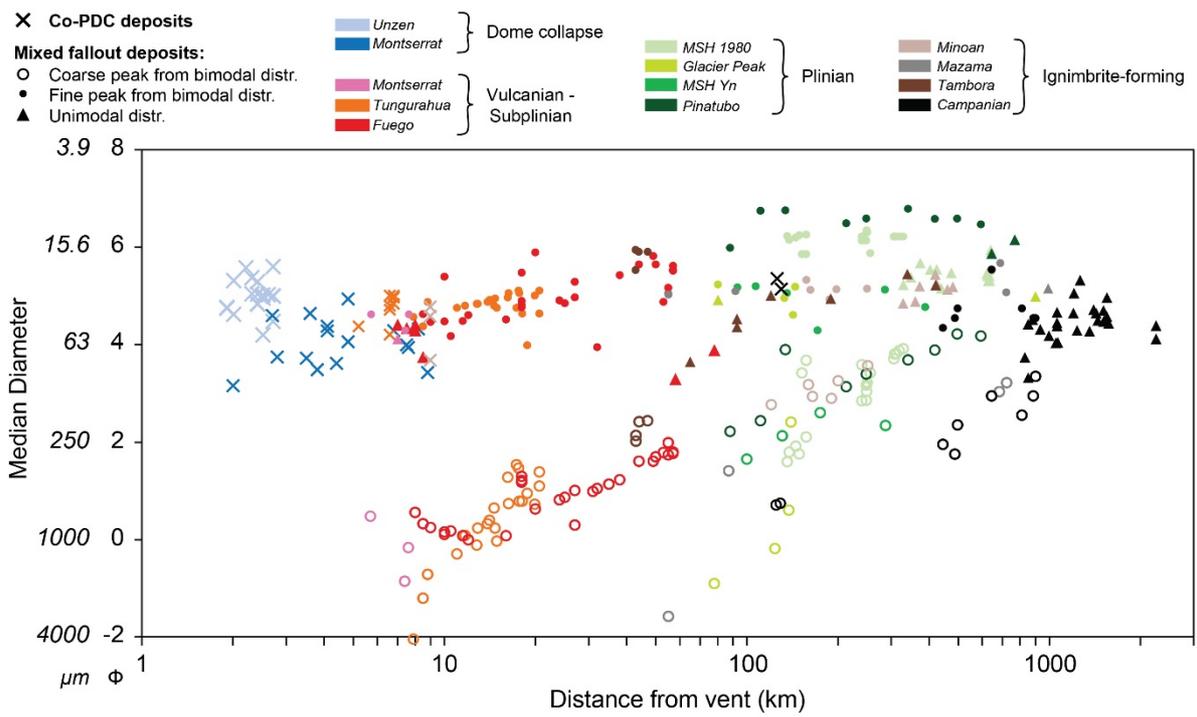
**Figure 3.** Grain size characteristics of the fallout deposit from the 39 ka Campanian ignimbrite-forming eruption from across the Mediterranean, showing the source location of the eruption (star), dispersal extent (dashed line, from Pyle et al. 2006), and locations of displayed grain size information. **A.** Grain size distributions of the basal vent-derived tephra layer and the upper co-PDC ash layer in a lake sediment core (photograph courtesy of Jens Mingram) 120 km from source. **B.** **C.** and **D.** Grain size changes with increasing distance from source within the unstratified Campanian ash layer observed in deep sea cores beyond 400 km from source. The red dashed line represents the size threshold beyond which particles are considered as fine ash (see Cashman and Rust chapter).



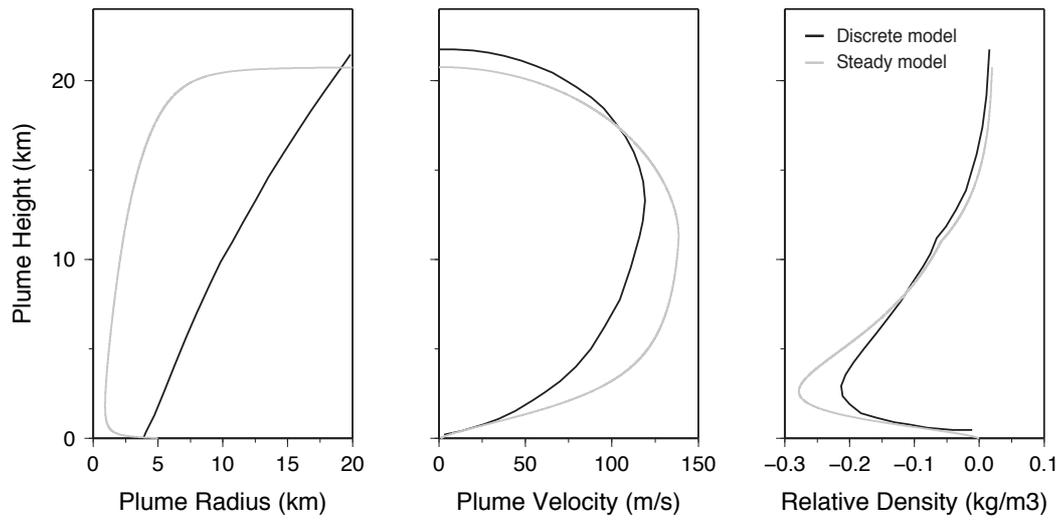
**Figure 4.** Dispersal of co-PDC deposits from dome collapse events at Soufriere Hills, Montserrat (**A** and **B**) and Mount Unzen, Japan (**C** and **D**), showing the controls of the geometry of the PDC deposits and the wind direction on the geometry and extent of the co-PDC deposits. **A** and **B**. Isopach maps (in mm) reproduced from Bonadonna et al. (2002a). **C** and **D**. Isomass maps (in g/m<sup>2</sup>) reproduced from Watanabe et al. (1999). tr: trace of ash-fall.



**Figure 5.** Dispersal of mixed fallout deposits from the May 18, 1980 Mount St. Helens (**A**) and August 16, 2006 Tungurahua (**B**) eruptions showing the effect of the co-PDC contributions on the geometry of the mixed deposits. Isomass lines from (**A**) Sarna-Wojcicki et al. (1981) and (**B**) Eychenne et al. (2013). **A.** Note the displacement of the overall deposit to the north of the vent (blue arrow) due to the northern momentum of the co-PDC plume generated by the blast. **B.** Note the steep high in topography (about 500 m to 1 km high) located perpendicular to the PDC propagation direction which enhanced the formation of co-PDC plumes.



**Figure 6.** Grain size variations of co-PDC and mixed fallout deposits with distance from vent. The Median Diameter of the unimodal particle size distributions and coarse and fine grain size subpopulations of mixed fallout deposits are distinguished. See Table 1 for the list of eruptions and references.



**Figure 7.** Comparison of model trends from the steady state plume model of Woods and Wohletz (1991), assuming an initial radius of 5 km, temperature of 832 K, an initial velocity of 1 m/s and initial plume density equal to that of the ambient, with the discrete buoyant thermal model of Woods and Kienle (1994), assuming an initial radius of 4 km and temperature of 700 K.

**Table 1.** Example of eruptions with

\* Magnitude data taken from LaMe  
+ Converted volume bulk to DRE a  
++ Converted from mass to volume |

1	Watanabe et al, 1999
2	Miyabuchi, 1999
3	Bonadonna et al, 2002
4	Druitt et al, 2002
5	Eychenne et al, 2013
6	Bernard et al, 2014
7	Eychenne et al, 2012
8	Rose et al, 2008
9	Davies et al. 1978
10	Eychenne et al, submitted
11	Sarna-Wojcicki et al, 1981;
12	Hoblitt et al, 1981
13	Durant et al, 2009
14	Gardner et al. 1998
15	Beget 1983
16	Brazier et al. 1983
17	Carey & Sigurdsson, 1986
18	Macias et al, 1998
19	Rose & Durant, 2008
20	Carey et al. 1995
21	Wiesner et al, 2004
22	Scott et al, 1996
23	Sigurdsson et al 1990
24	Sparks & Huang 1980
25	Williams 1942
26	Williams and Goles 1968
27	Bacon 1983
28	Kandlbauer et al. 2014
29	Sigurdsson & Carey 1989
30	Pyle et al. 2006
31	Fisher et al. 1993
32	Engwell et al. 2014
33	Rose & Chesner 1987
34	Gatti and Oppenheimer 2012
35	Knight et al. 1986

associated coPDC deposits.

ive database:

ssuming a deposit density of  $1000 \text{ kg/m}^3$

DRE assuming a solid density of 2500 (Mount St. Helens), 2600 (Tungurahua<sup>1</sup> and Moi

ntserrat<sup>y</sup>) and 2200 kg/m<sup>3</sup> (Unzen<sup>x</sup>)